

# TBA

Patrick Fox



Patrick Fox



# Tim By

# Approximation

Patrick Fox



Patrick Fox

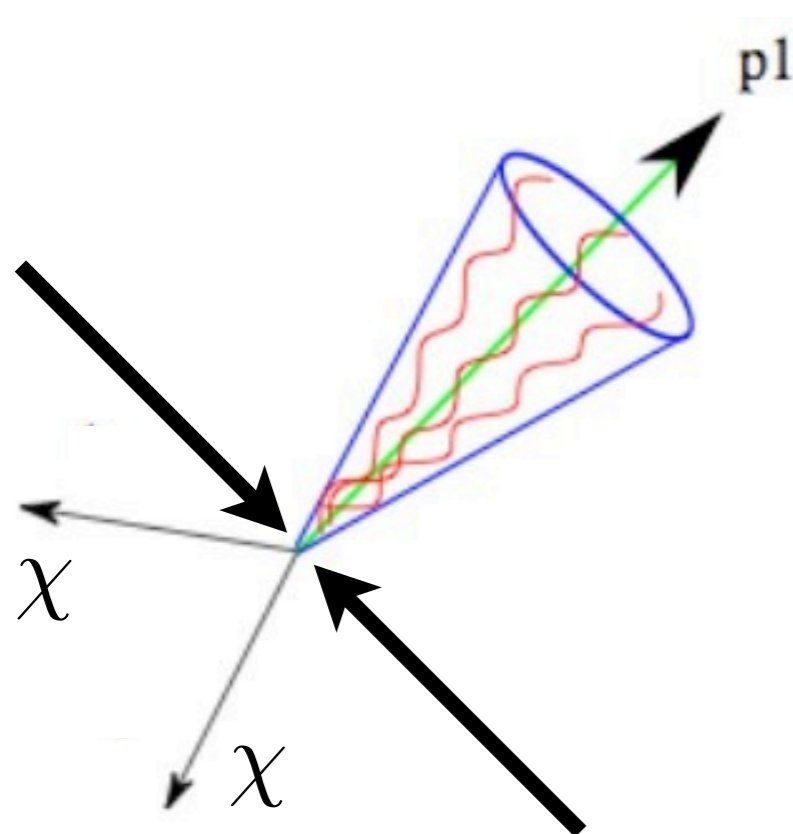




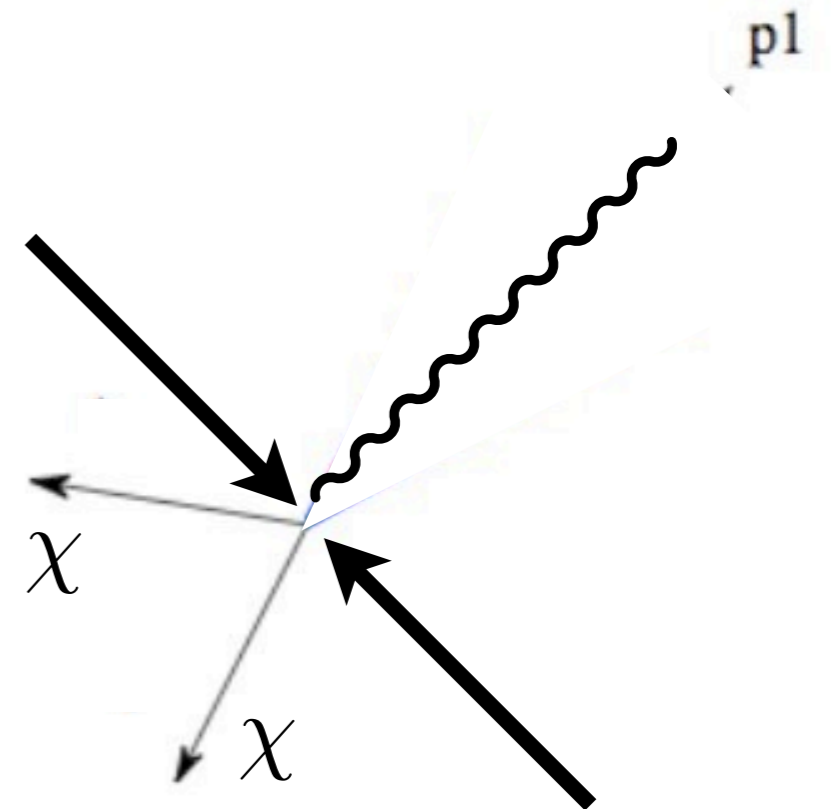
Patrick Fox



# Hunting for Dark Matter at Colliders



Patrick Fox  
 **Fermilab**

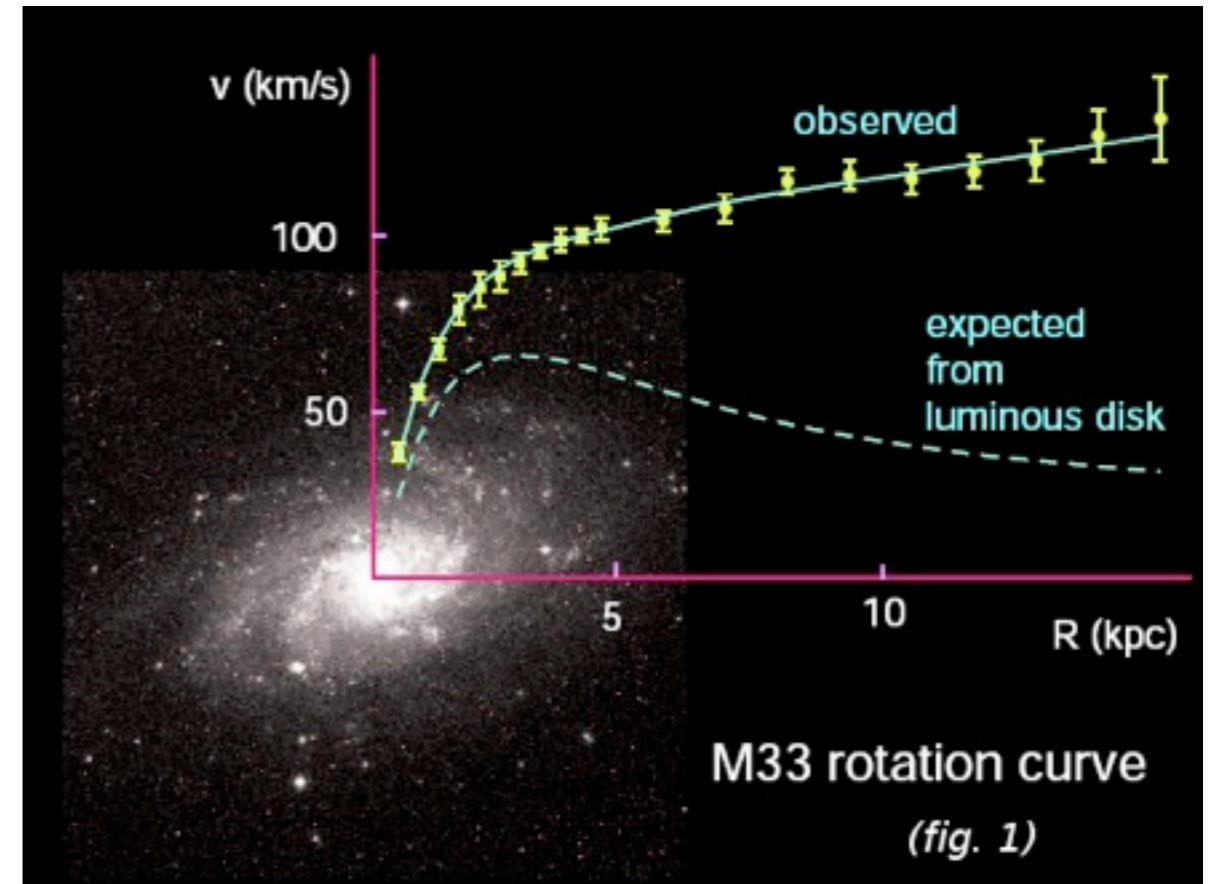
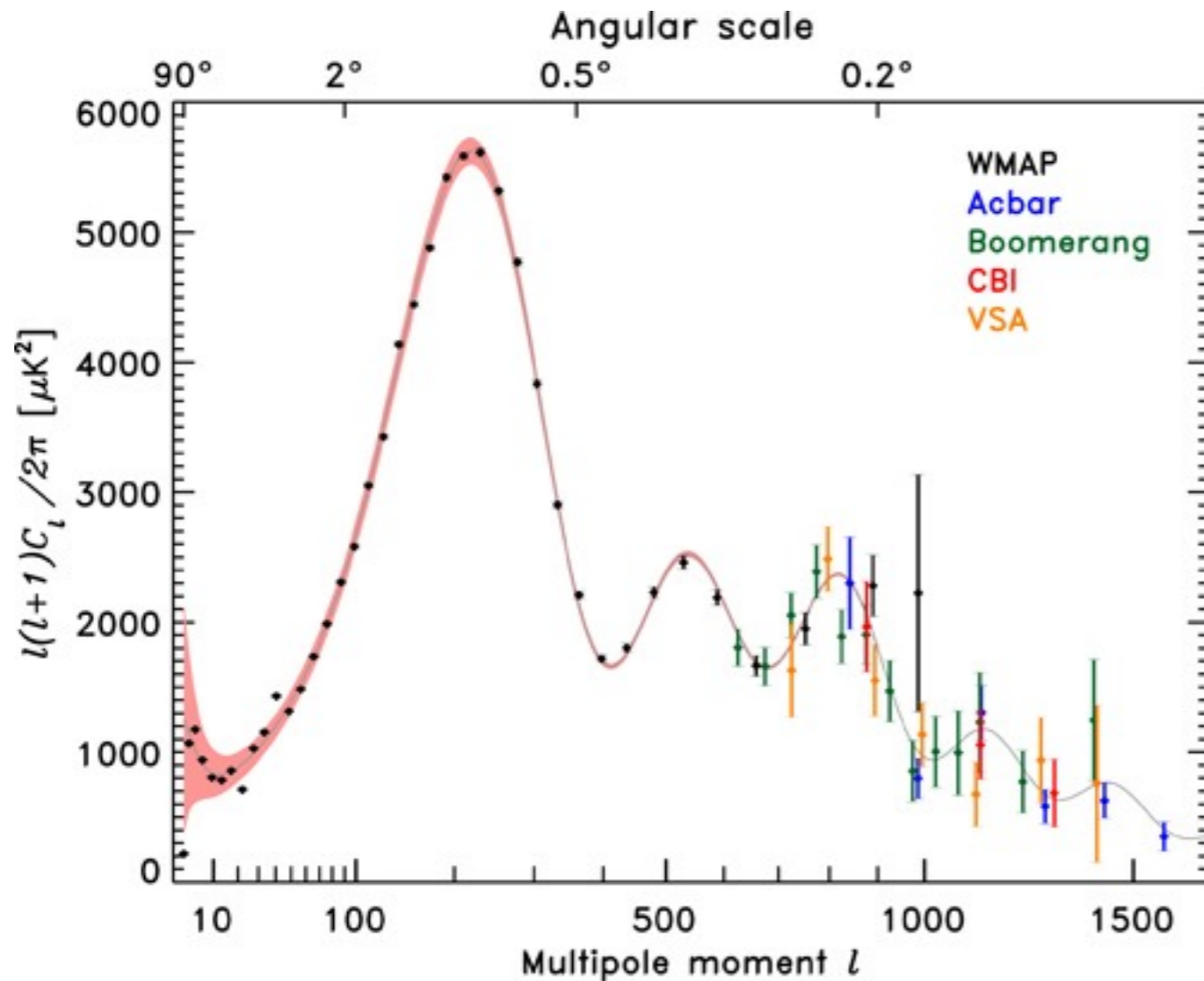


with Yang Bai and Roni Harnik  
(arXiv:1005.3797)

with Roni Harnik, Joachim Kopp and  
Yuhsin Tsai  
(to appear)

# Dark Matter

Lots of evidence for non-baryonic matter:



Cosmological abundance

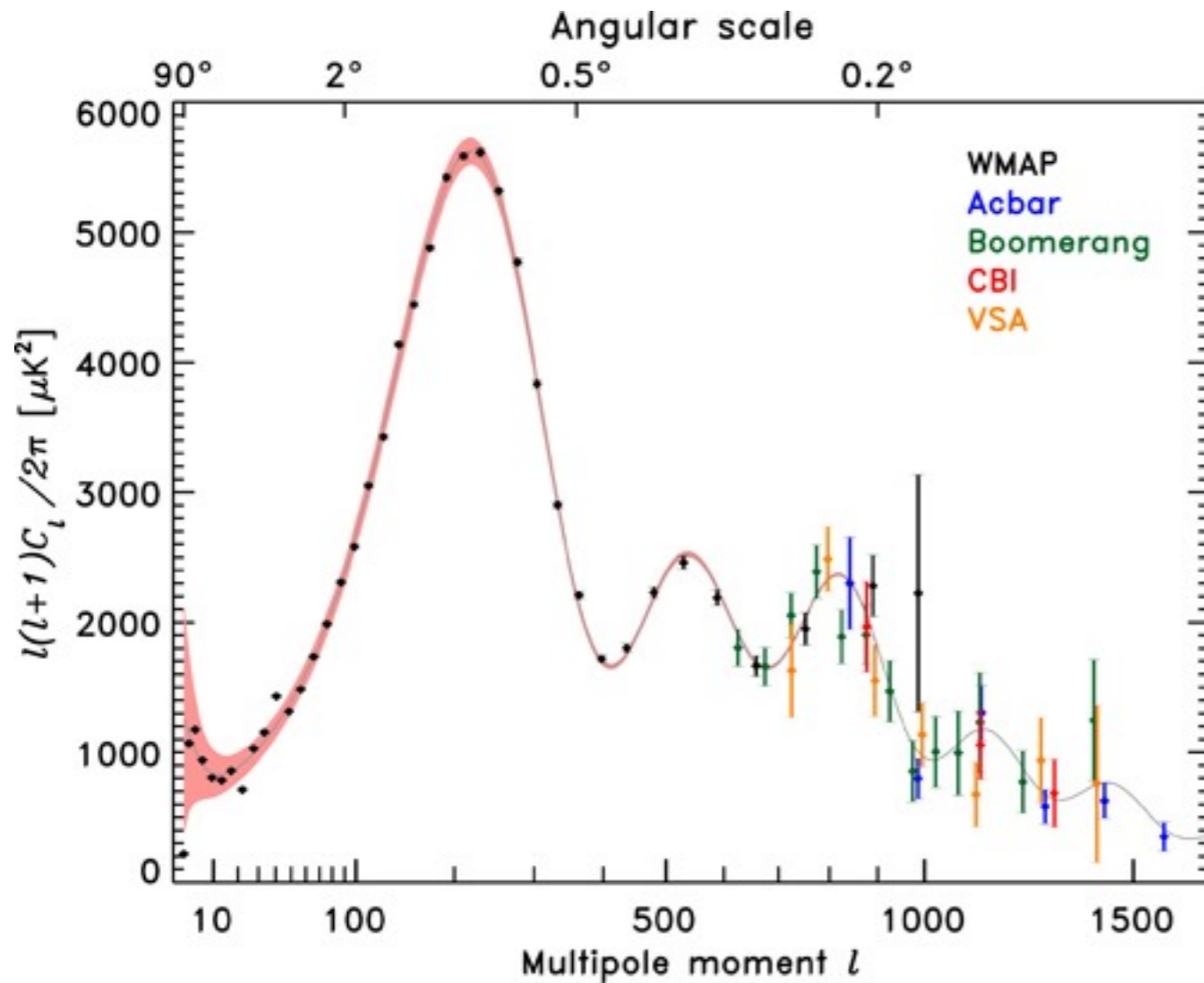
$$\Omega_{DM} = 0.213$$

Local abundance\*

$$\rho_{DM} \sim 0.3 \text{ GeV cm}^{-3}$$

# Dark Matter

Lots of evidence for non-baryonic matter:



Cosmological abundance

$$\Omega_{DM} = 0.213$$

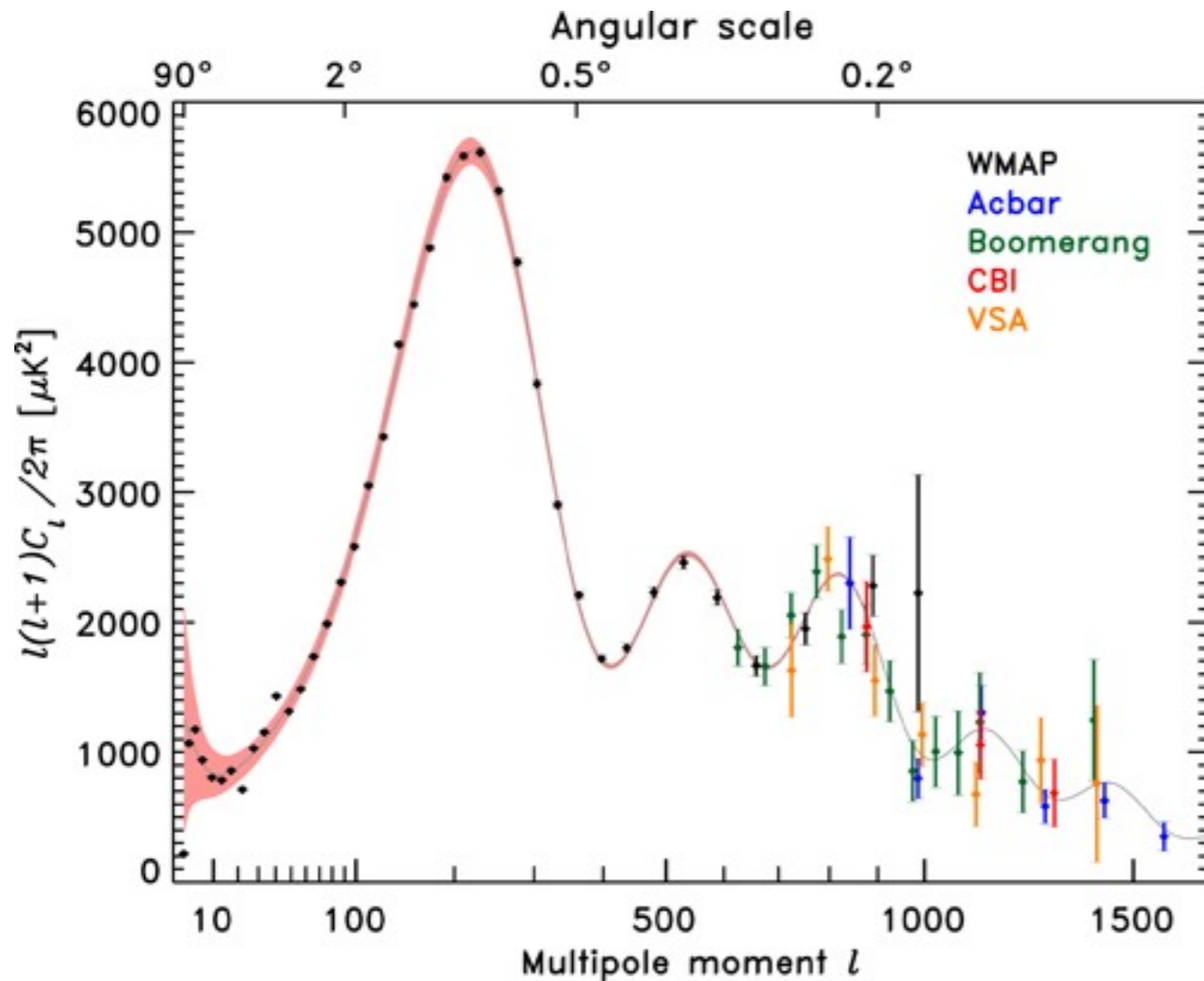
Local abundance\*

$$\rho_{DM} \sim 0.3 \text{ GeV cm}^{-3}$$



# Dark Matter

Lots of evidence for non-baryonic matter:



Cosmological abundance

$$\Omega_{DM} = 0.213$$

Local abundance\*

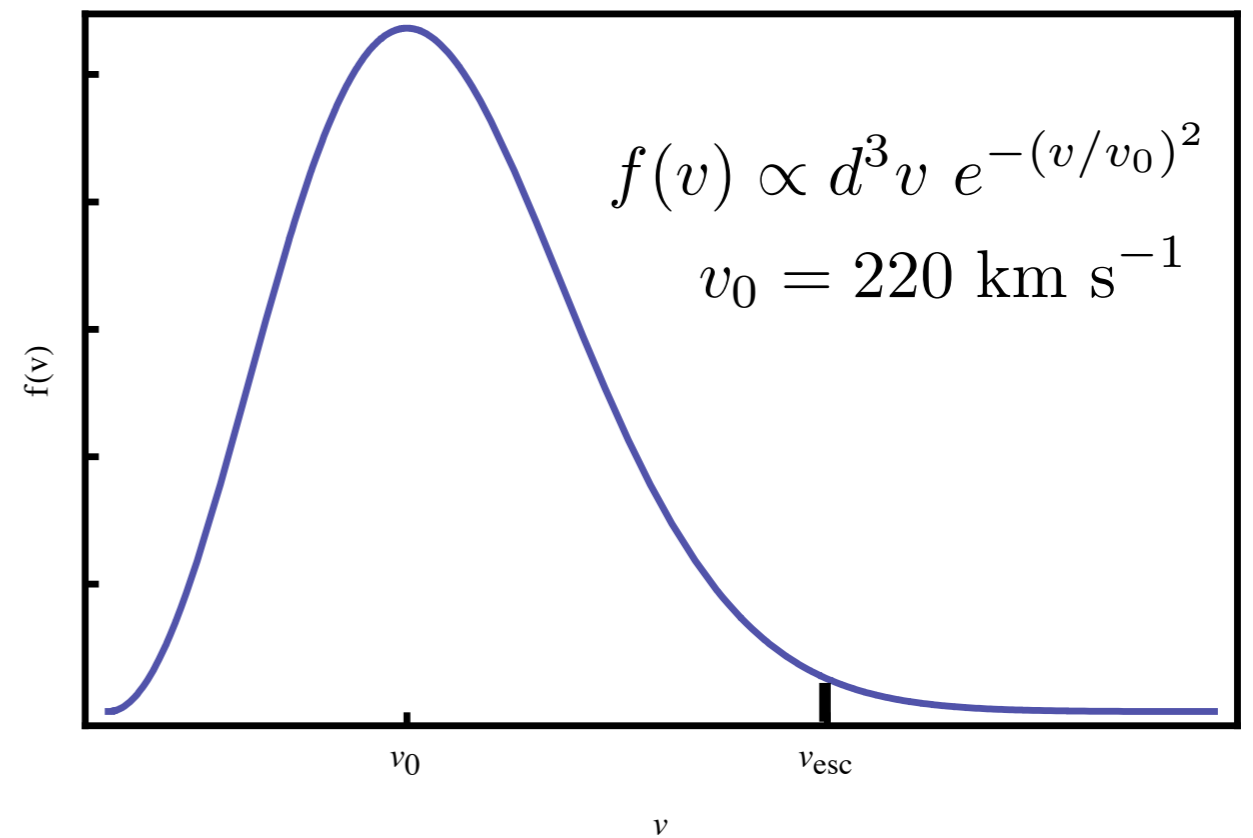
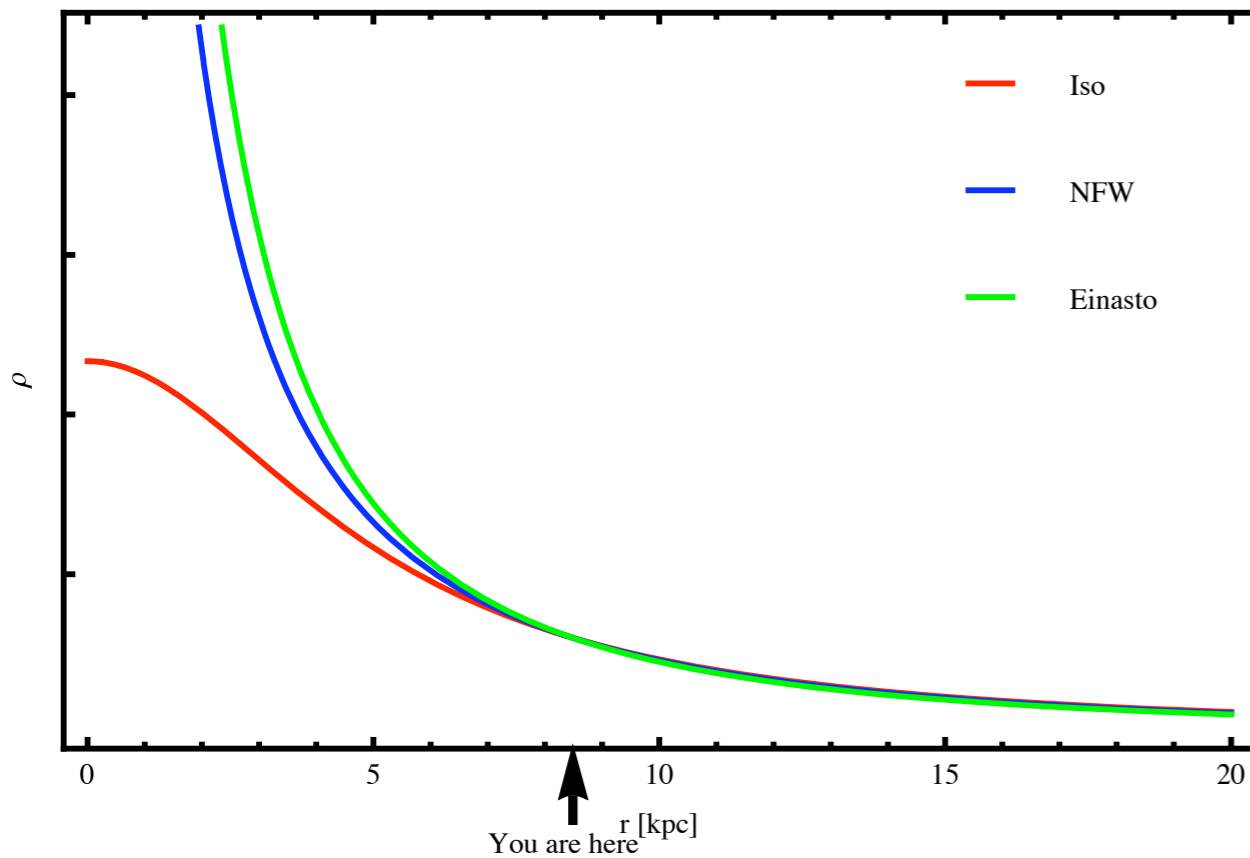
$$\rho_{DM} \sim 0.3 \text{ GeV cm}^{-3}$$

\*  $\pm$  a factor of two

# Dark Matter

Near us:  $\rho_{DM} \sim 0.3 \text{ GeV cm}^{-3}$

## Maxwell-Boltzmann velocity distribution



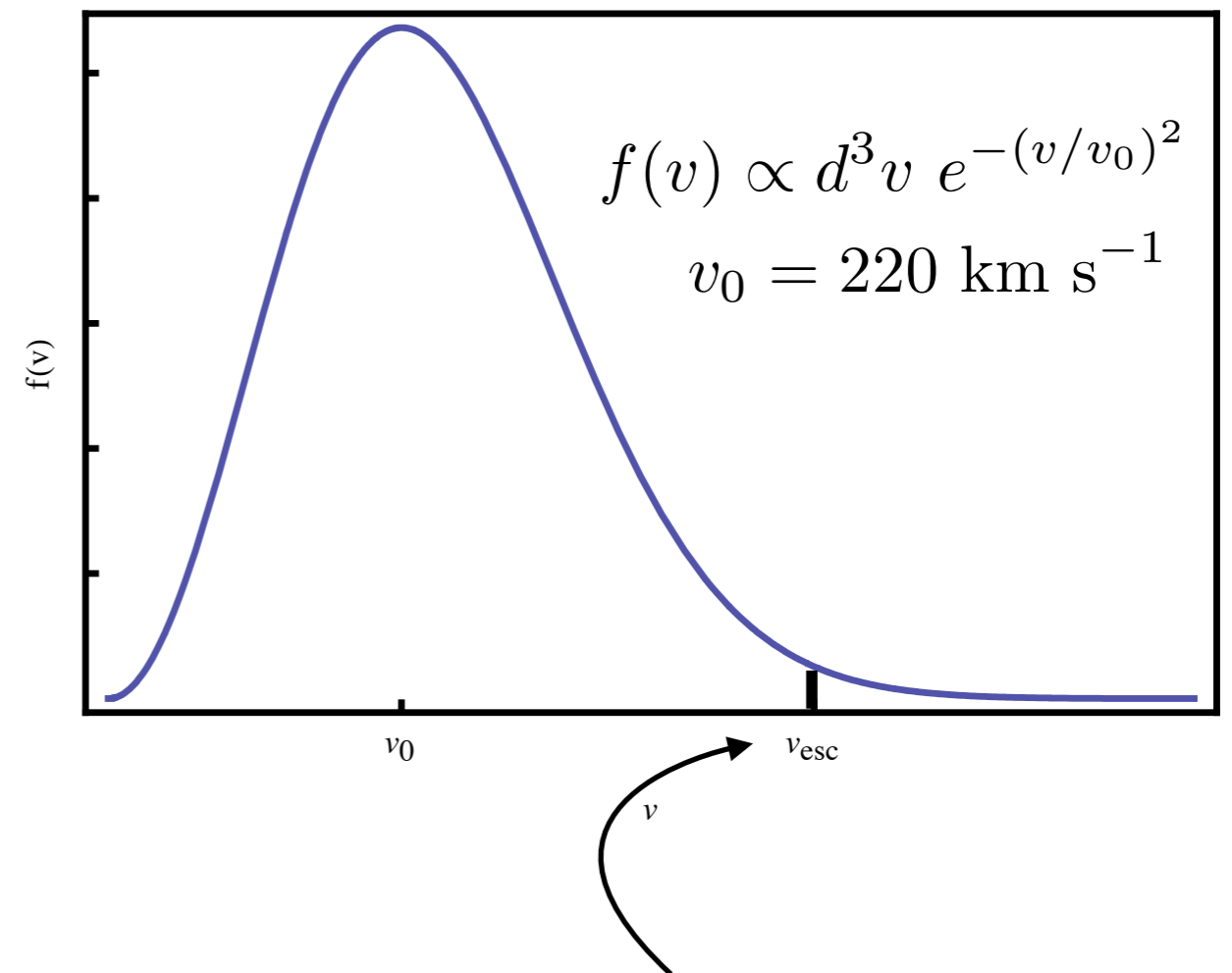
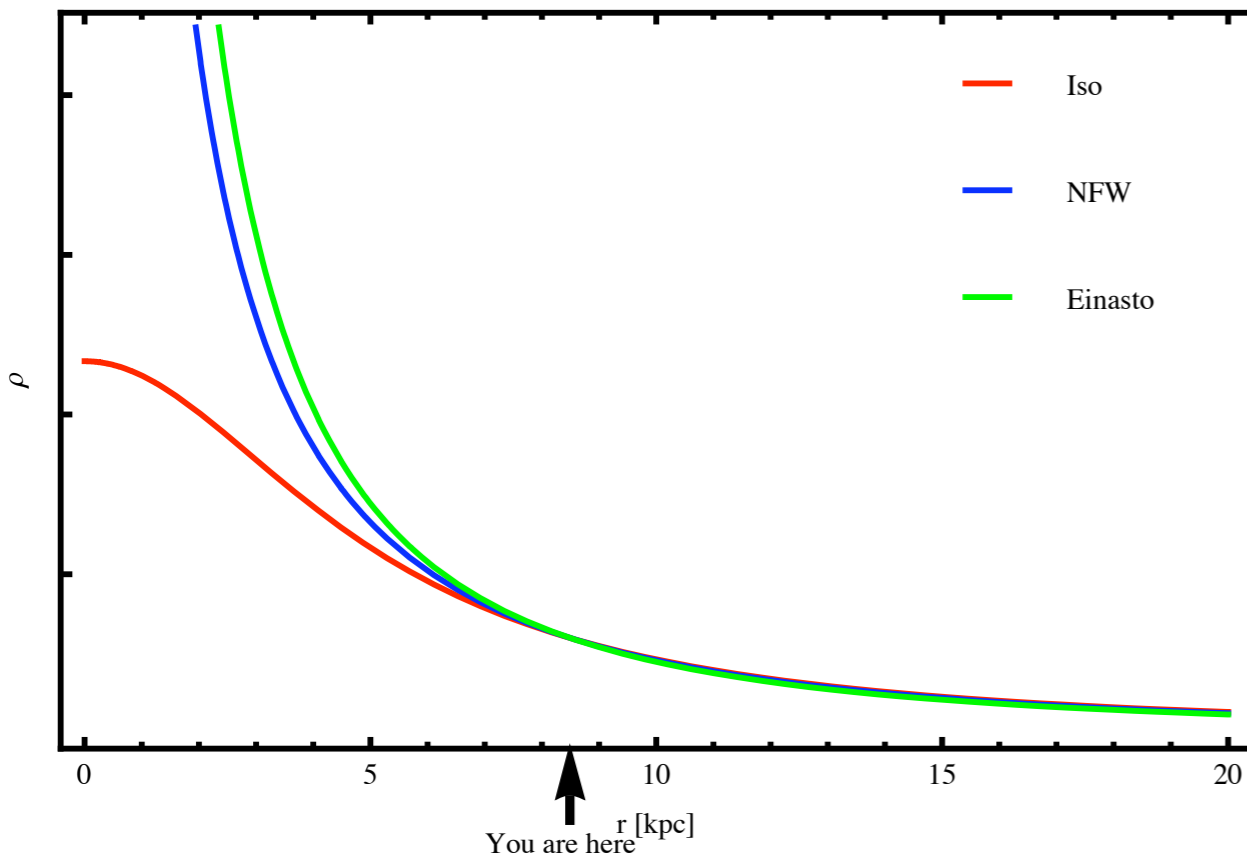
Escape velocity in galactic frame  $498 \text{ km/s} \leq v_{esc} \leq 608$

$$f(v) = \frac{1}{(\pi v_0^2)^{3/2}} e^{-v^2/v_0^2}$$

# Dark Matter

Near us:  $\rho_{DM} \sim 0.3 \text{ GeV cm}^{-3}$

## Maxwell-Boltzmann velocity distribution

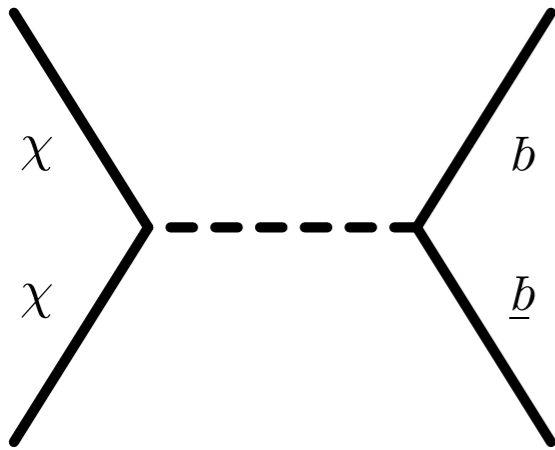


Escape velocity in galactic frame  $498 \text{ km/s} \leq v_{esc} \leq 608$

$$f(v) = \frac{1}{(\pi v_0^2)^{3/2}} e^{-v^2/v_0^2}$$

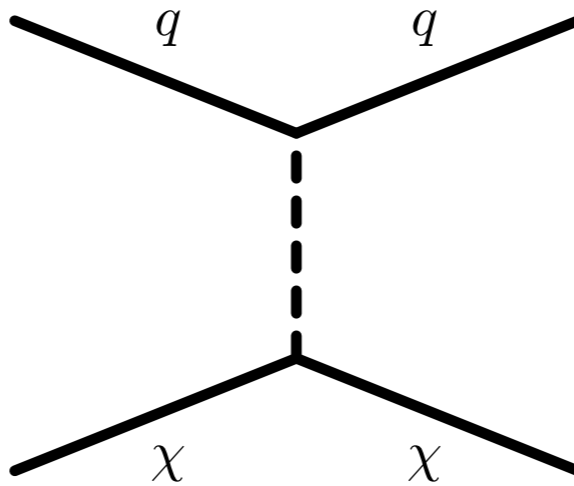
# Searching for dark matter

(here, there and everywhere)



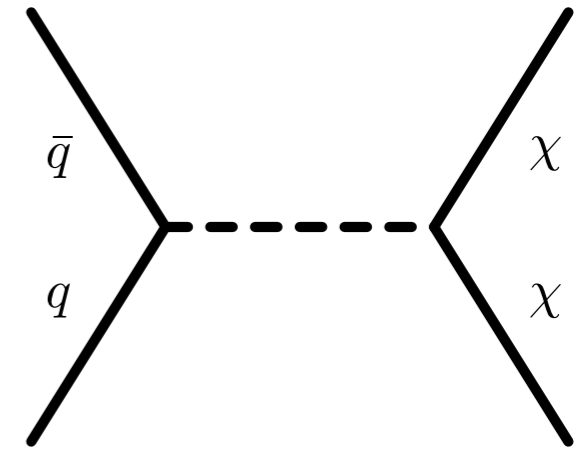
Indirect detection

Look up  
Anti-matter  
excesses in  
cosmic rays,  
photons from  
centre of galaxy



Direct detection

Look down  
Low rate, low  
energy recoil  
events in  
underground  
labs

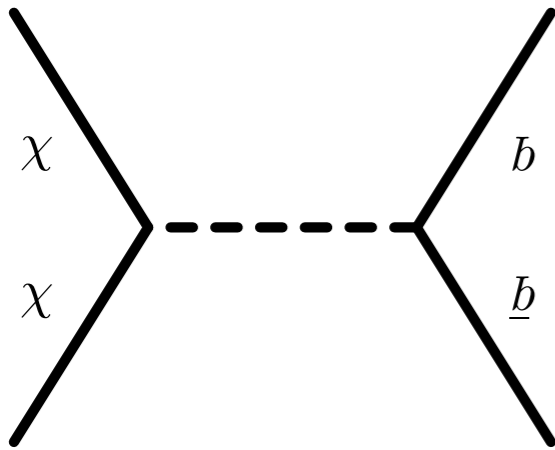


Collider searches

Look small  
Missing energy  
events at  
colliders

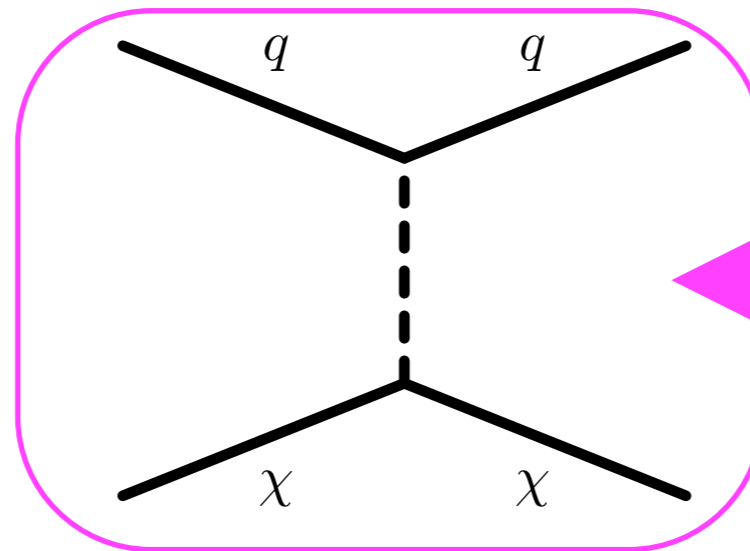
# Searching for dark matter

(here, there and everywhere)



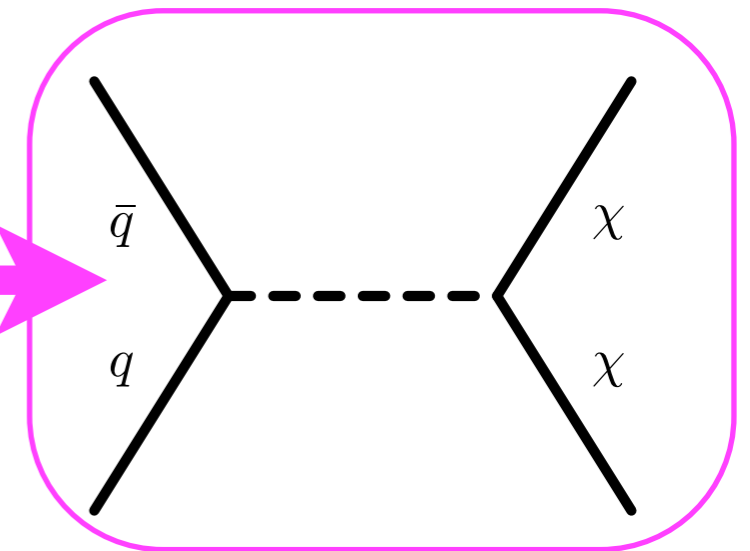
Indirect detection

Look up  
Anti-matter  
excesses in  
cosmic rays,  
photons from  
centre of galaxy



Direct detection

Look down  
Low rate, low  
energy recoil  
events in  
underground  
labs

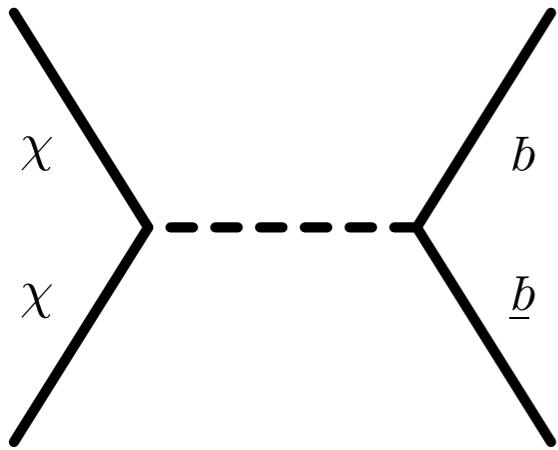


Collider searches

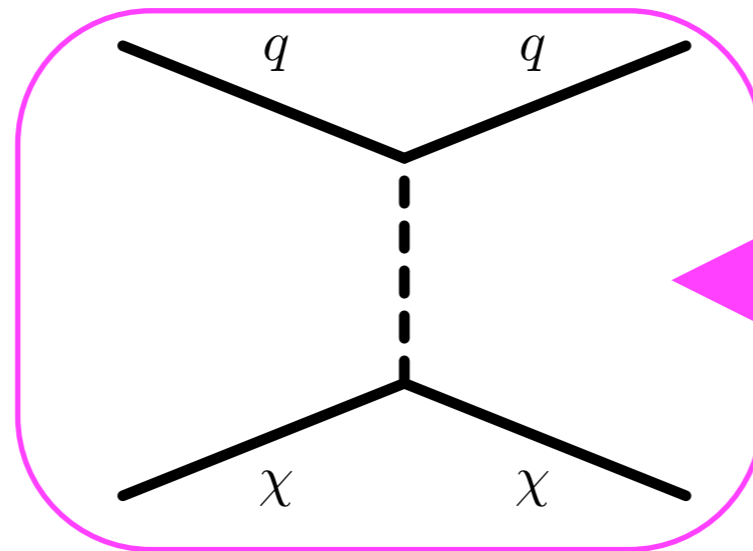
Look small  
Missing energy  
events at  
colliders

# Searching for dark matter

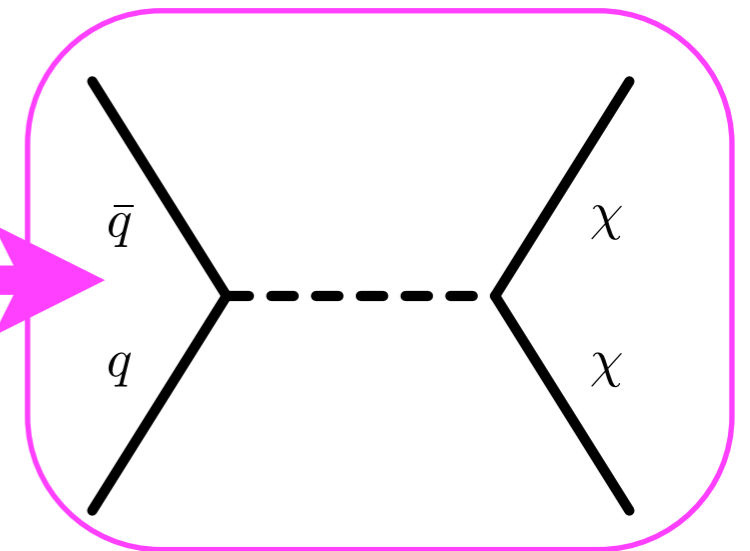
(here, there and everywhere)



Indirect detection



Direct detection



Collider searches

Look up

Anti-matter  
excesses in  
cosmic rays,  
photons from  
centre of galaxy

Look down

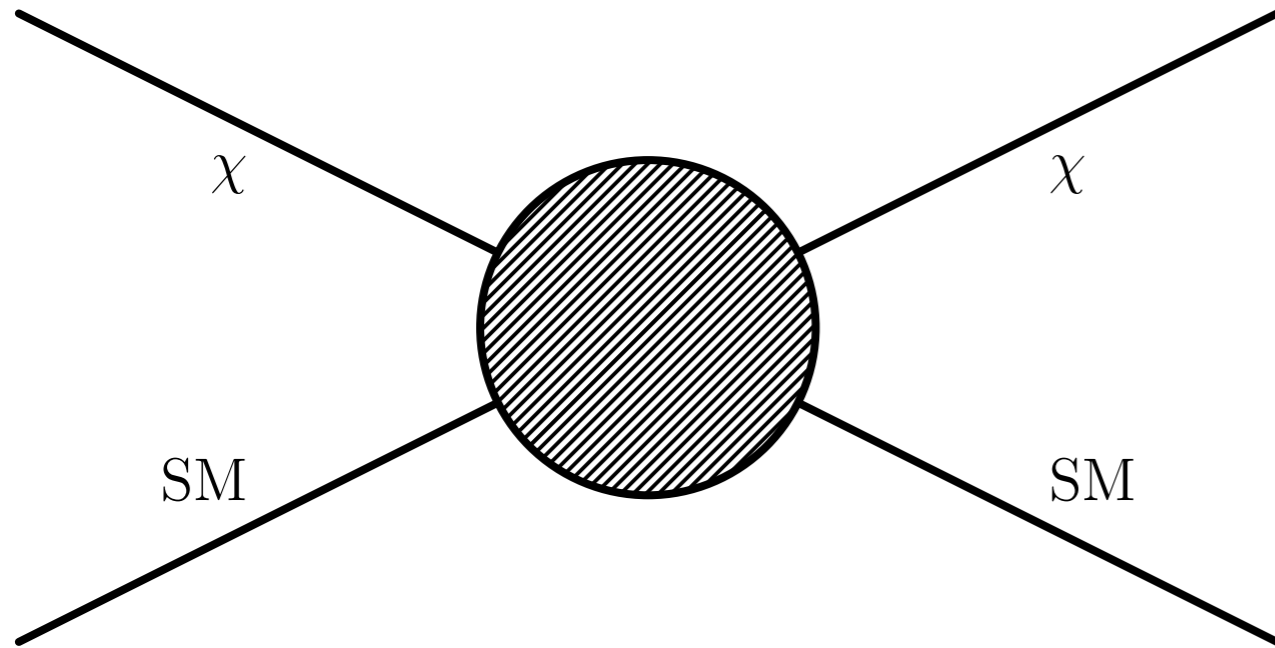
Low rate, low  
energy recoil  
events in  
underground  
labs

Look small

Missing energy  
events at  
colliders

← Thermal relic? Predicts  $\sigma\text{-sec} \sim 1 \text{ pb}$

# Direct Detection



$$E_R \sim \frac{q_\chi^2}{2 M_T} \sim 100 \text{ keV}$$

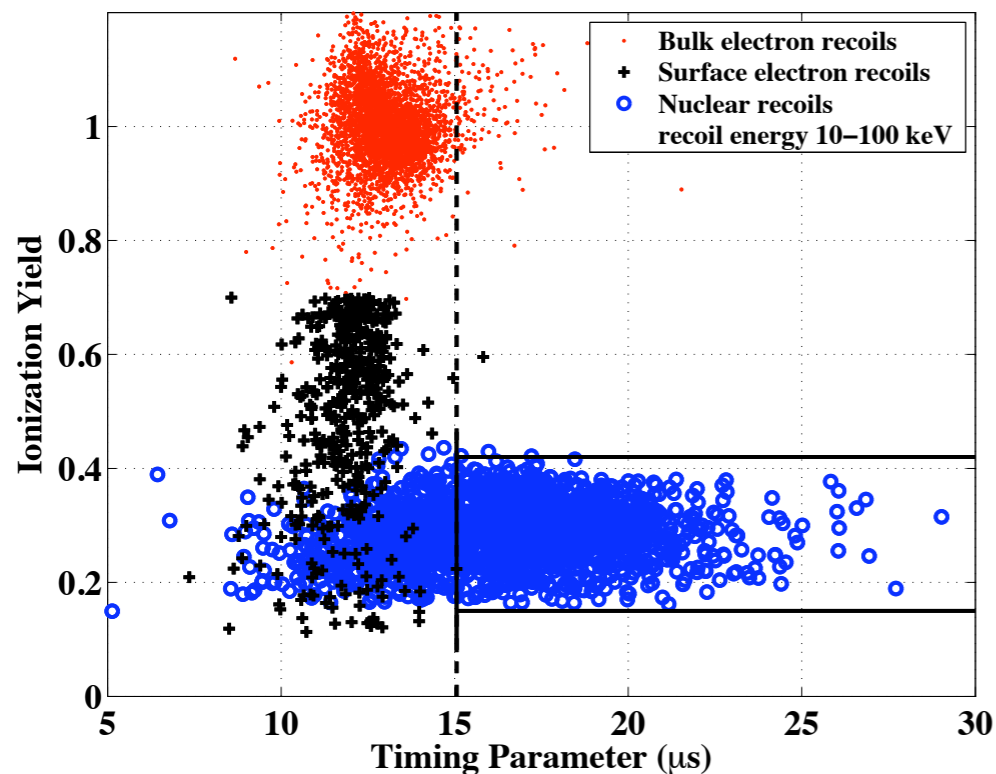
$$R \sim N_T \frac{\rho_\chi}{m_\chi} \langle \sigma v \rangle \approx 1 \text{ event/day/kg}$$

How to distinguish this small number of low energy events from backgrounds?

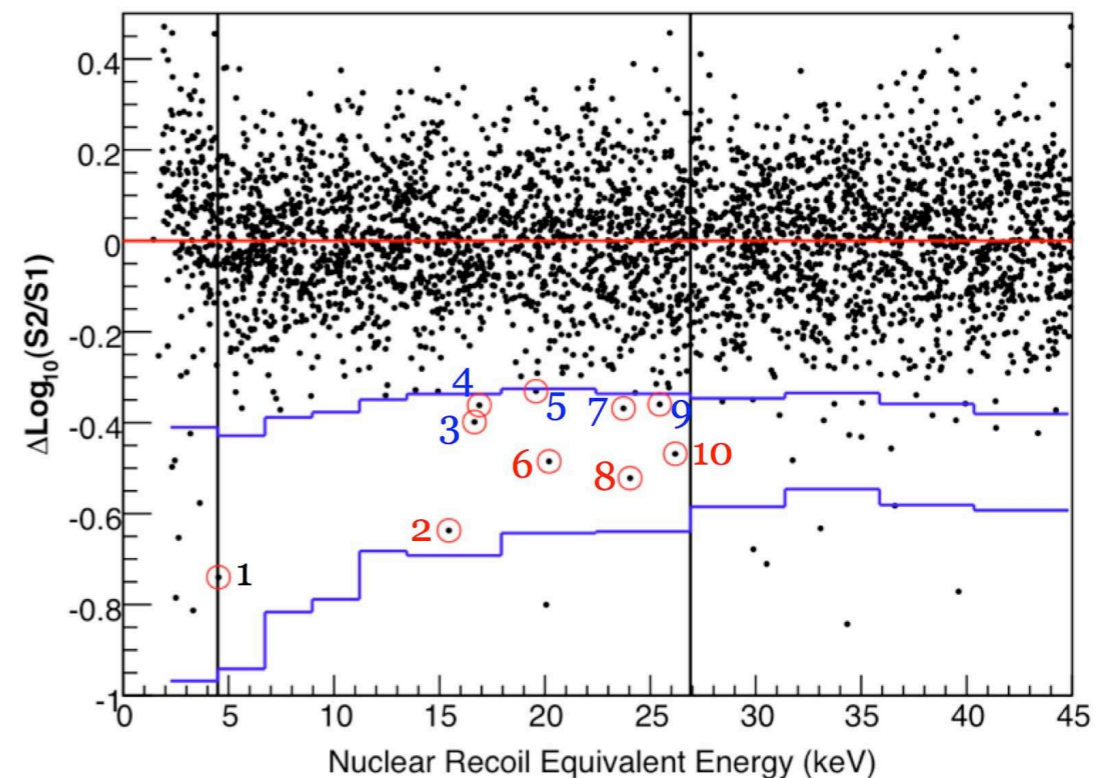
# Direct Detection

One Way:

- Remove cosmic backgrounds by going underground
- Shield experiment from radioactive elements
- Cool equipment
- Take multiple measurements to distinguish background from nuclear recoils e.g. ionization, scintillation, phonons



[CDMS collaboration]

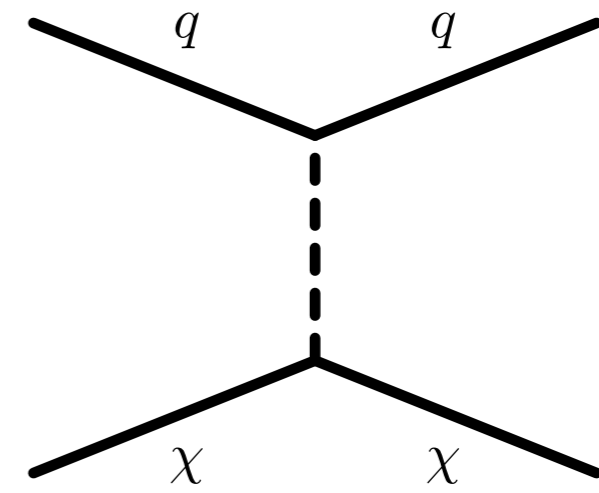


[XENON10 collaboration]

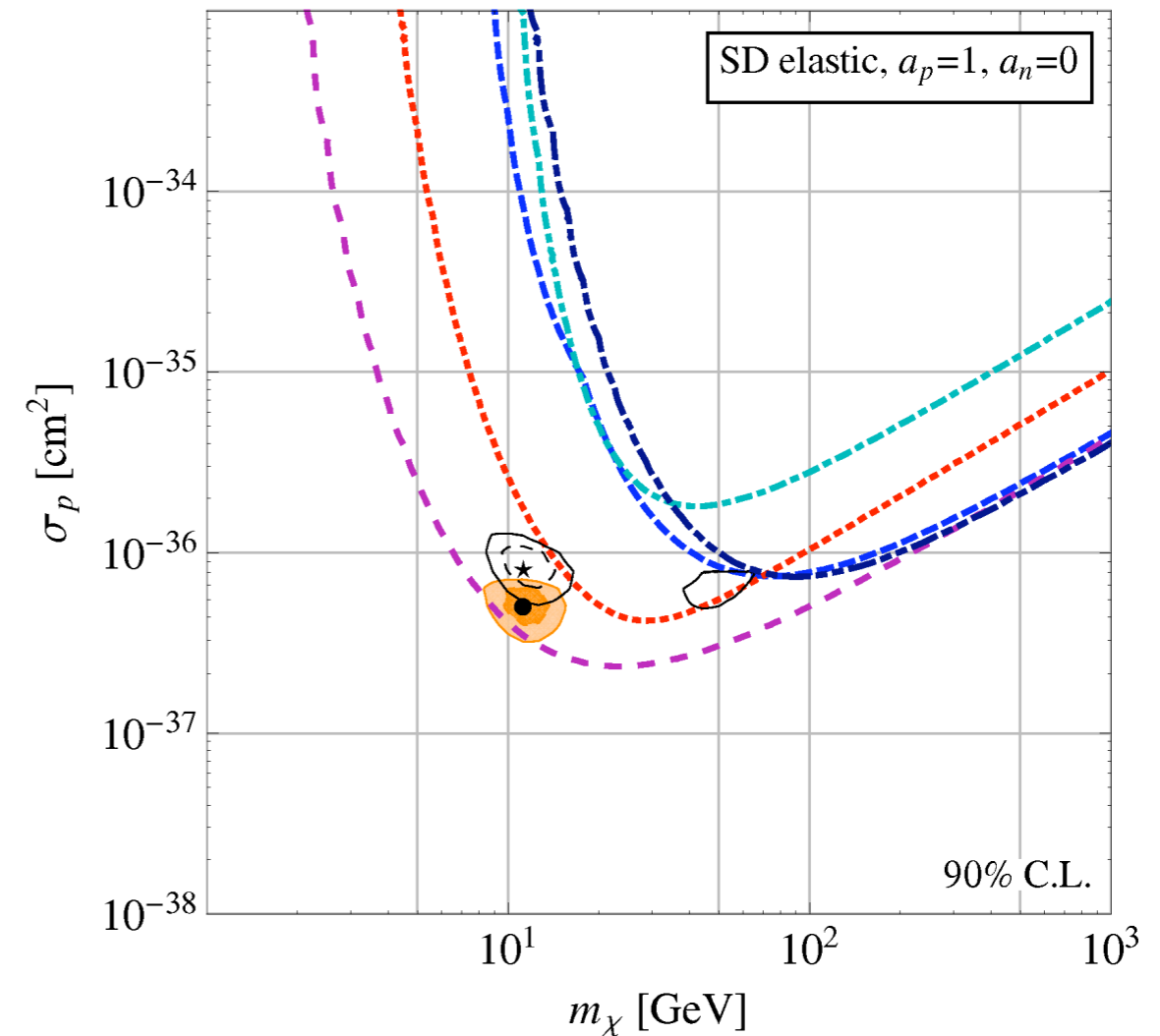
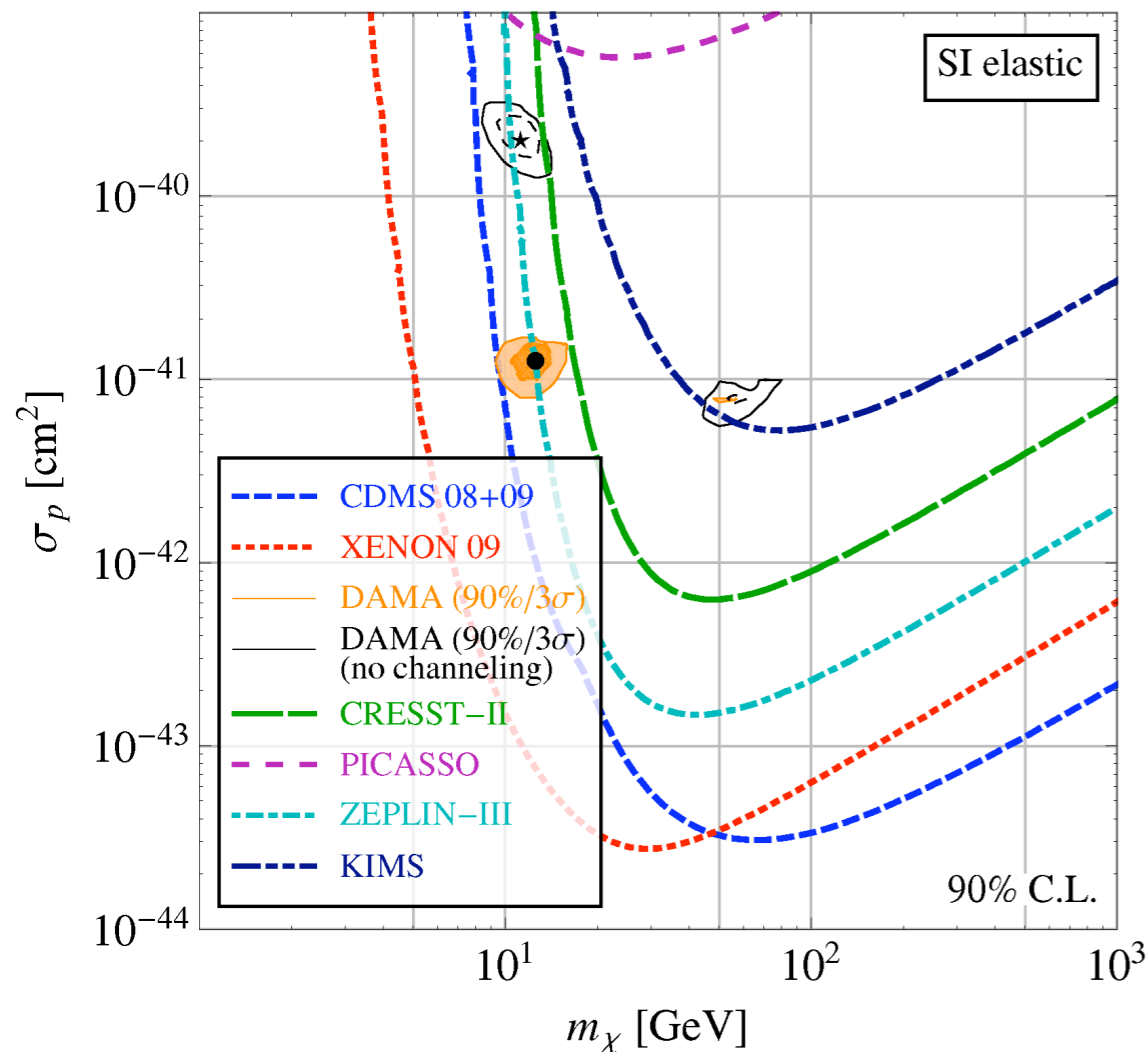


# Existing DD bounds

CDMS, XENON, DAMA,  
CoGeNT, COUPP,  
CRESST, .....

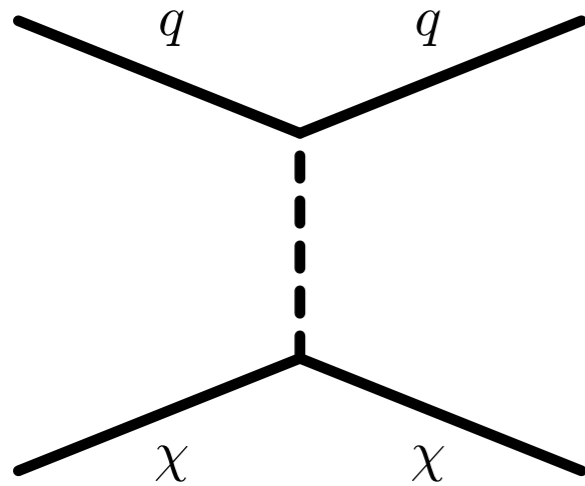


[Kopp, Schwetz and Zupan]



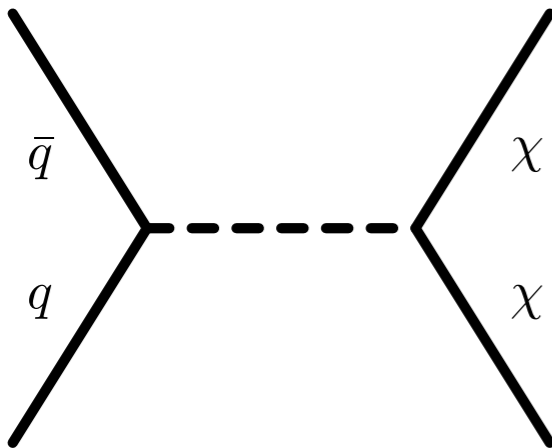
(Assume local abundance is  $0.3 \text{ GeV/cm}^3$ )

# Direct detection vs Collider production



Direct detection

$$q \sim 100 \text{ MeV}$$



Collider searches

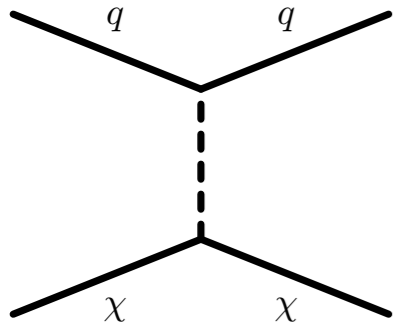
$$q \sim 10 - 100 \text{ GeV}$$

How does one search impact the other?

[Birkedal, Matchev and Perelstein]

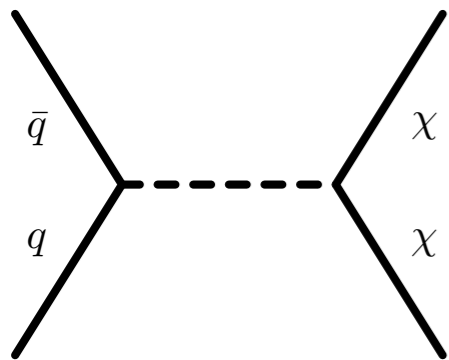
# Mediator Mass dependence

Only consider mediators with mass  $\gtrsim 100$  MeV



$$\sigma_{DD} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\mu = \frac{m_{\chi} m_N}{m_N + m_{\chi}}$$



**Mono-jet +  $\cancel{E}_T$**

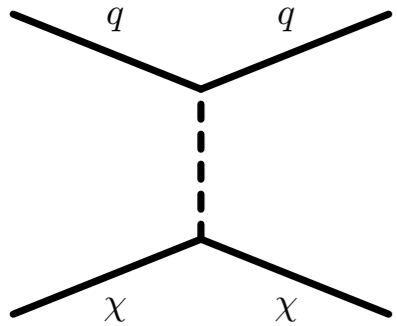
$$\sigma_{1j} \sim \begin{cases} \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2} & M \lesssim 100 \text{ GeV} \\ \alpha_s g_{\chi}^2 g_q^2 \frac{p_T^2}{M^4} & M \gtrsim 100 \text{ GeV} \end{cases}$$

CDF analysed  $1 \text{ fb}^{-1}$  and saw no significant deviation

[<http://www-cdf.fnal.gov/physics/exotic/r2a/20070322.monojet/public/ykk.html>]

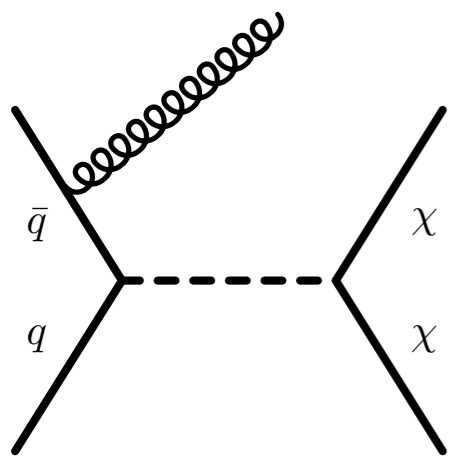
# Mediator Mass dependence

Only consider mediators with mass  $\gtrsim 100$  MeV



$$\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\mu = \frac{m_{\chi} m_N}{m_N + m_{\chi}}$$



**Mono-jet +  $\cancel{E}_T$**

$$\sigma_{1j} \sim \begin{cases} \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2} & M \lesssim 100 \text{ GeV} \\ \alpha_s g_{\chi}^2 g_q^2 \frac{p_T^2}{M^4} & M \gtrsim 100 \text{ GeV} \end{cases}$$

CDF analysed  $1 \text{ fb}^{-1}$  and saw no significant deviation

[<http://www-cdf.fnal.gov/physics/exotic/r2a/20070322.monojet/public/ykk.html>]

Consider massive mediator:

$$(p_T \sim 100 \text{ GeV})$$

$$\sigma_{1j} \sim \alpha_s g_\chi^2 g_q^2 \frac{p_T^2}{M^4}$$

$$(\mu \sim 1 \text{ GeV})$$

$$\sigma_{DD} \sim g_\chi^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\frac{\sigma_{1j}}{\sigma_{DD}} \sim \mathcal{O}(1000)$$

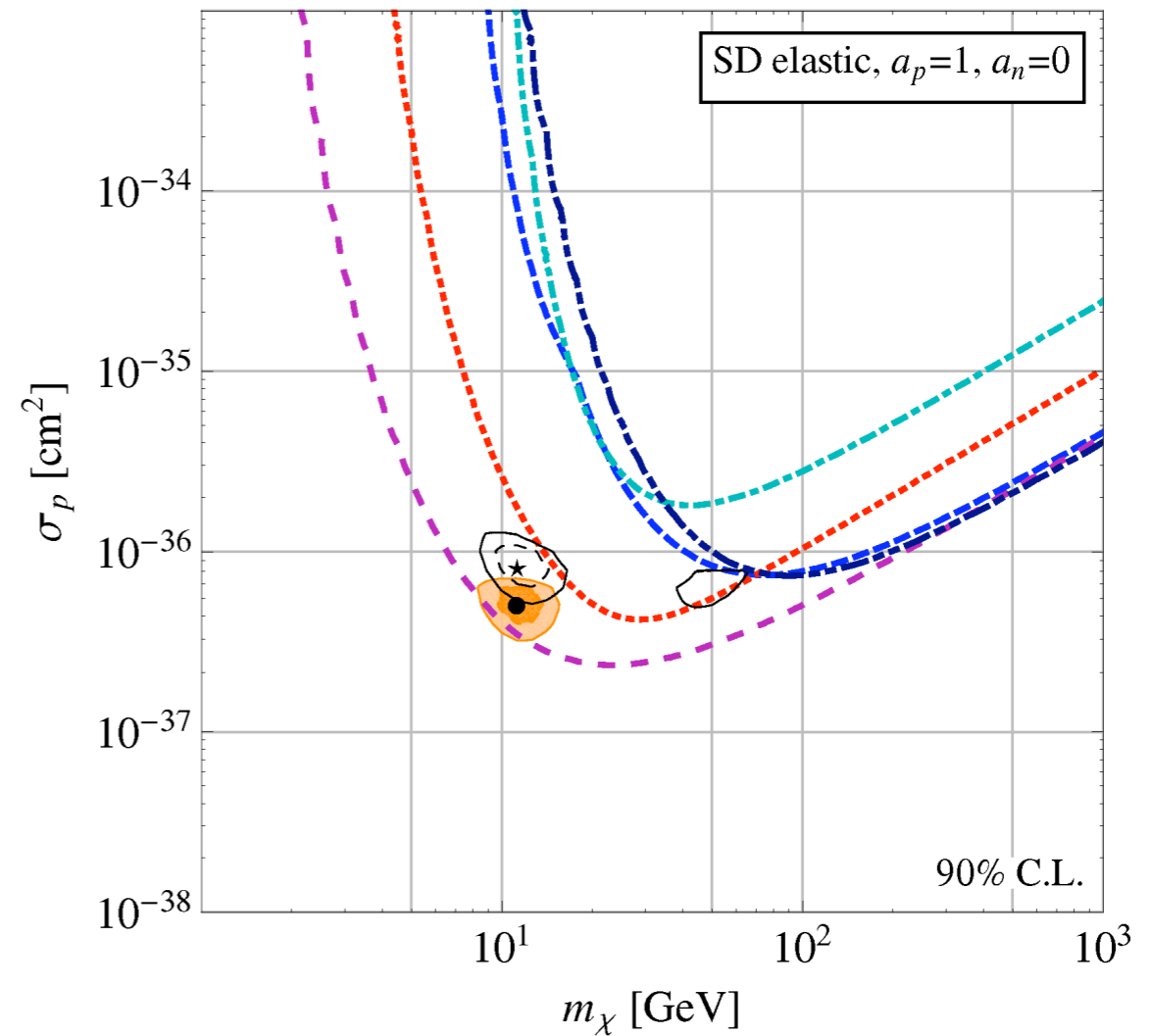
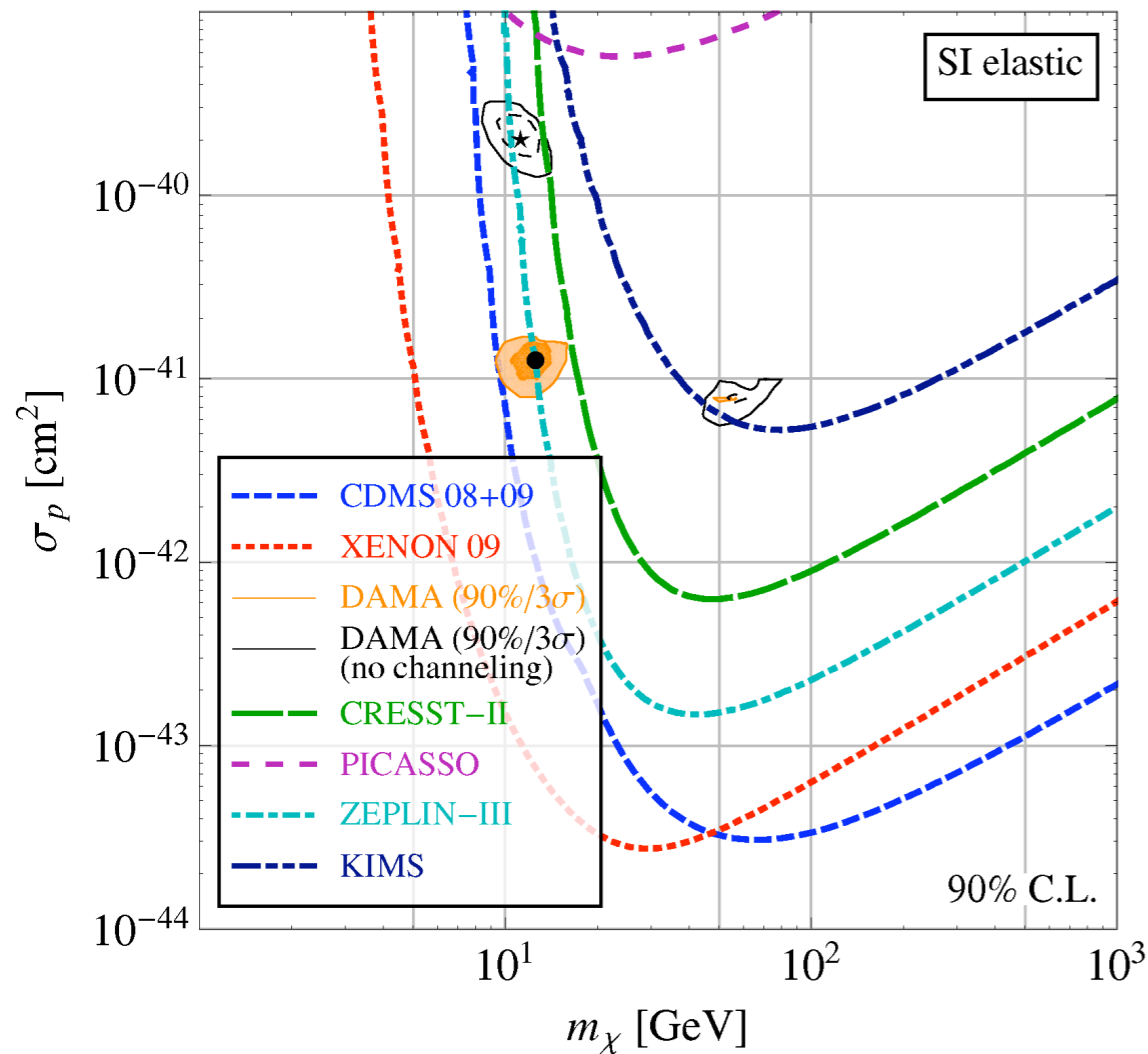
In 1 invfb CDF saw 8449 mono-jet events, expected  $8663 \pm 332$

$$\Rightarrow \sigma_{1j} \lesssim 500 \text{ fb}$$

$$\sigma_{DD} \lesssim 0.5 \text{ fb} = 5 \times 10^{-40} \text{ cm}^2$$



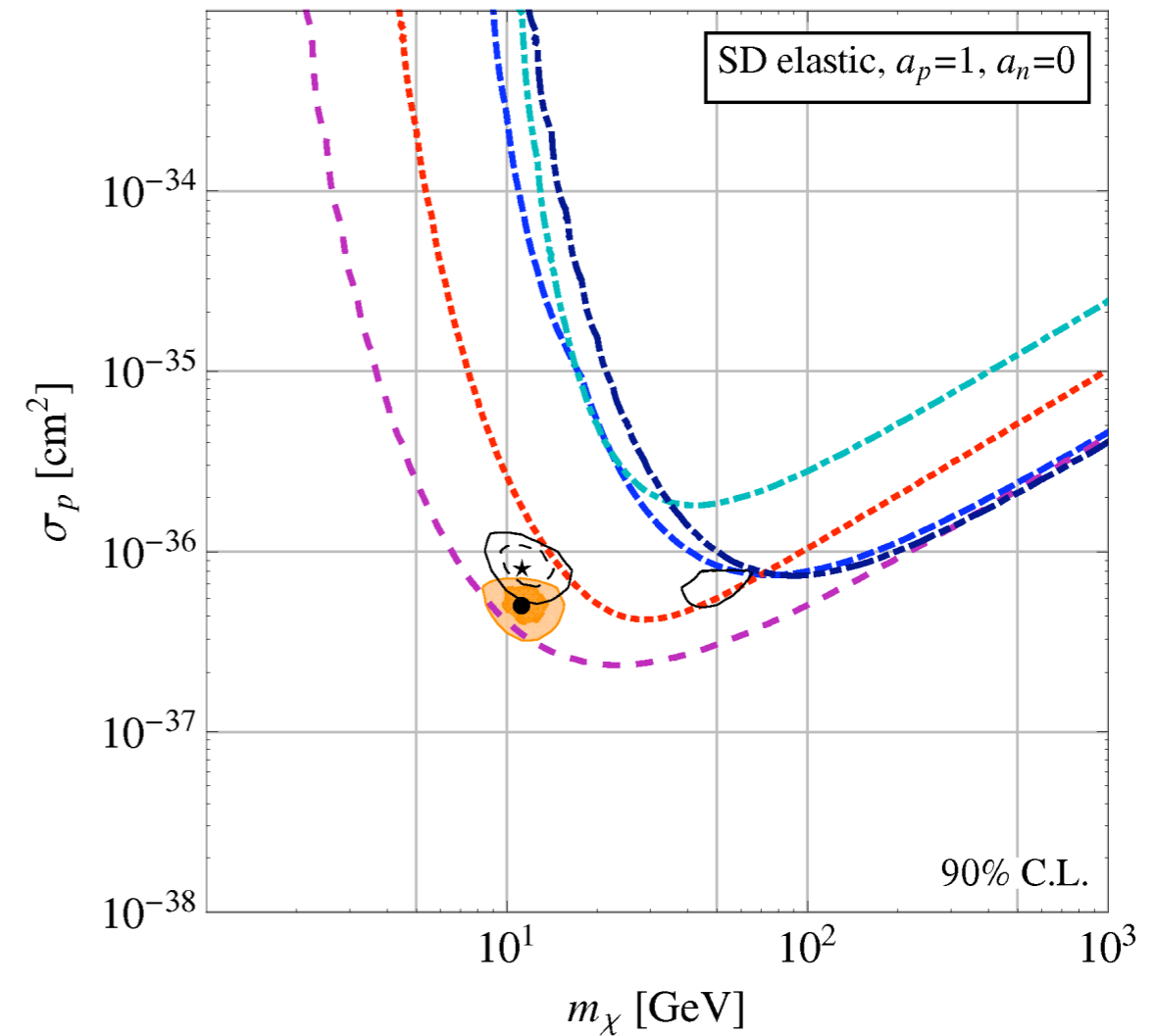
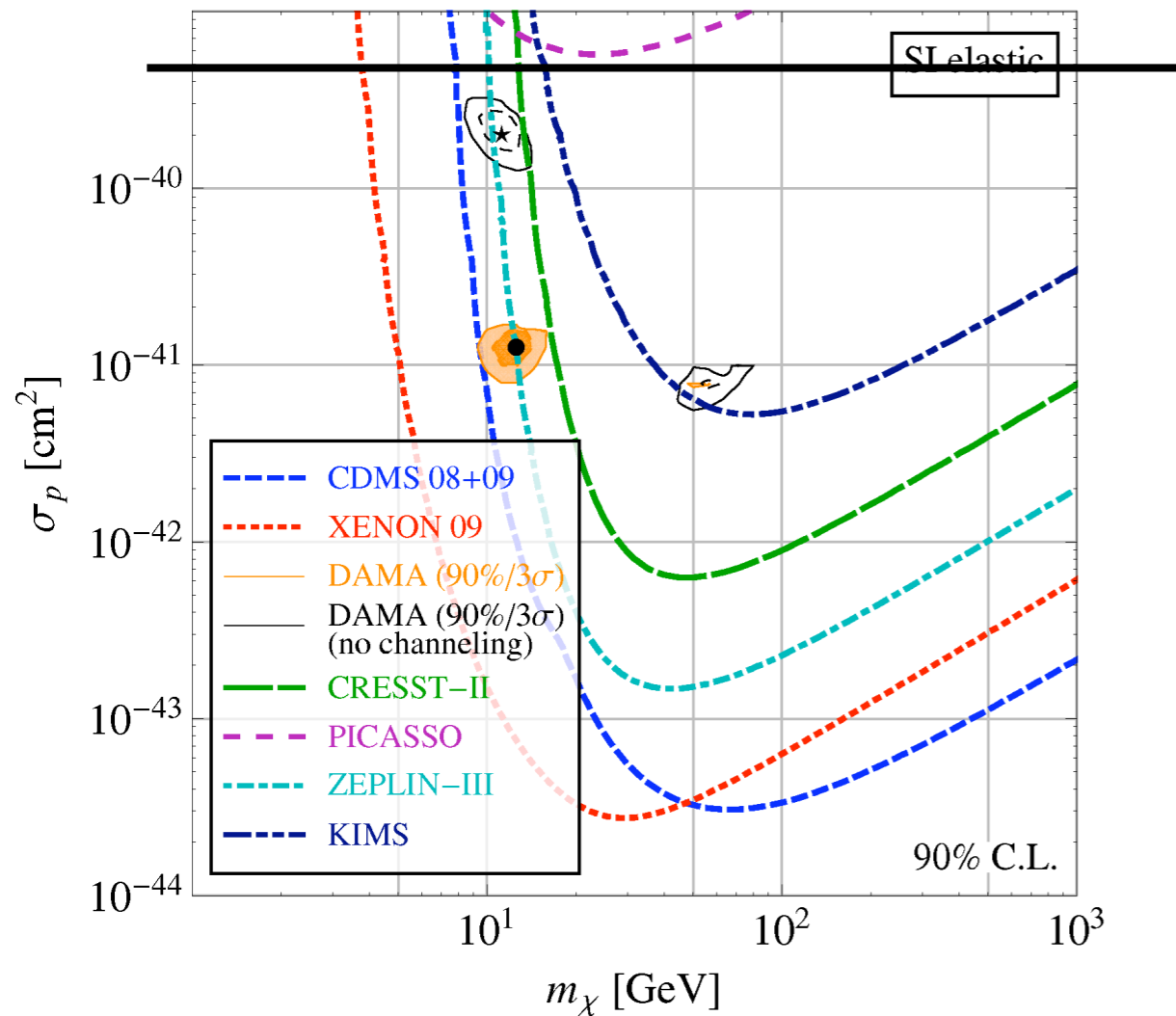
# Existing DD bounds



## ROI's

- Light mass DM
- Non-standard DM introduced to explain DAMA
- Velocity, momentum or spin suppression

# Existing DD bounds



## ROI's

- Light mass DM
- Non-standard DM introduced to explain DAMA
- Velocity, momentum or spin suppression

# Outline

- Motivation and estimation
- Operator analysis
- Heavy mediators
- Collider bounds
- Light mediators
- LEP
- Conclusions



# Outline

- ~~• Motivation and estimation~~
- Operator analysis
- Heavy mediators
- Collider bounds
- Light mediators
- LEP
- Conclusions

# Operators

$$\mathcal{O}_1 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}q) (\bar{\chi}\chi) , \quad \text{SI, scalar exchange}$$

$$\mathcal{O}_2 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}\gamma_\mu q) (\bar{\chi}\gamma^\mu \chi) , \quad \text{SI, vector exchange}$$

$$\mathcal{O}_3 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}\gamma_\mu \gamma_5 q) (\bar{\chi}\gamma^\mu \gamma_5 \chi) , \quad \text{SD, axial-vector exchange}$$

$$\mathcal{O}_4 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}\gamma_5 q) (\bar{\chi}\gamma_5 \chi) , \quad \text{SD and mom. dep., psuedo-scalar exchange}$$

- DM a Dirac fermion
- Consider each operator, and each flavour separately

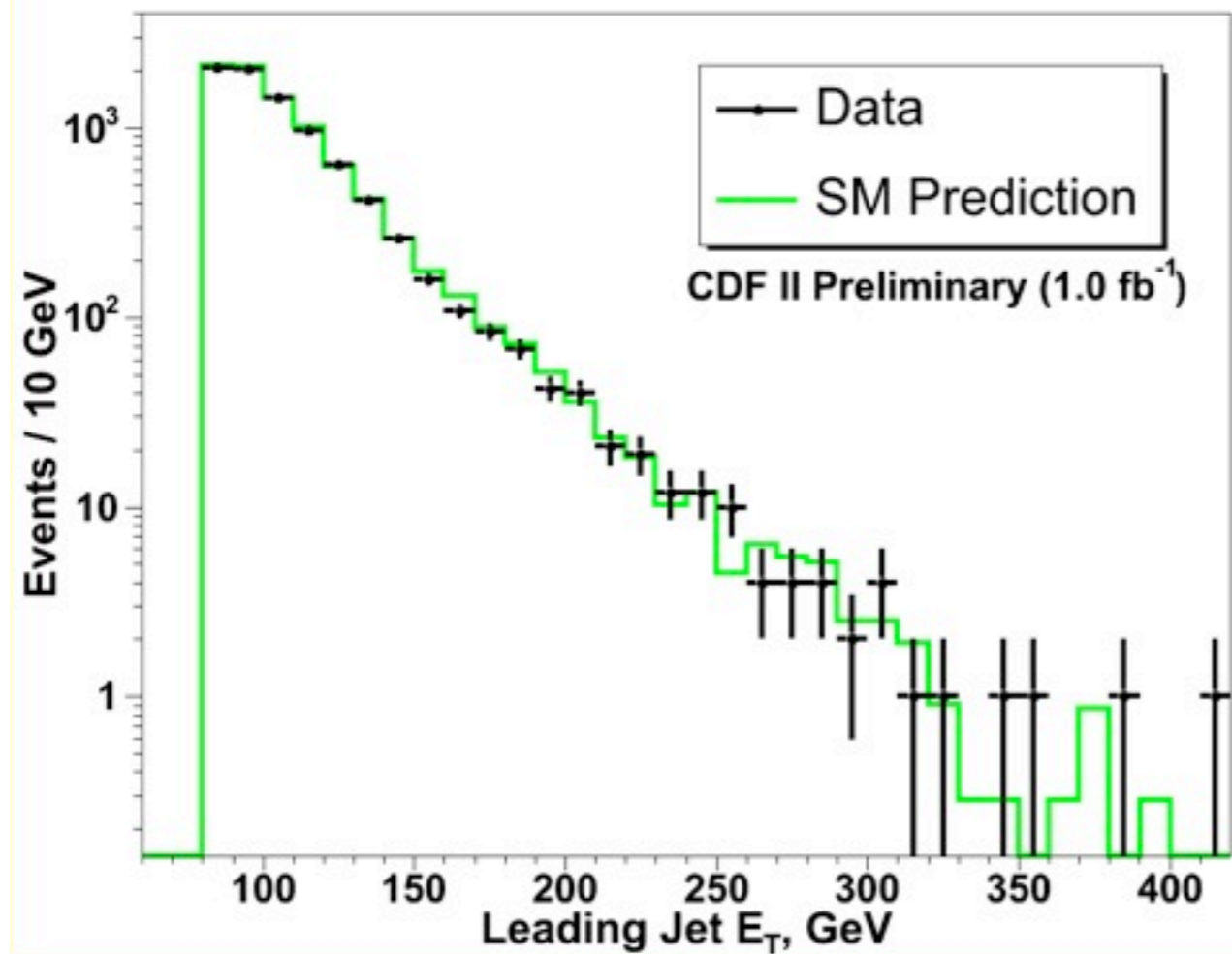
# CDF mono-jet search

[<http://www-cdf.fnal.gov/physics/exotic/r2a/20070322.monojet/public/ykk.html>]

- 1/fb analysed

- $\cancel{E}_T > 80 \text{ GeV}$
- $p_T(j1) > 80 \text{ GeV}$
- $p_T(j2) < 30 \text{ GeV}$
- $p_T(j3) < 20 \text{ GeV}$

Background	Number of Events
Z -> nu nu	3203 +/- 137
W -> tau nu	2010 +/- 69
W -> mu nu	1570 +/- 54
W -> e nu	824 +/- 28
Z -> ll	87 +/- 3
QCD	708 +/- 146
Gamma plus Jet	209 +/- 41
Non-Collision	52 +/- 52
<b>Total Predicted</b>	<b>8663 +/- 332</b>
<b>Data Observed</b>	<b>8449</b>

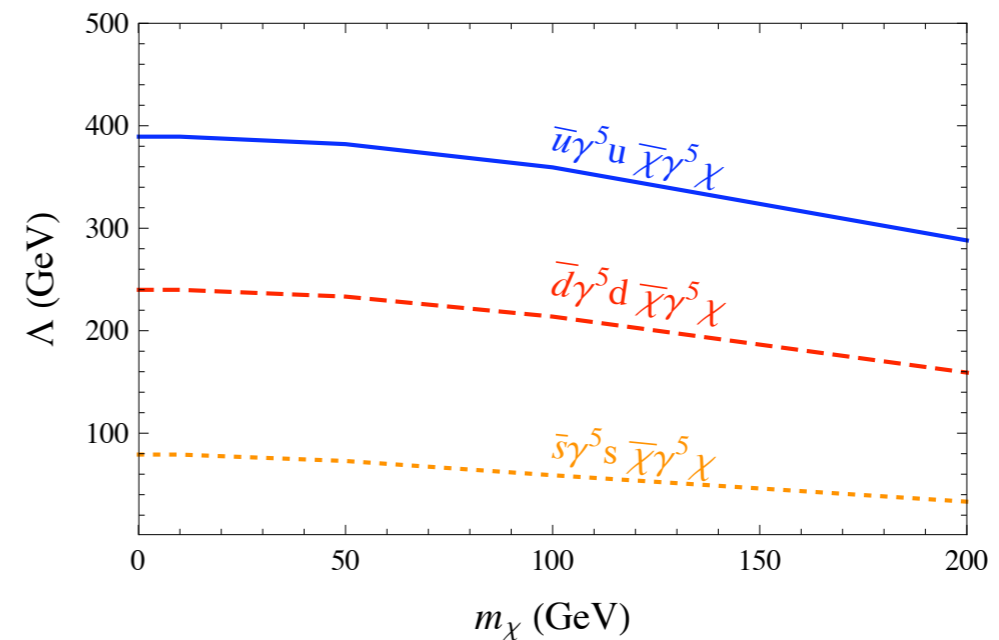
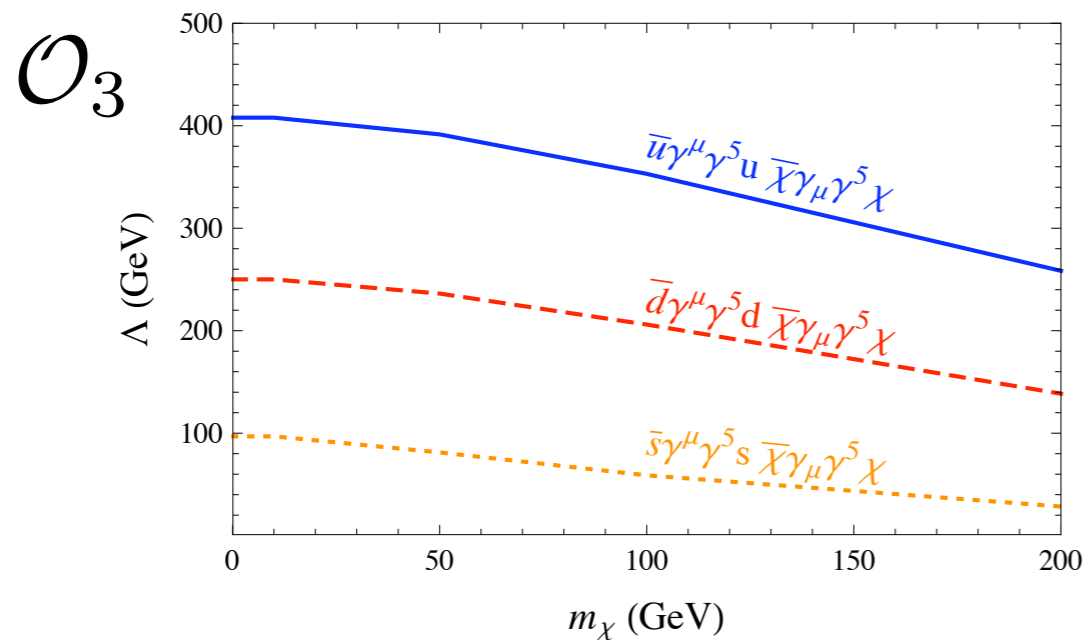
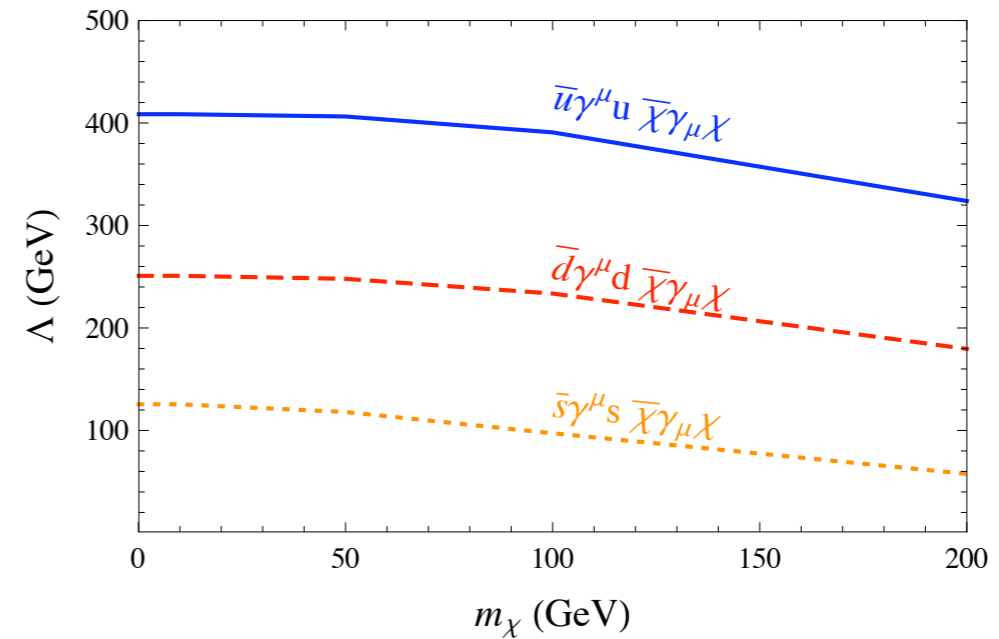
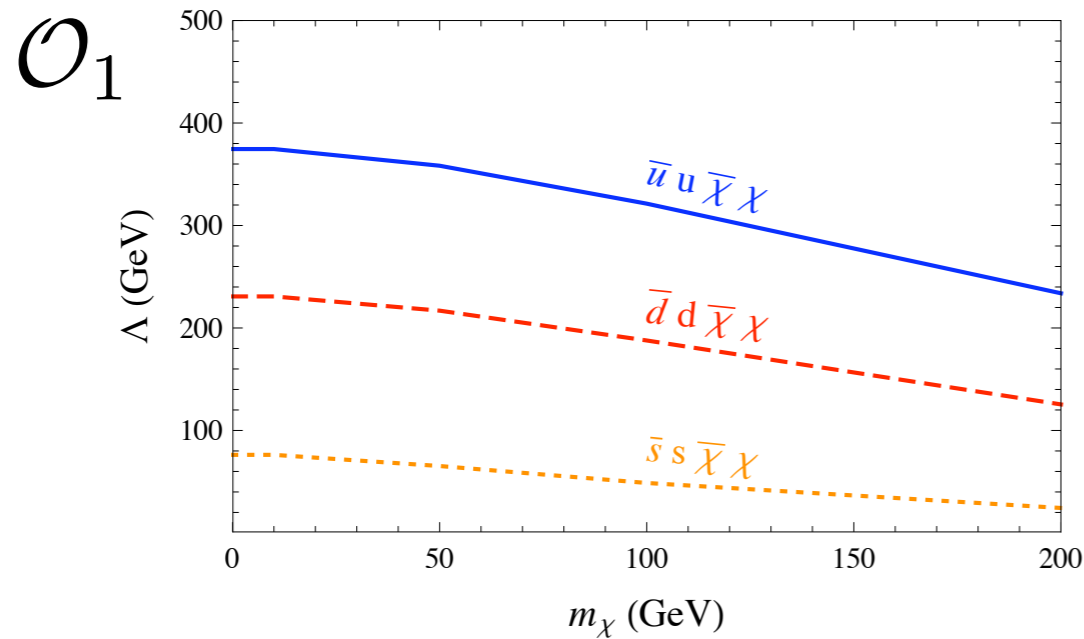


**Observed: 8449 events**

# Bounds on operators

Assume a heavy mediator:  $\Lambda = \frac{M}{\sqrt{g_\chi g_1}}$

Simulate events in calcHEP, one operator at a time



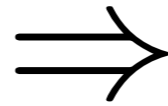
# Collider bounds on direct detection

- Up quark bounds typically strongest
- Collider bounds relatively strongest when DD suppressed e.g. SD, MDDM, light, ....
- iDM splitting not important at colliders
- Tevatron not constrained by velocity distribution - low mass DM
- DM with vector couplings to 2 or 3 gen. quarks
- .....

# Spin independent

$$\mathcal{O}_1 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}q) (\bar{\chi}\chi) ,$$

$$\mathcal{O}_2 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}\gamma_\mu q) (\bar{\chi}\gamma^\mu \chi)$$



$$\sigma_1^{Nq} = \frac{\mu^2}{\pi \Lambda^4} B_{Nq}^2 ,$$

$$\sigma_2^{Nq} = \frac{\mu^2}{\pi \Lambda^4} f_{Nq}^2 ,$$

$$B_u^p = B_d^n = 8.22 \pm 2.26 ,$$

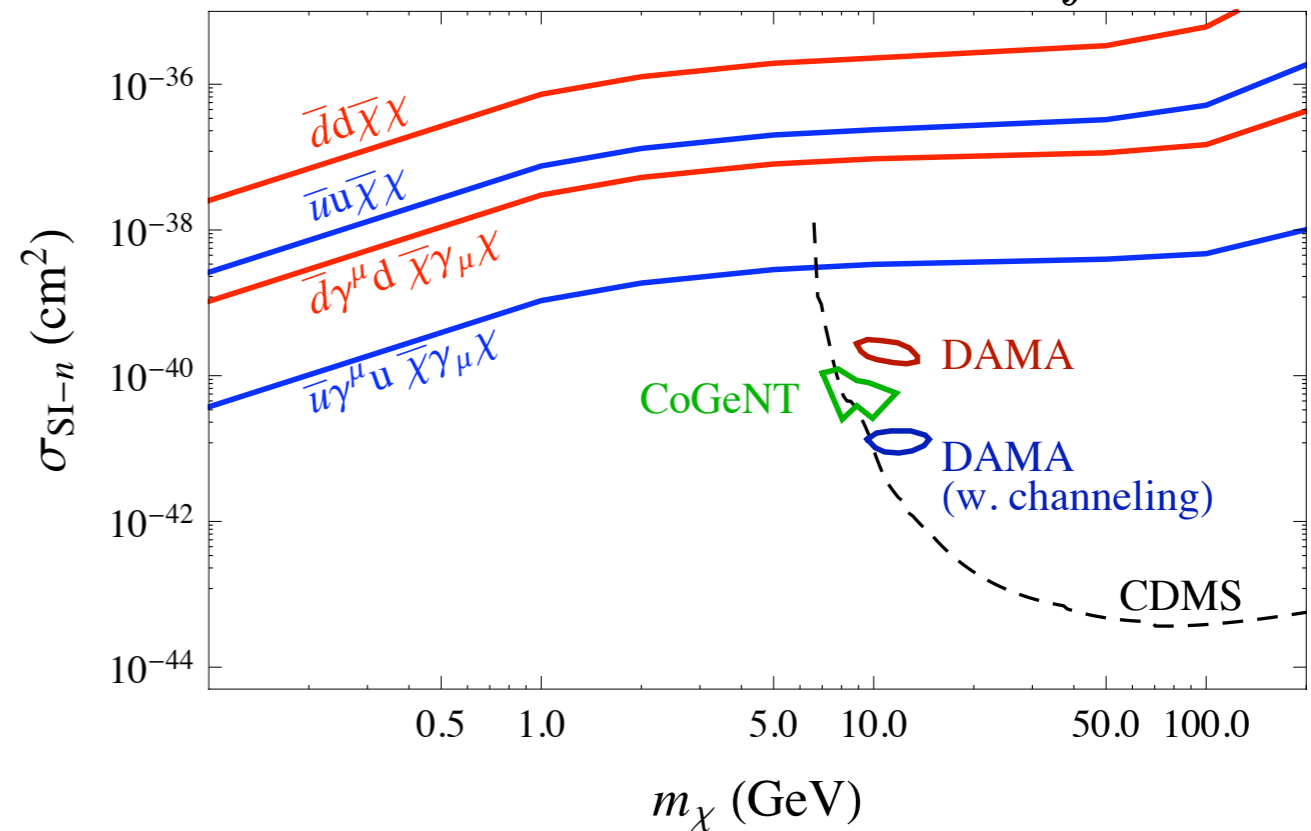
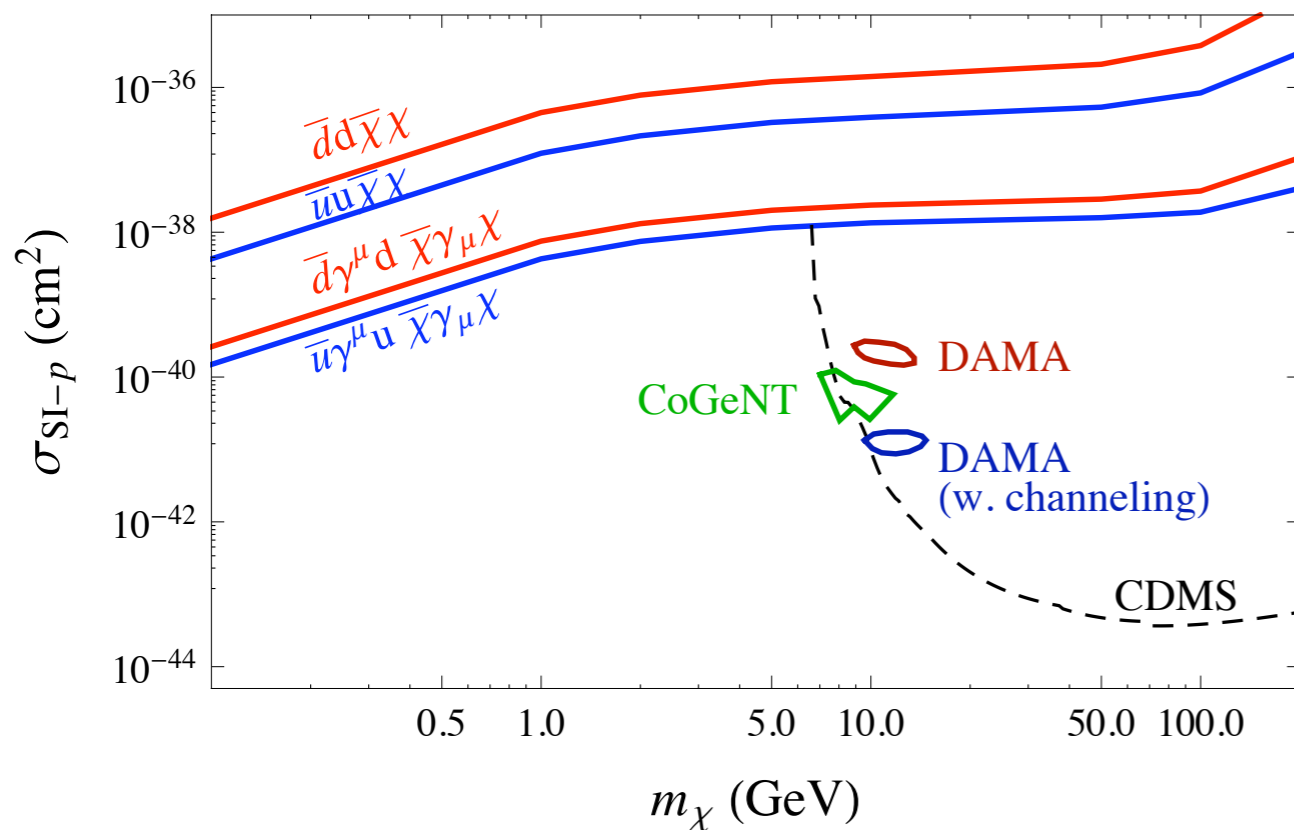
$$f_u^p = f_d^n = 2$$

$$B_d^p = B_u^n = 6.62 \pm 1.92 ,$$

$$f_d^p = f_u^n = 1$$

$$B_s^p = B_s^n = 3.36 \pm 1.45$$

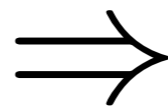
otherwise  $f = 0$



# Spin independent

$$\mathcal{O}_1 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}q) (\bar{\chi}\chi) ,$$

$$\mathcal{O}_2 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q}\gamma_\mu q) (\bar{\chi}\gamma^\mu \chi)$$



$$\sigma_1^{Nq} = \frac{\mu^2}{\pi \Lambda^4} B_{Nq}^2 ,$$

$$\sigma_2^{Nq} = \frac{\mu^2}{\pi \Lambda^4} f_{Nq}^2 ,$$

$$B_u^p = B_d^n = 8.22 \pm 2.26 ,$$

$$f_u^p = f_d^n = 2$$

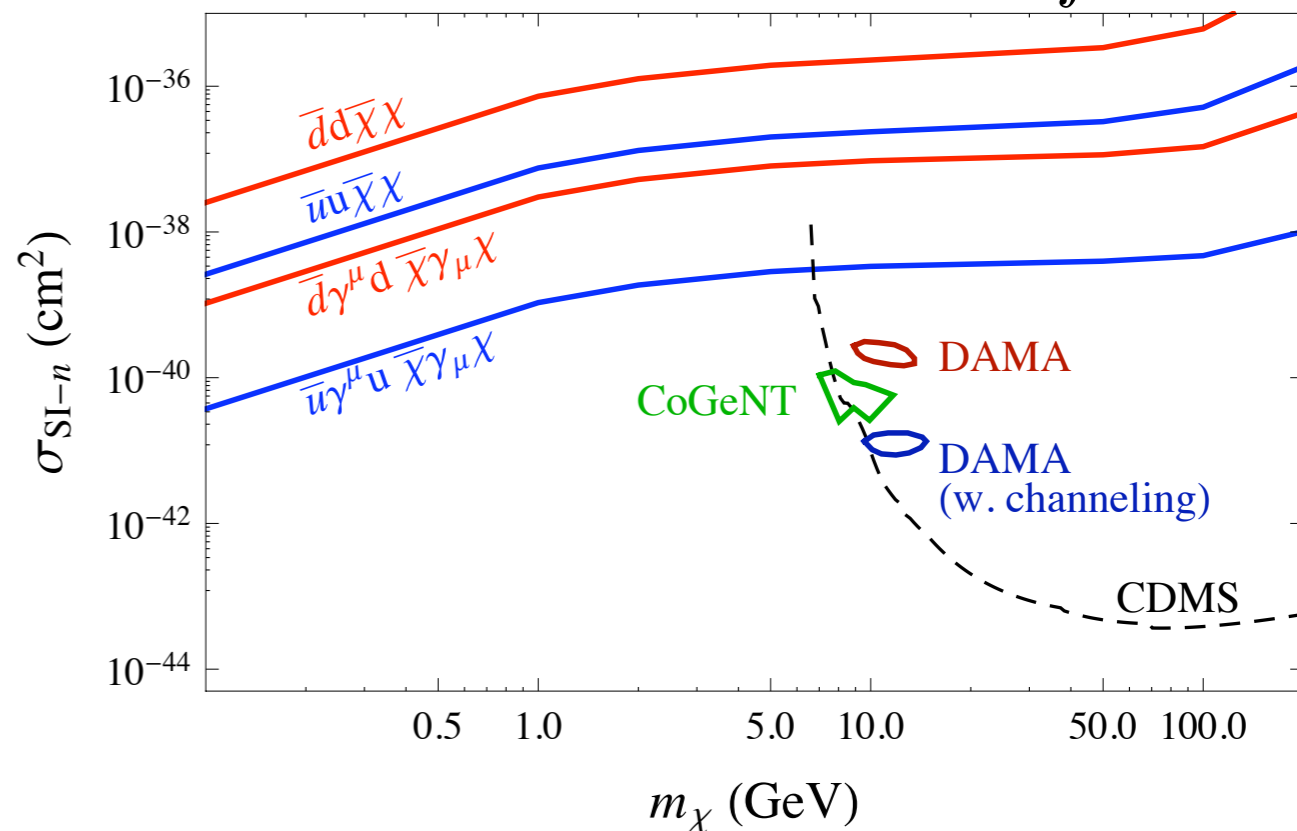
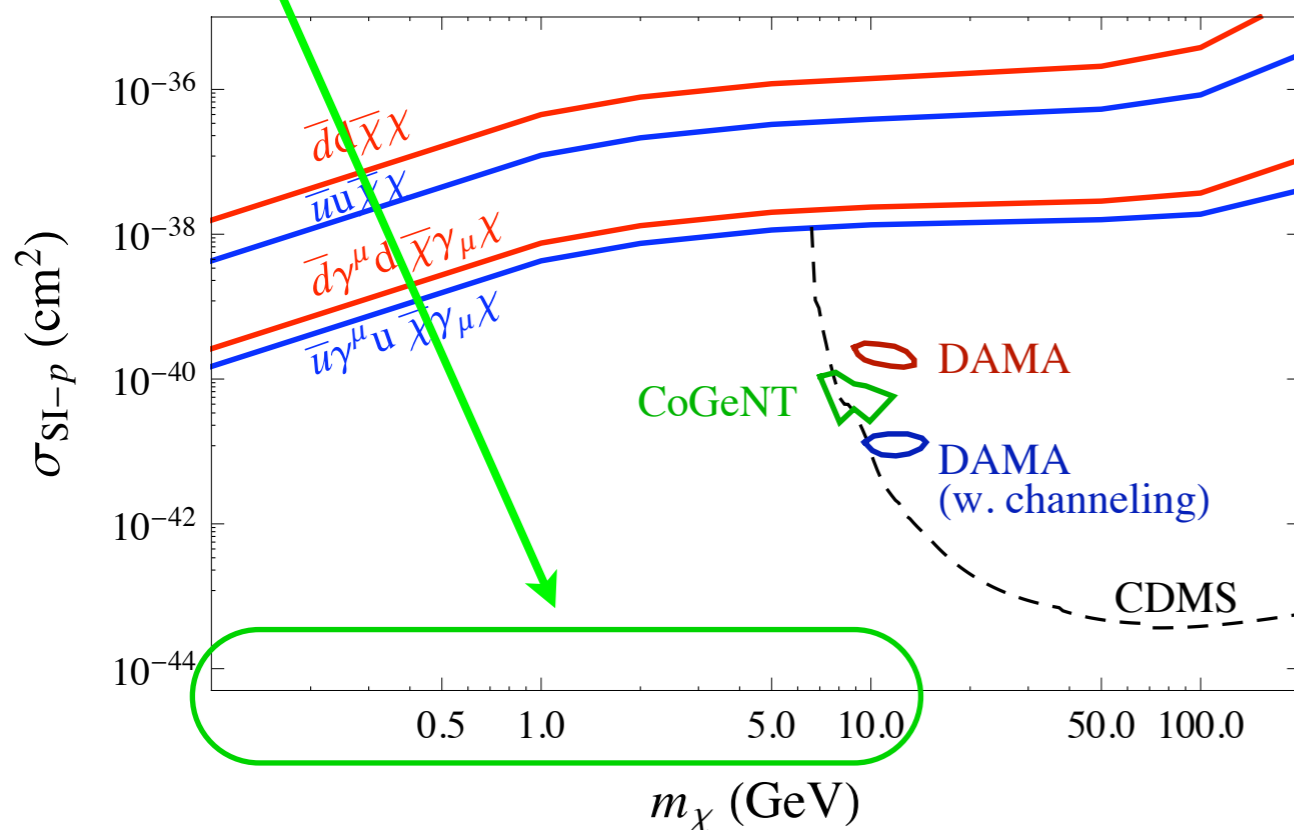
$$B_d^p = B_u^n = 6.62 \pm 1.92 ,$$

$$f_d^p = f_u^n = 1$$

$$B_s^p = B_s^n = 3.36 \pm 1.45$$

otherwise  $f = 0$

World's best limits  
at low mass



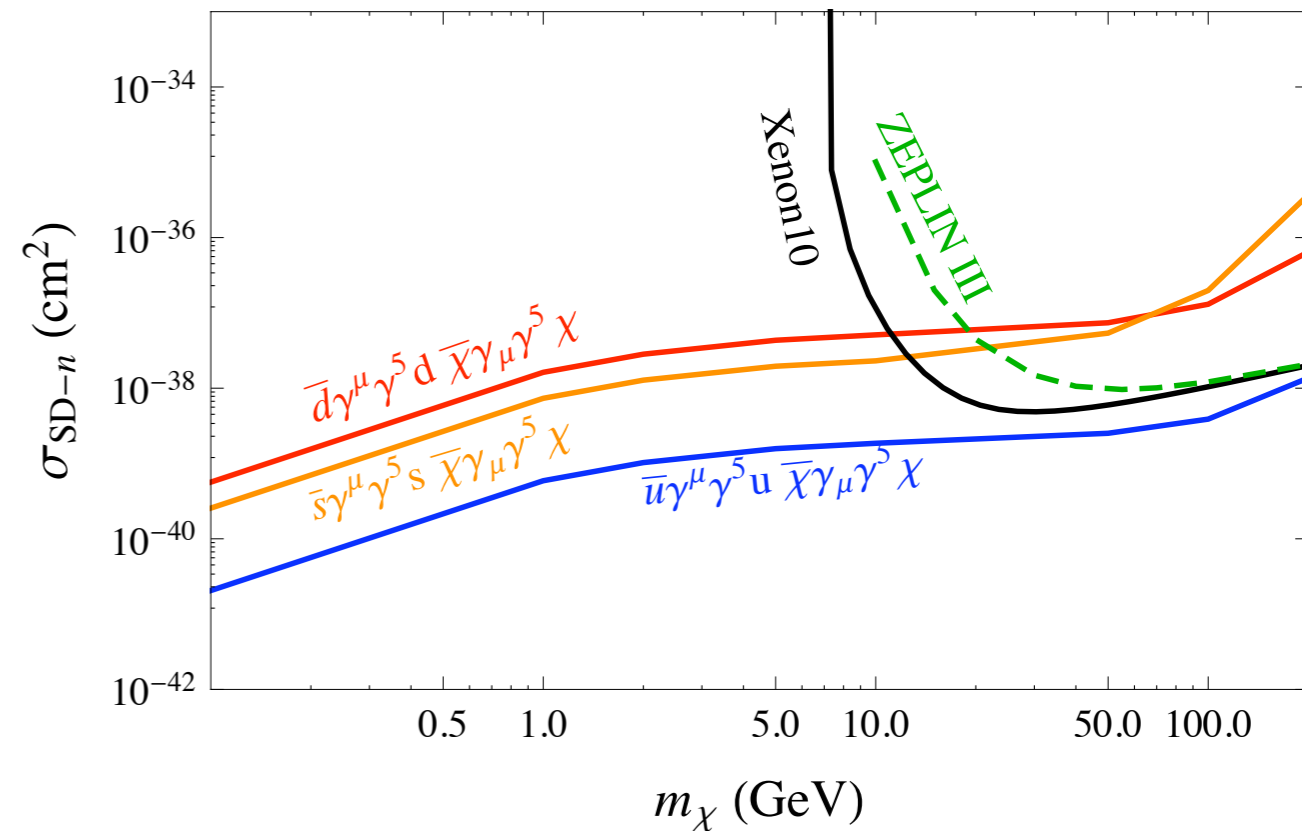
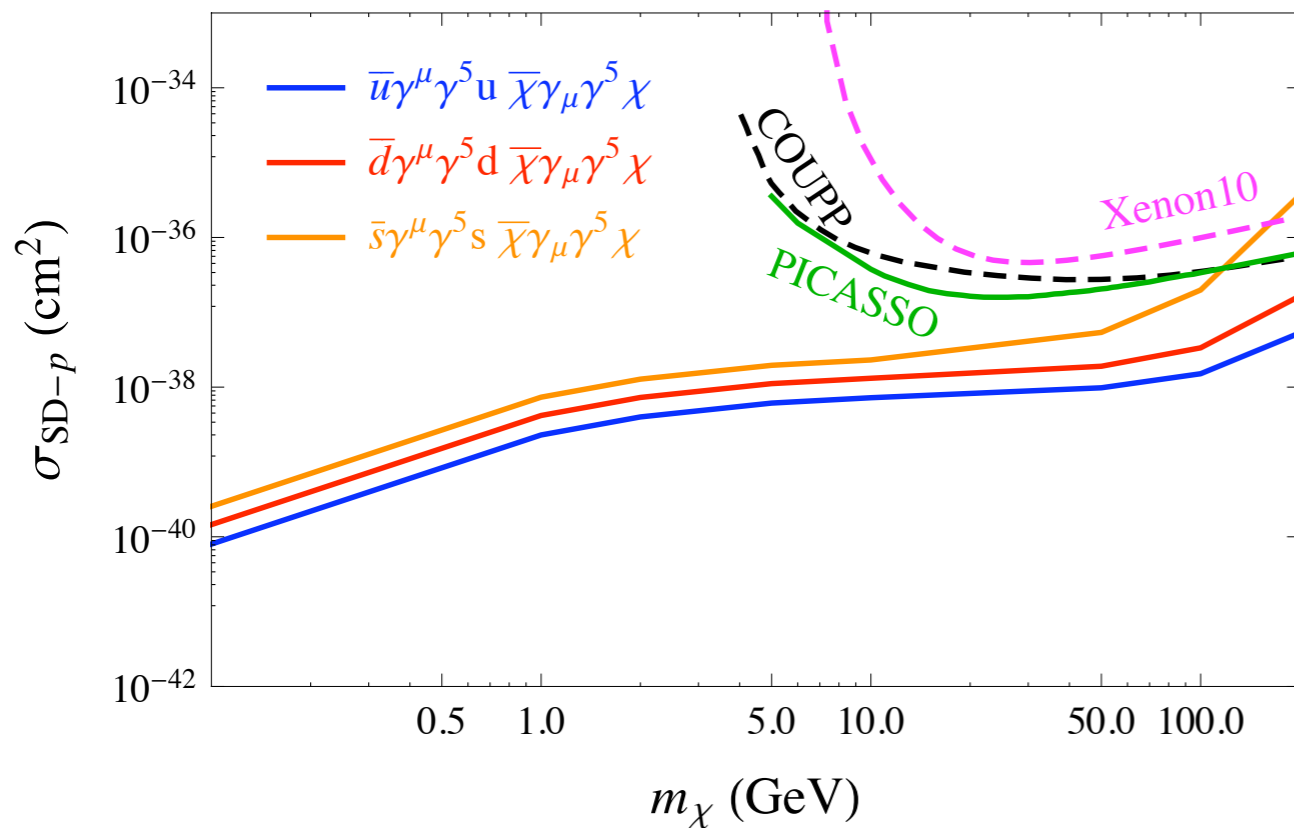
# Spin dependent

$$\mathcal{O}_3 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q} \gamma_\mu \gamma_5 q) (\bar{\chi} \gamma^\mu \gamma_5 \chi)$$

$$\mathcal{O}_3^{Nq} = \Delta_q^N \frac{(\bar{N} \gamma^\mu \gamma_5 N) (\bar{\chi} \gamma_\mu \gamma_5 \chi)}{\Lambda^2}$$

$$\sigma_3^{Nq} = \frac{3 \mu^2}{\pi \Lambda^4} (\Delta_q^N)^2$$

$$\begin{aligned} \Delta_u^p &= \Delta_d^n = 0.842 \pm 0.012, \\ \Delta_d^p &= \Delta_u^n = -0.427 \pm 0.013, \\ \Delta_s^p &= \Delta_s^n = -0.085 \pm 0.018. \end{aligned}$$





# Spin dependent

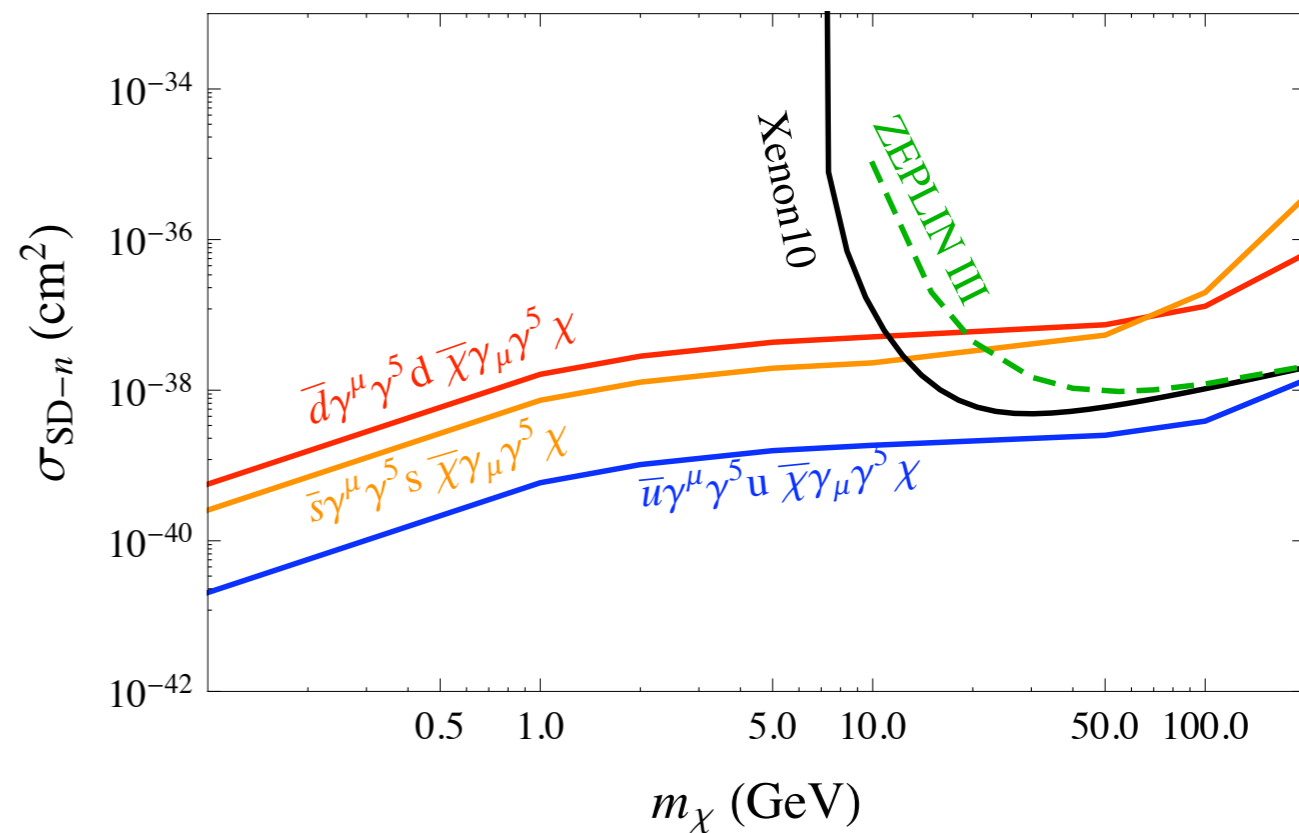
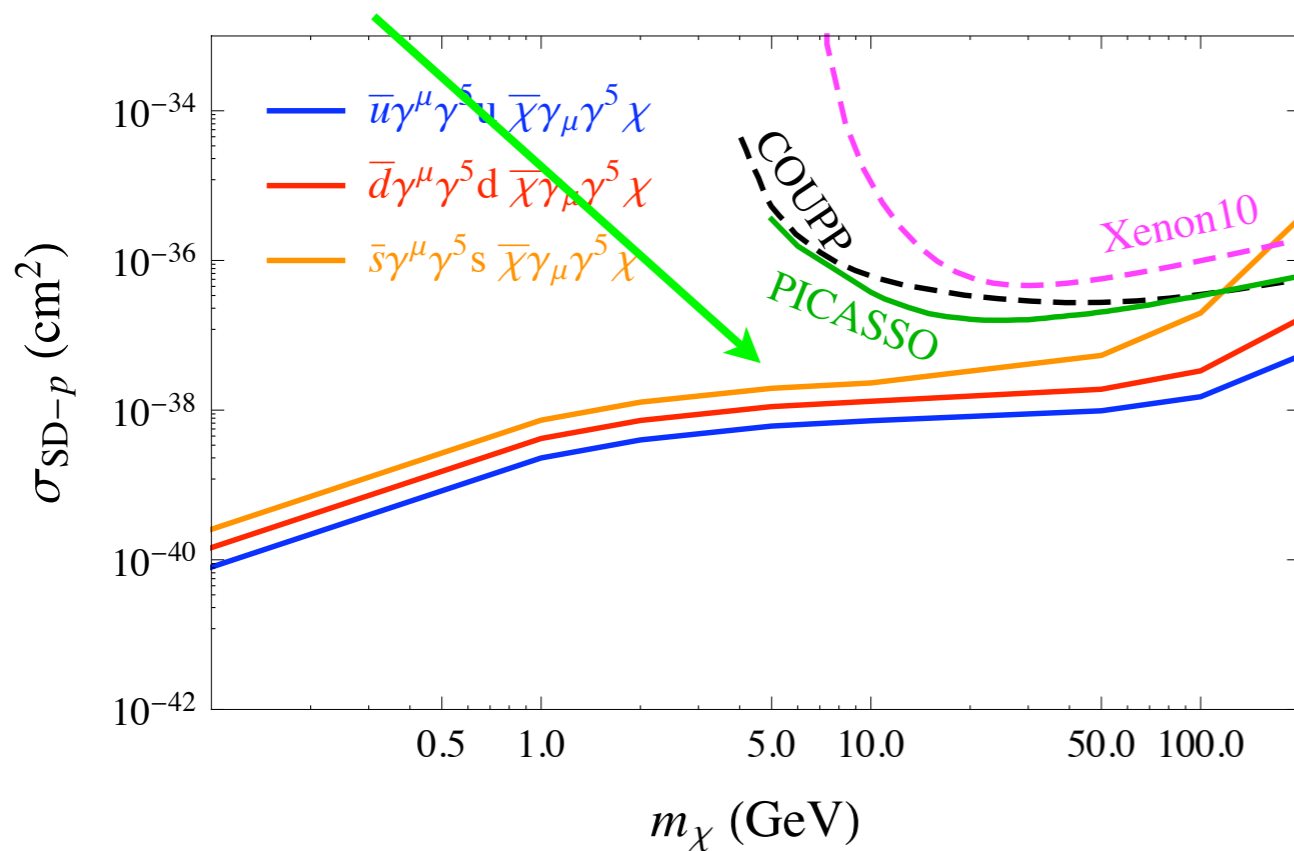
$$\mathcal{O}_3 = \frac{i g_\chi g_q}{q^2 - M^2} (\bar{q} \gamma_\mu \gamma_5 q) (\bar{\chi} \gamma^\mu \gamma_5 \chi)$$

$$\mathcal{O}_3^{Nq} = \Delta_q^N \frac{(\bar{N} \gamma^\mu \gamma_5 N) (\bar{\chi} \gamma_\mu \gamma_5 \chi)}{\Lambda^2}$$

$$\sigma_3^{Nq} = \frac{3 \mu^2}{\pi \Lambda^4} (\Delta_q^N)^2$$

$$\begin{aligned} \Delta_u^p &= \Delta_d^n = 0.842 \pm 0.012, \\ \Delta_d^p &= \Delta_u^n = -0.427 \pm 0.013, \\ \Delta_s^p &= \Delta_s^n = -0.085 \pm 0.018. \end{aligned}$$

World's best limits, up to  
~200 GeV

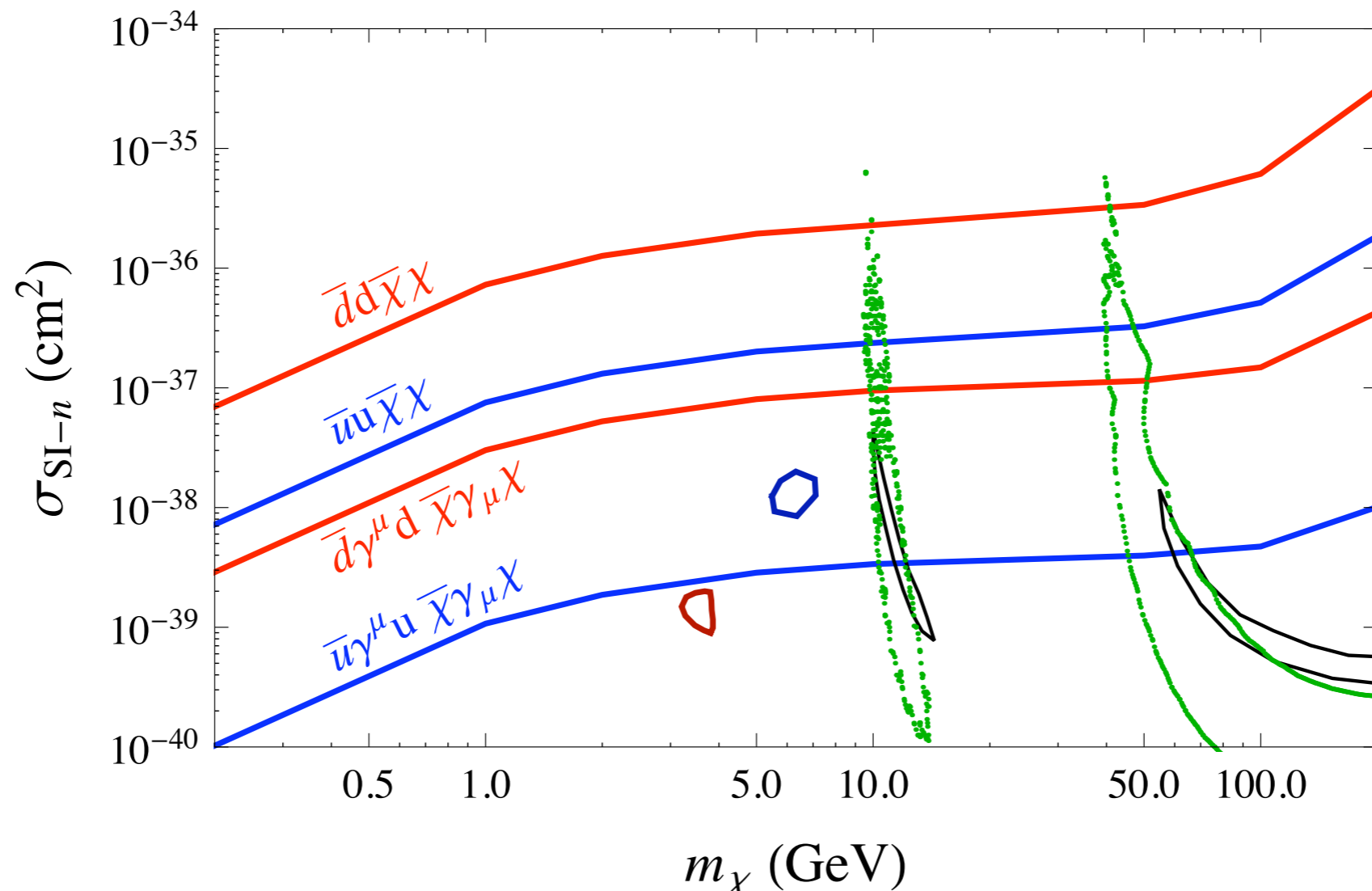


# iDM, exothermic

$$\frac{dR}{dE_R} \propto n_\chi \sigma_N \int_{v_{min}}^{v_{esc}} \frac{f(v)}{v} dv,$$

$\delta$

$$v_{min} = \sqrt{\frac{1}{2m_T E_R} \left( \frac{m_T E_R}{\mu_T} + \delta \right)}$$

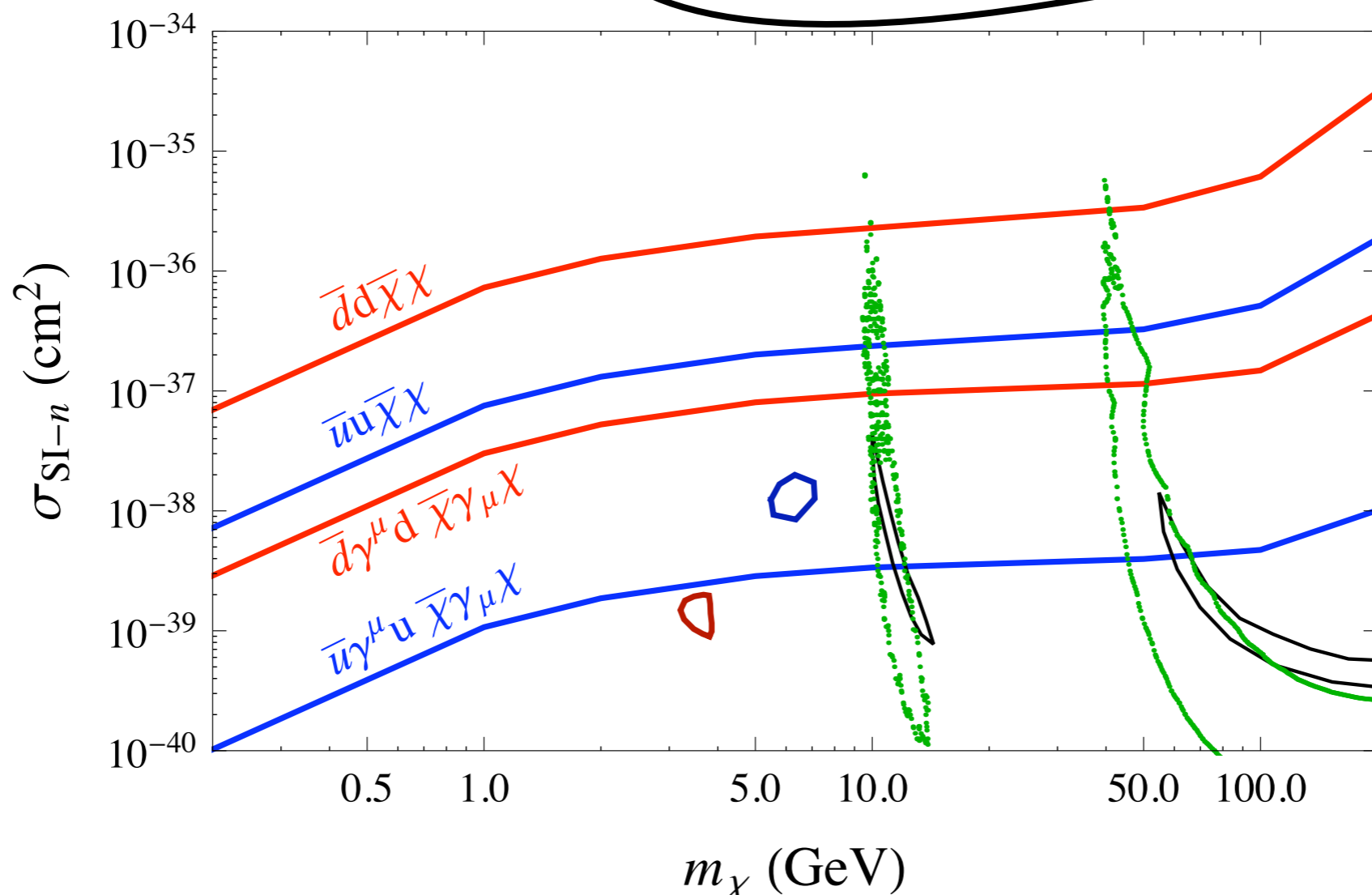


# iDM, exothermic

$$\frac{dR}{dE_R} \propto n_\chi \sigma_N \int_{v_{min}}^{v_{esc}} \frac{f(v)}{v} dv, \quad \delta \begin{array}{c} \xrightarrow{\chi^*} \\ \xleftarrow{\chi} \end{array}$$

$$v_{min} = \sqrt{\frac{1}{2m_T E_R} \left( \frac{m_T E_R}{\mu_T} + \delta \right)}$$

$\delta \sim 100 \text{ keV}, 5 \text{ keV}$

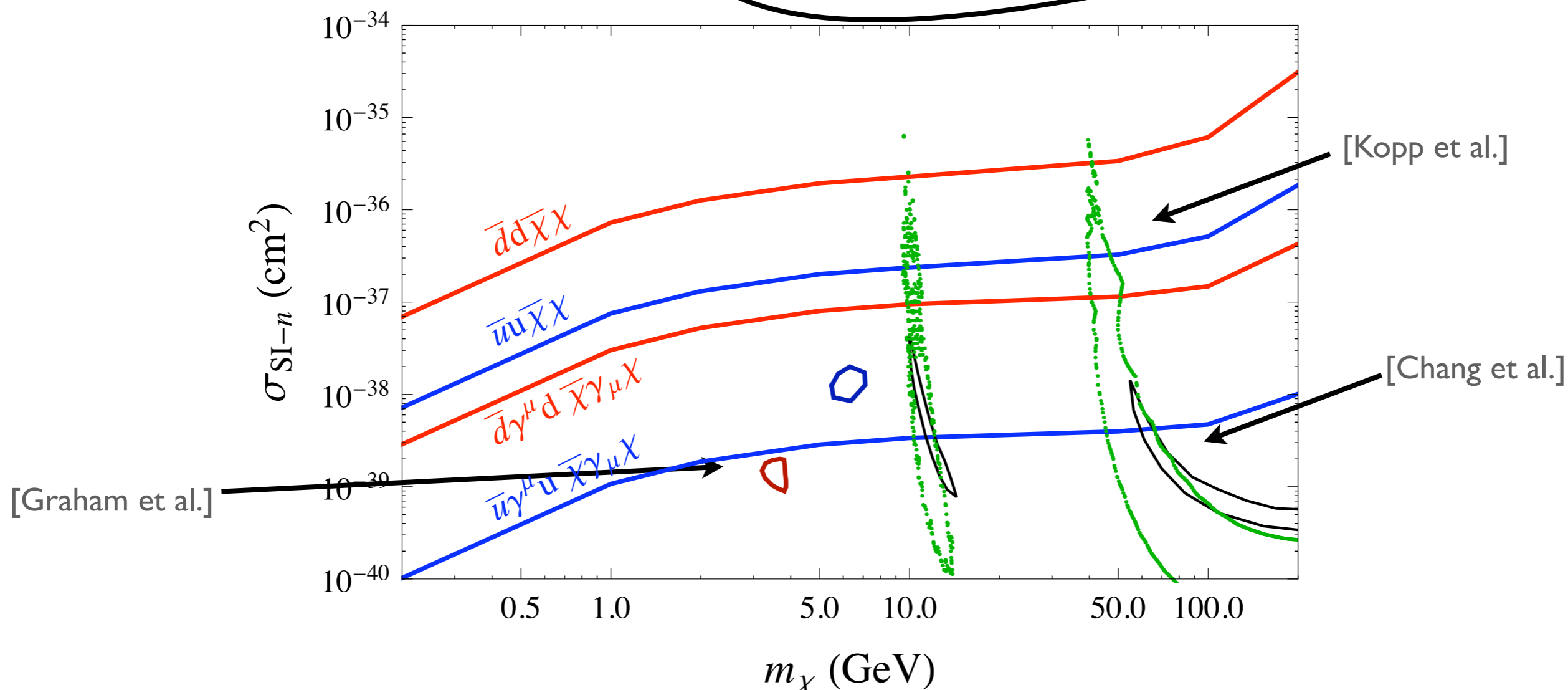


# iDM, exothermic

$$\frac{dR}{dE_R} \propto n_\chi \sigma_N \int_{v_{min}}^{v_{esc}} \frac{f(v)}{v} dv, \quad \delta \begin{array}{c} \xrightarrow{\chi^*} \\ \xleftarrow{\chi} \end{array}$$

$$v_{min} = \sqrt{\frac{1}{2m_T E_R} \left( \frac{m_T E_R}{\mu_T} + \delta \right)}$$

$\delta \sim 100 \text{ keV}, 5 \text{ keV}$



# iDM, exothermic

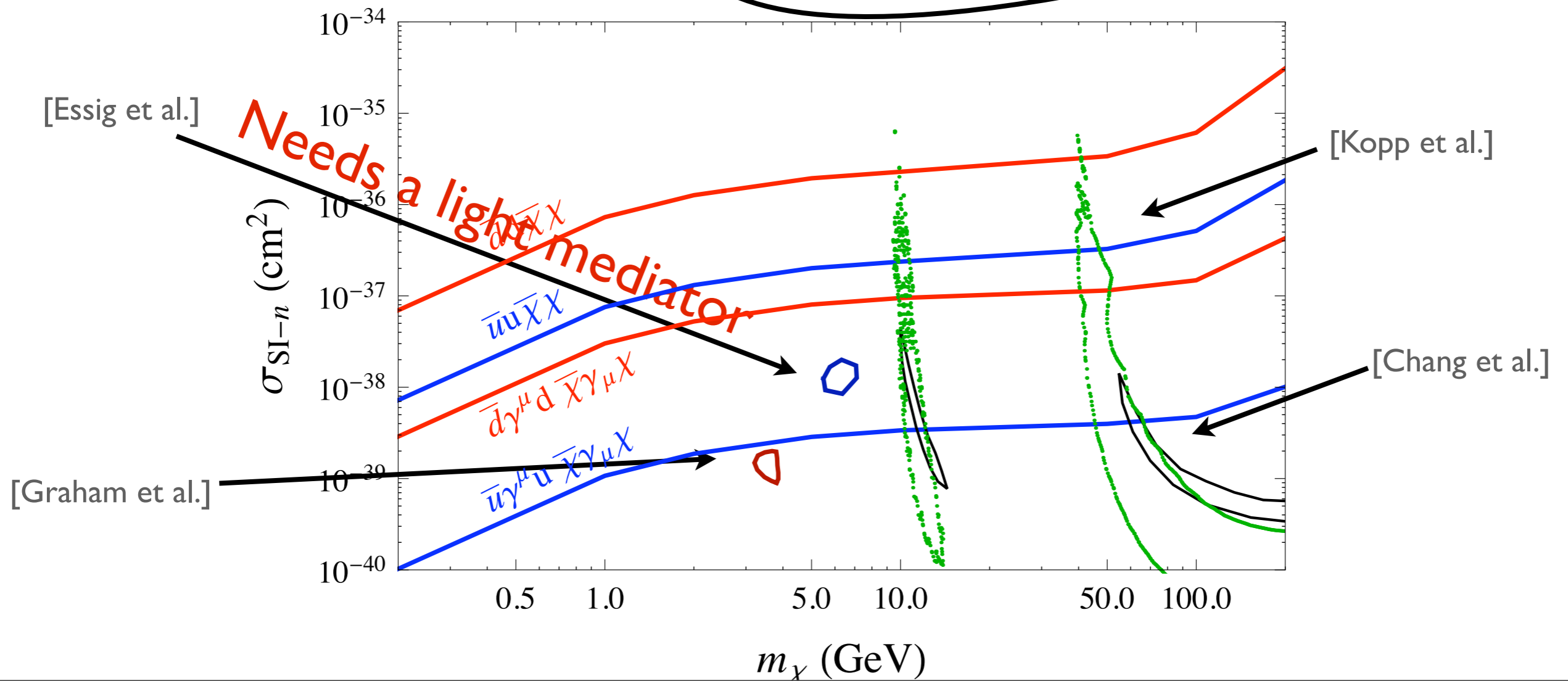
$$\frac{dR}{dE_R} \propto n_\chi \sigma_N \int_{v_{min}}^{v_{esc}} \frac{f(v)}{v} dv,$$

$\delta \begin{matrix} \uparrow \text{iDM} \\ \downarrow \text{exoDM} \end{matrix}$

$\frac{\chi^*}{\chi}$

$$v_{min} = \sqrt{\frac{1}{2m_T E_R} \left( \frac{m_T E_R}{\mu_T} + \delta \right)}$$

$\delta \sim 100 \text{ keV}, 5 \text{ keV}$



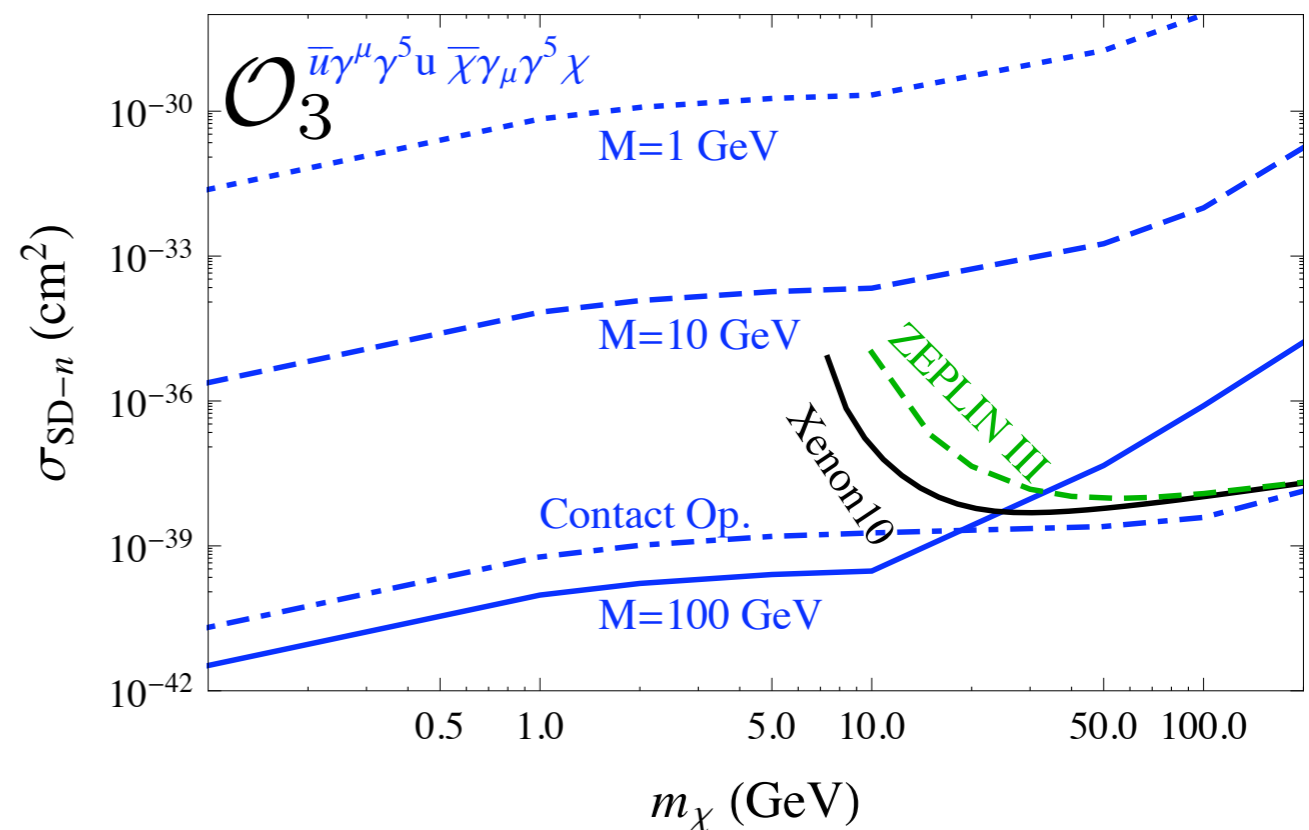
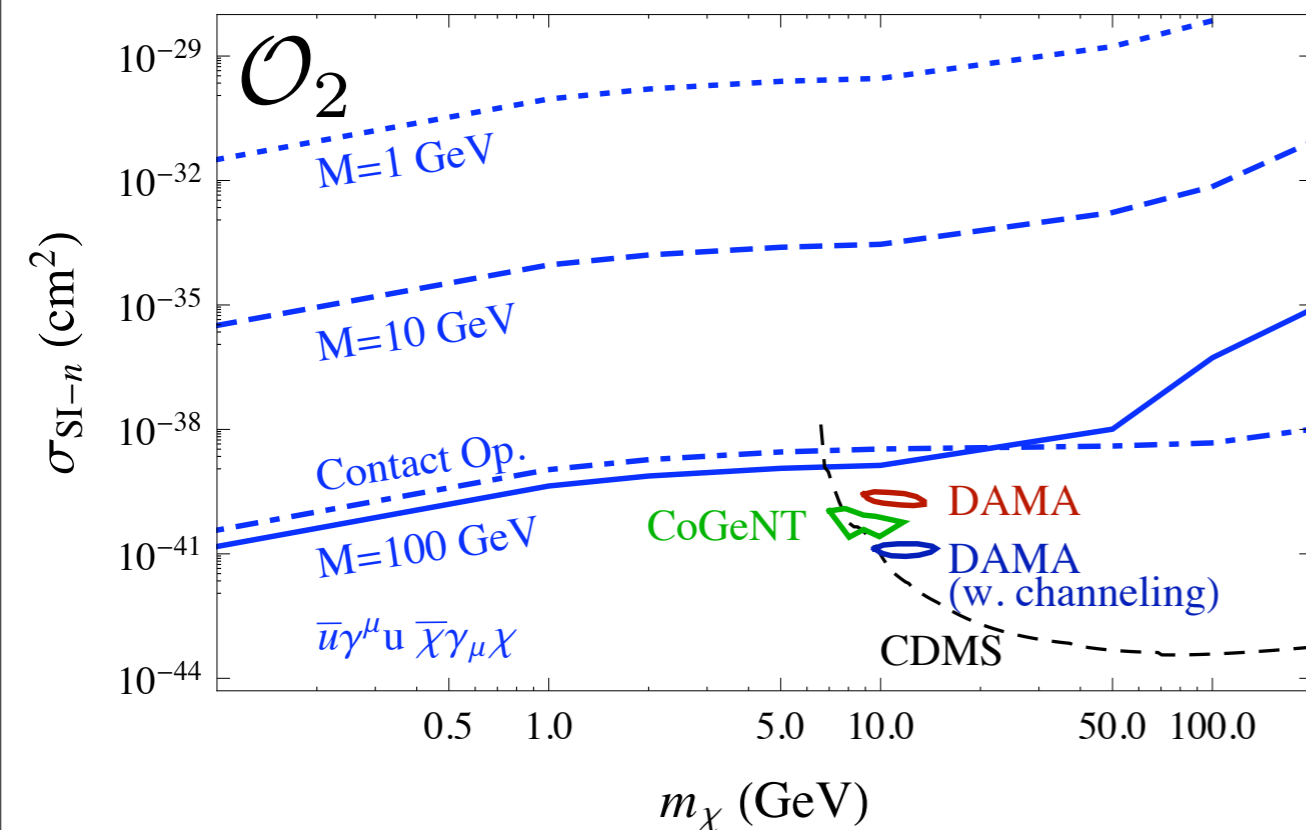
# Light mediators

$$\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\sigma_{1j} \sim \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2}$$

Direct detection wins

Two body vs three body production:  $2 m_{\chi} < M < s^{1/2}$



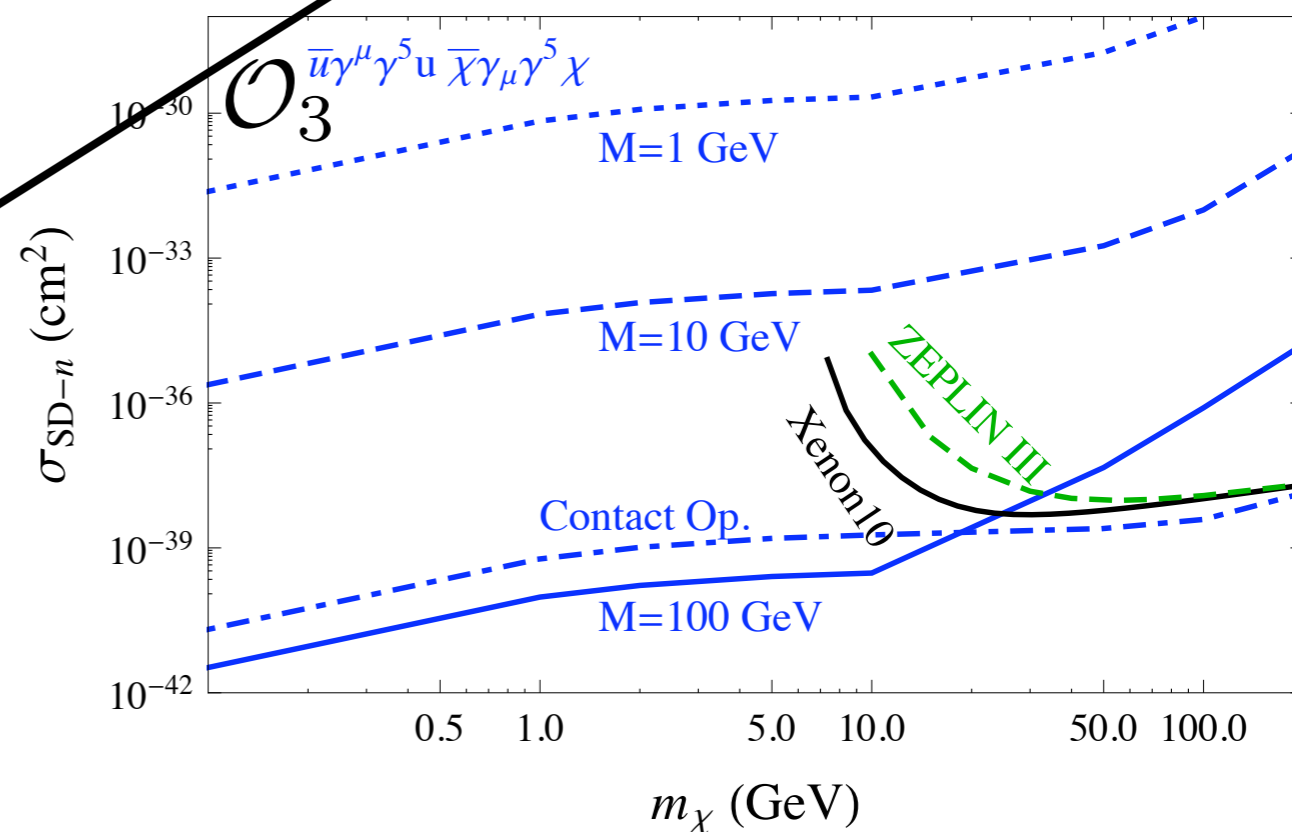
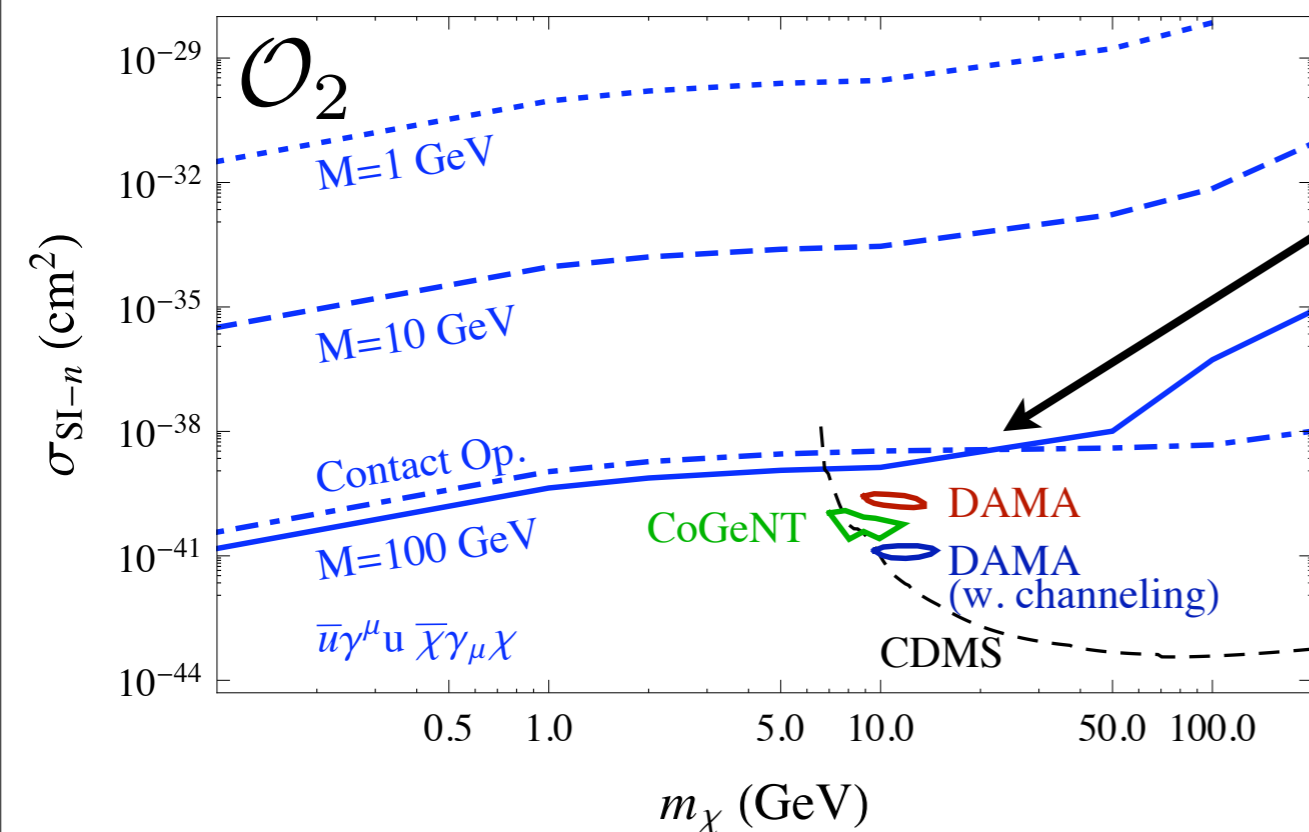
# Light mediators

$$\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\sigma_{1j} \sim \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2}$$

Direct detection wins

Two body vs three body production:  $2 m_{\chi} < M < s^{1/2}$

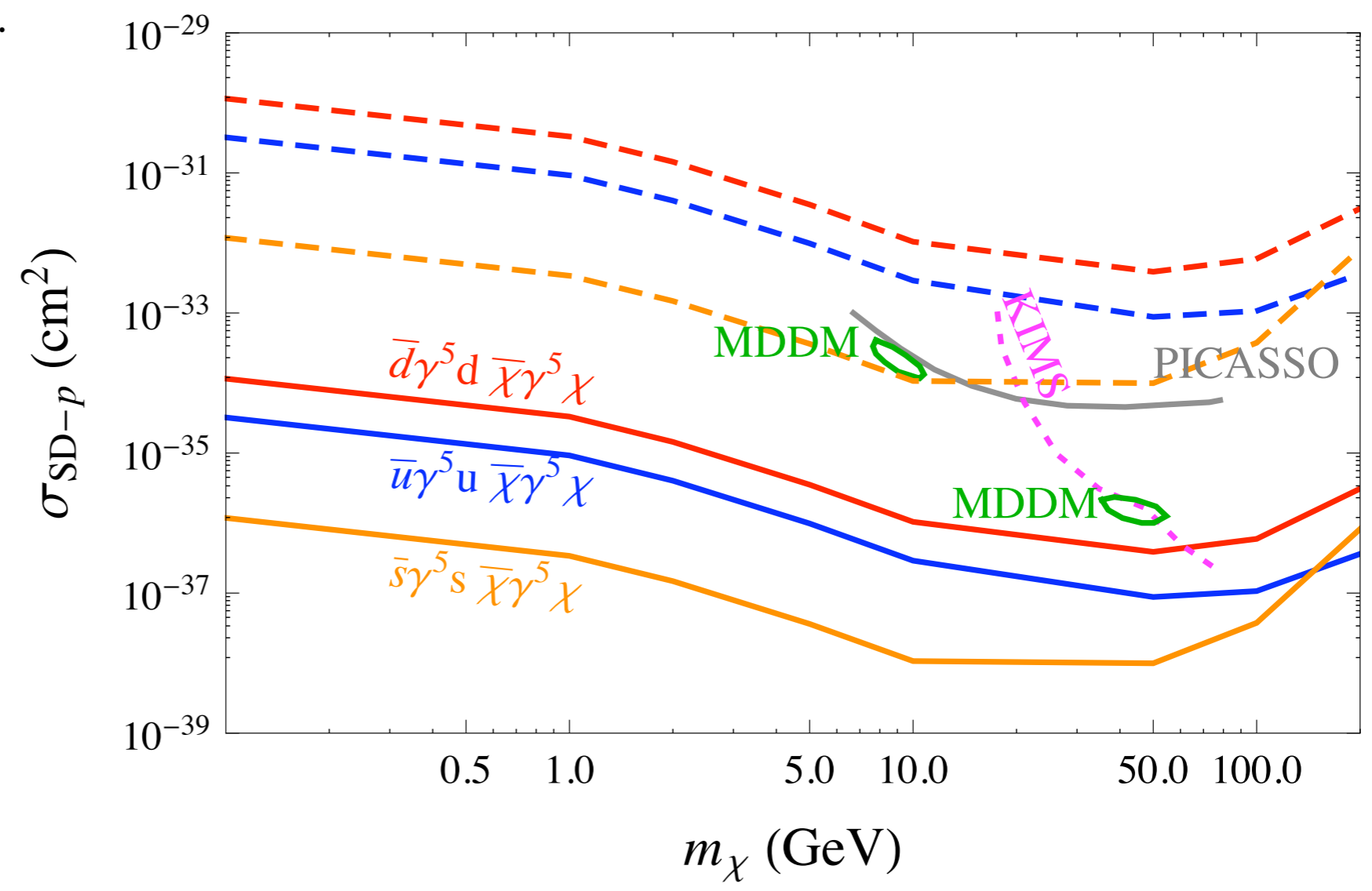


# Momentum dependent

$$\mathcal{O}_4^{Nq} = -i C_q^N \frac{(\bar{N} \gamma_5 N) (\bar{\chi} \gamma_5 \chi)}{\Lambda^2}$$

$$\frac{d\sigma_4^{Nq}}{d\cos\theta} = \frac{1}{32\pi\Lambda^4} \frac{q^4}{(m_\chi + m_N)^2} (C_q^N)^2$$

$$\begin{aligned} C_u^p &= 168.5, & C_u^n &= -165.2, \\ C_d^p &= -164.2, & C_d^n &= 165.8, \\ C_s^p &= -4.3, & C_s^n &= -0.67. \end{aligned}$$



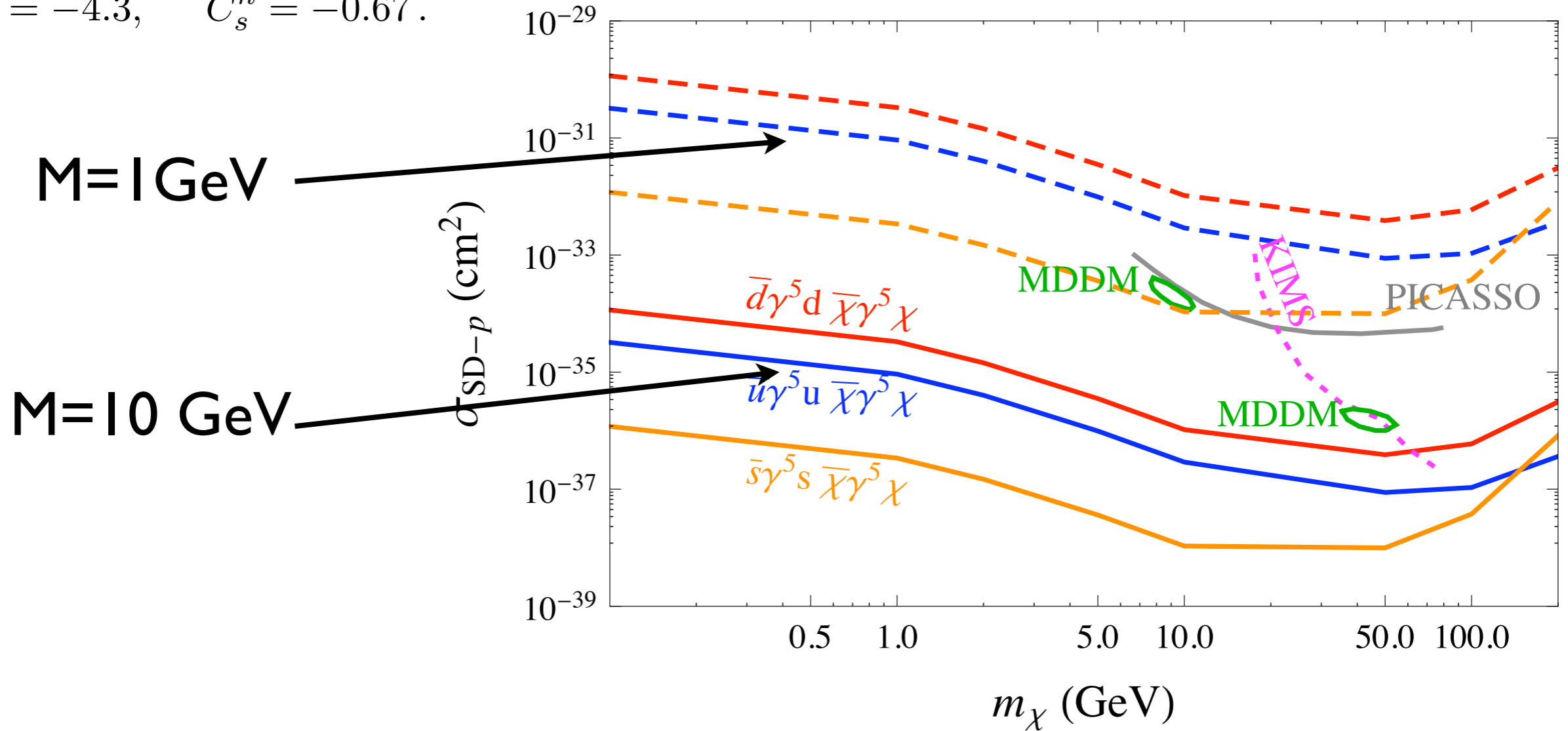


# Momentum dependent

$$O_4^{Nq} = -i C_q^N \frac{(\bar{N} \gamma_5 N) (\bar{\chi} \gamma_5 \chi)}{\Lambda^2}$$

$$\frac{d\sigma_4^{Nq}}{d \cos \theta} = \frac{1}{32\pi \Lambda^4} \frac{q^4}{(m_\chi + m_N)^2} (C_q^N)^2$$

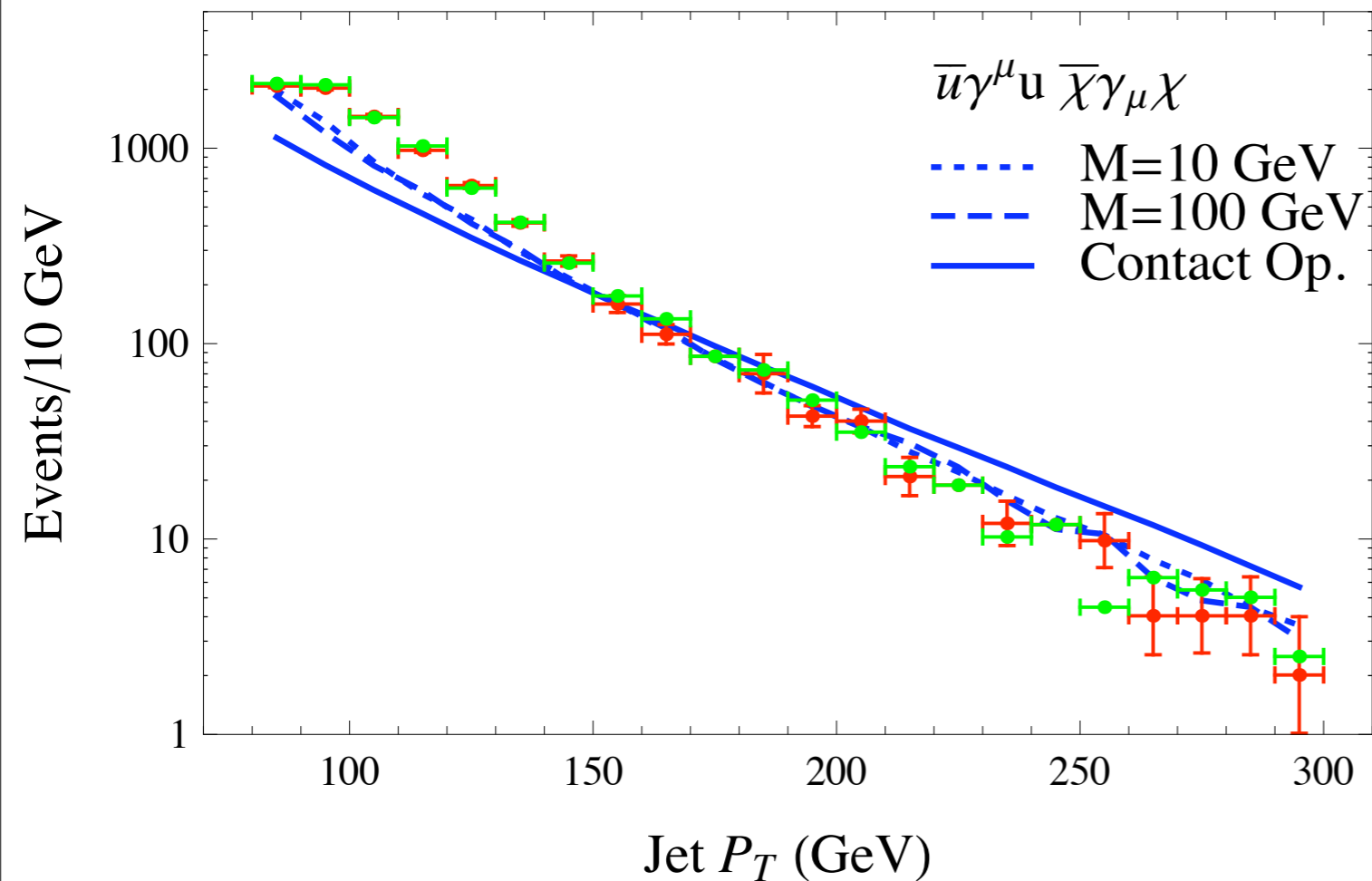
- $C_u^p = 168.5, \quad C_u^n = -165.2,$
- $C_d^p = -164.2, \quad C_d^n = 165.8,$
- $C_s^p = -4.3, \quad C_s^n = -0.67.$



# Improvements?

So far only CDF analysis on 1/fb  
Mono-photon could also be done

$$m_\chi = 10 \text{ GeV}$$



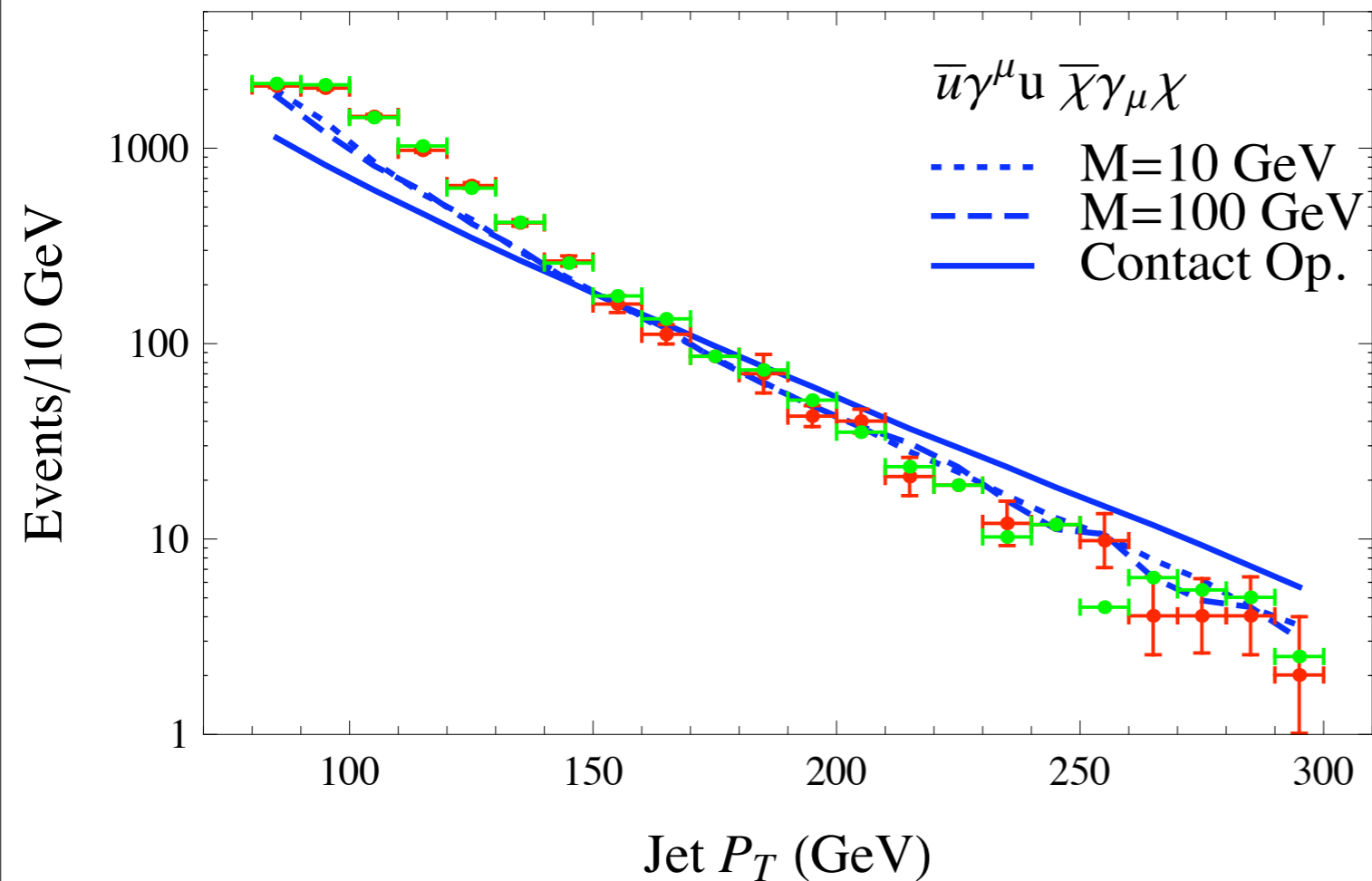
Use shape information

Tevatron reach limited to  $\sim 300$  GeV

# Improvements?

So far only CDF analysis on 1/fb  
Mono-photon could also be done

$$m_\chi = 10 \text{ GeV}$$



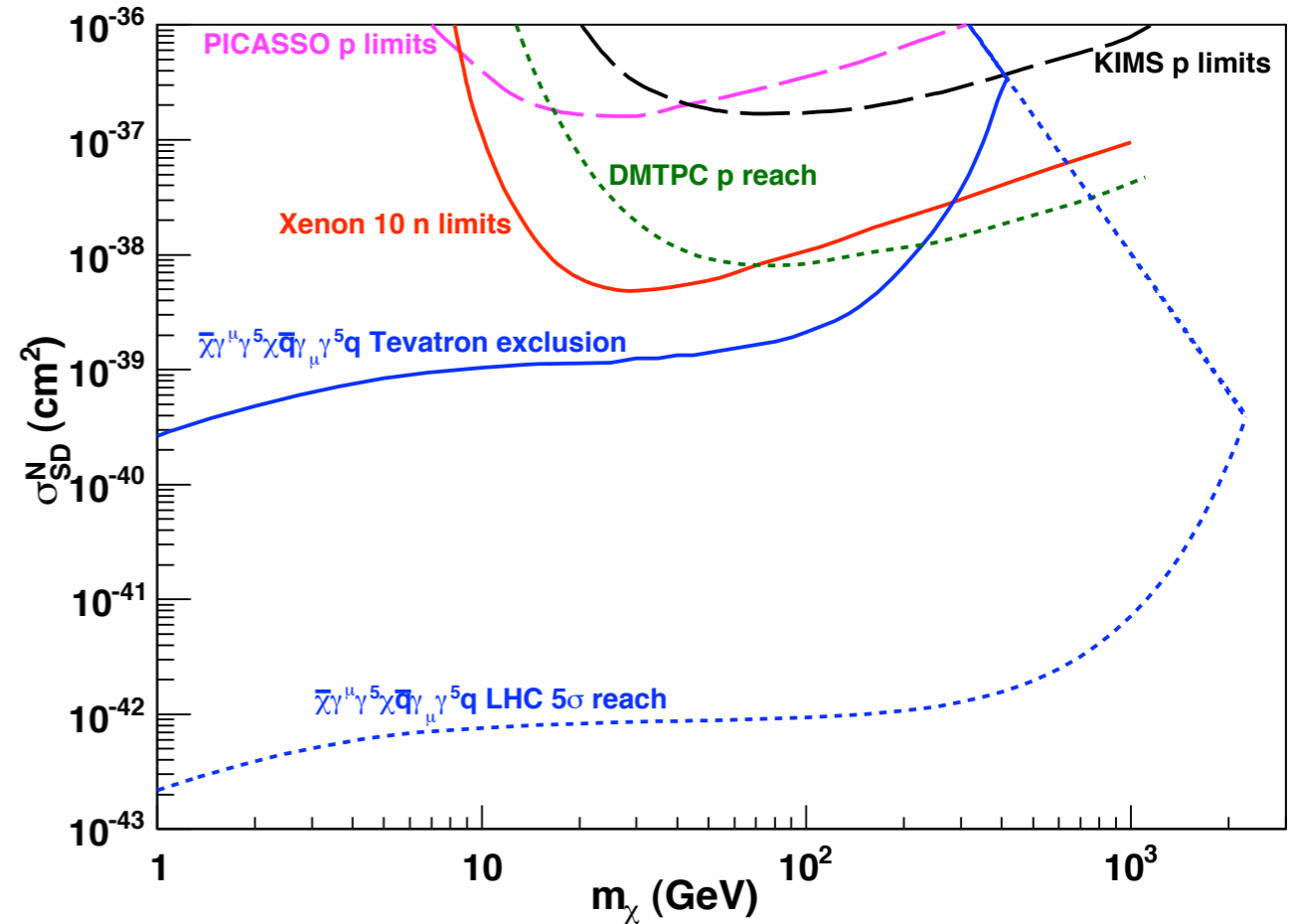
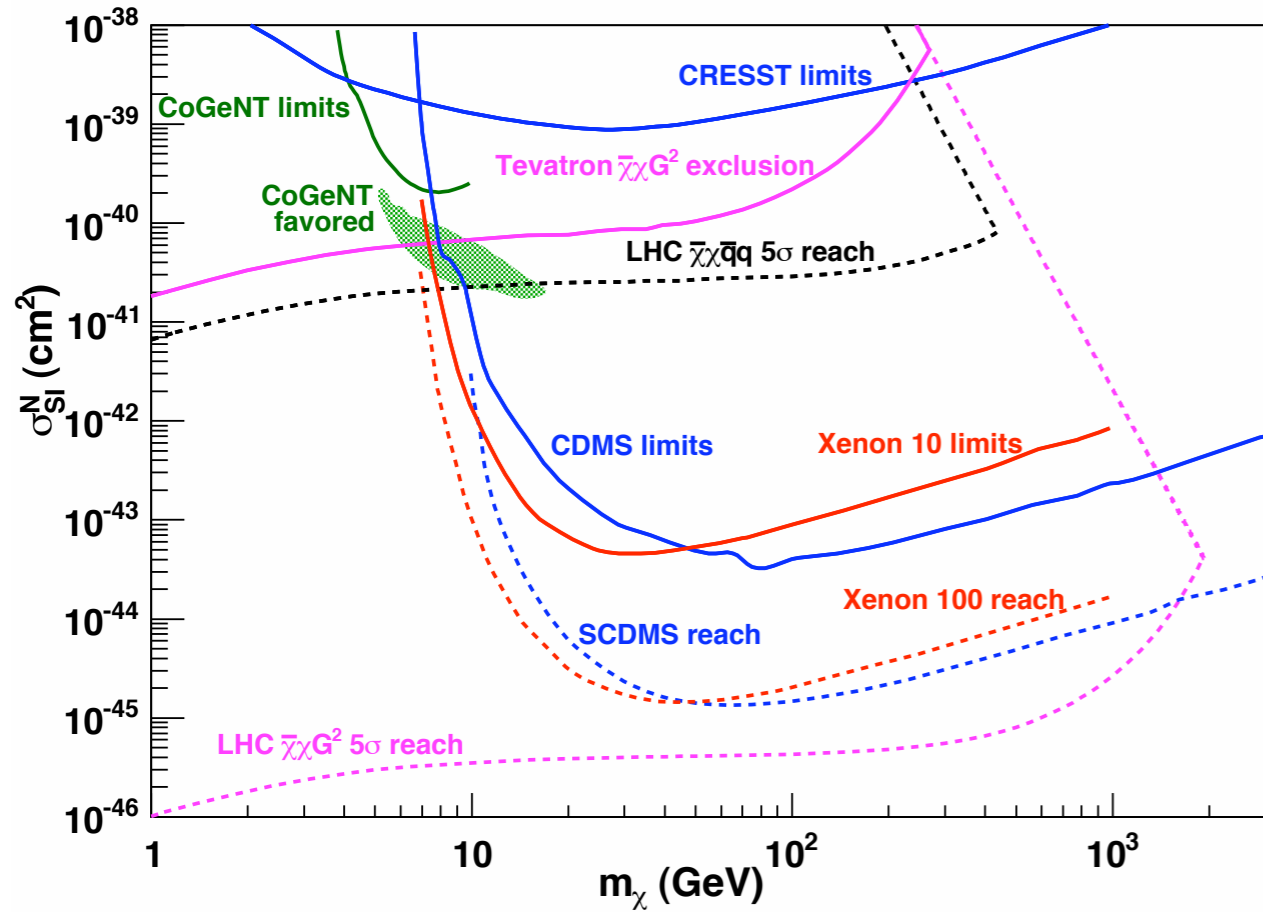
Use shape information

Recently CDF + Bai,  
Harnik, PJF have  
started a “real”  
analysis on full data  
set!

Tevatron reach limited to  $\sim 300$  GeV

# Improvements

[Goodman, Ibe, Rajaraman, Shepherd, Tait, Yu]



## LHC

$$\sqrt{s} = 14 \text{ TeV}$$

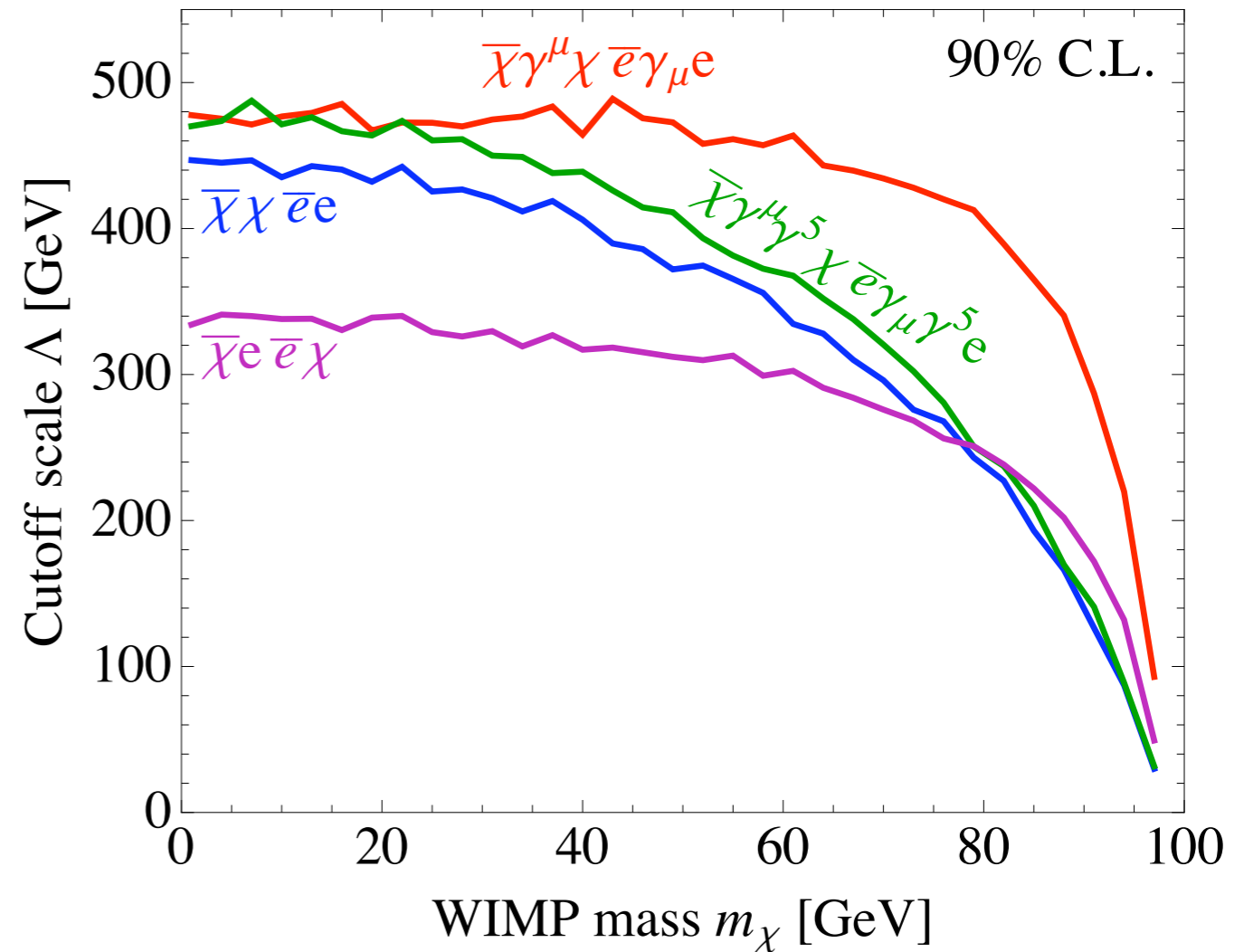
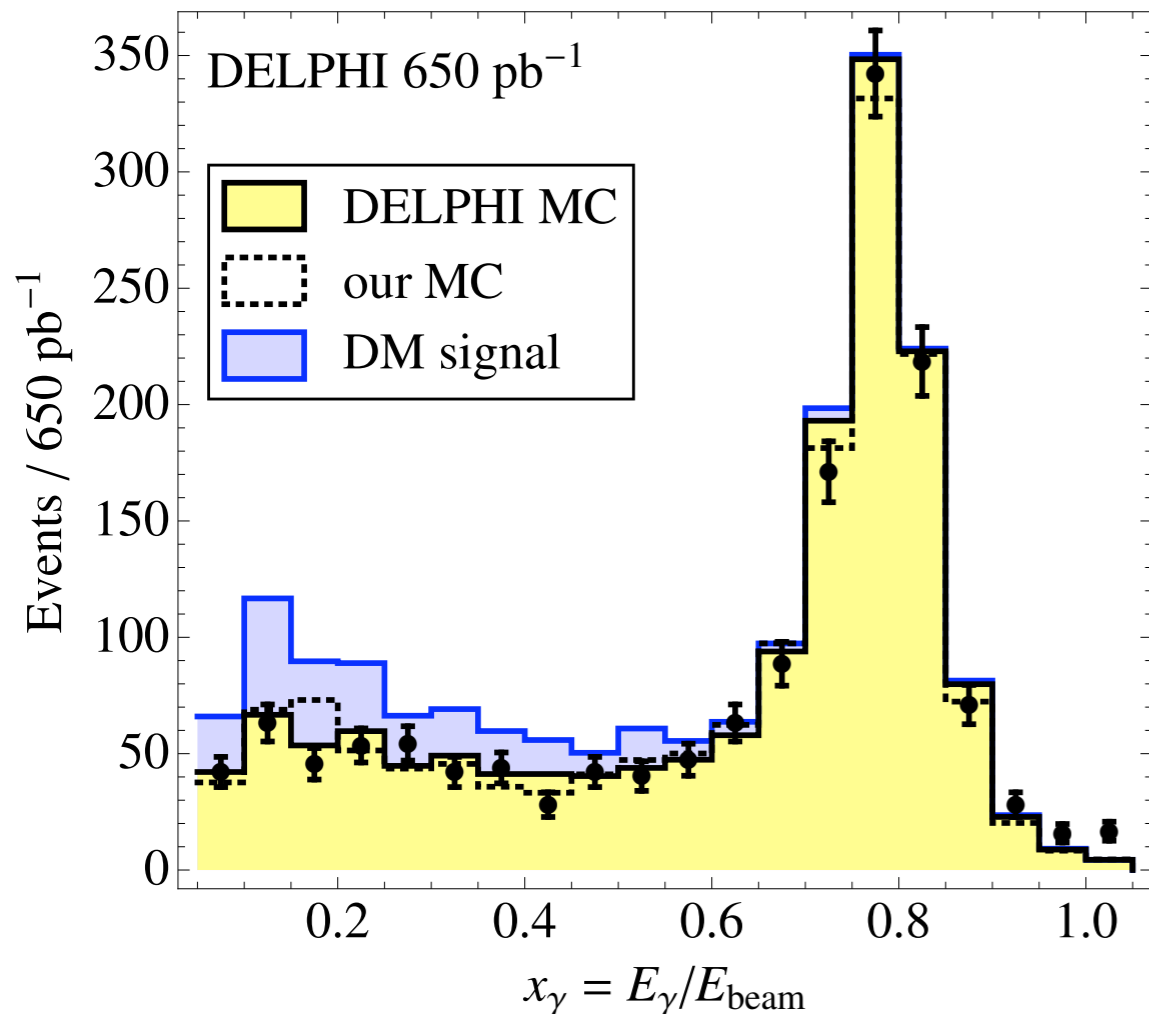
$$\mathcal{L} = 100 \text{ fb}^{-1}$$

$$\cancel{E}_T > 500 \text{ GeV}$$

No longer monojet search  
BSM backgrounds?

# Leptophilic dark matter

## Mono-photons at LEP



More model dependent than Tevatron constraints

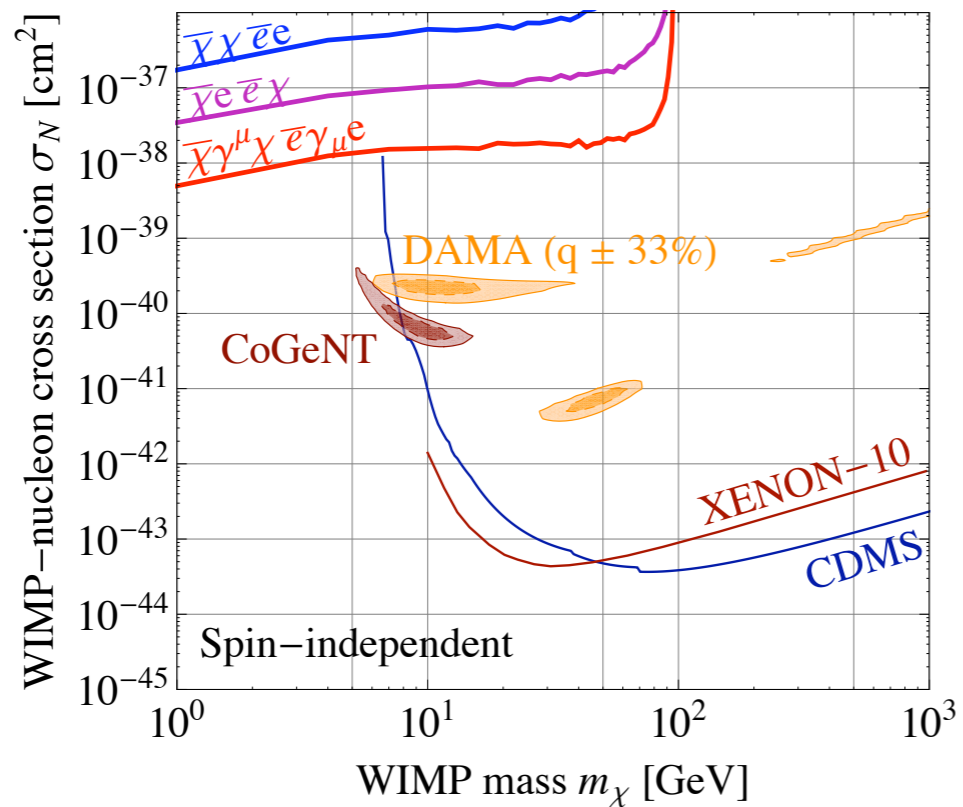
Consider two “extreme” hypotheses:

DM has equal coupling to all SM fermions

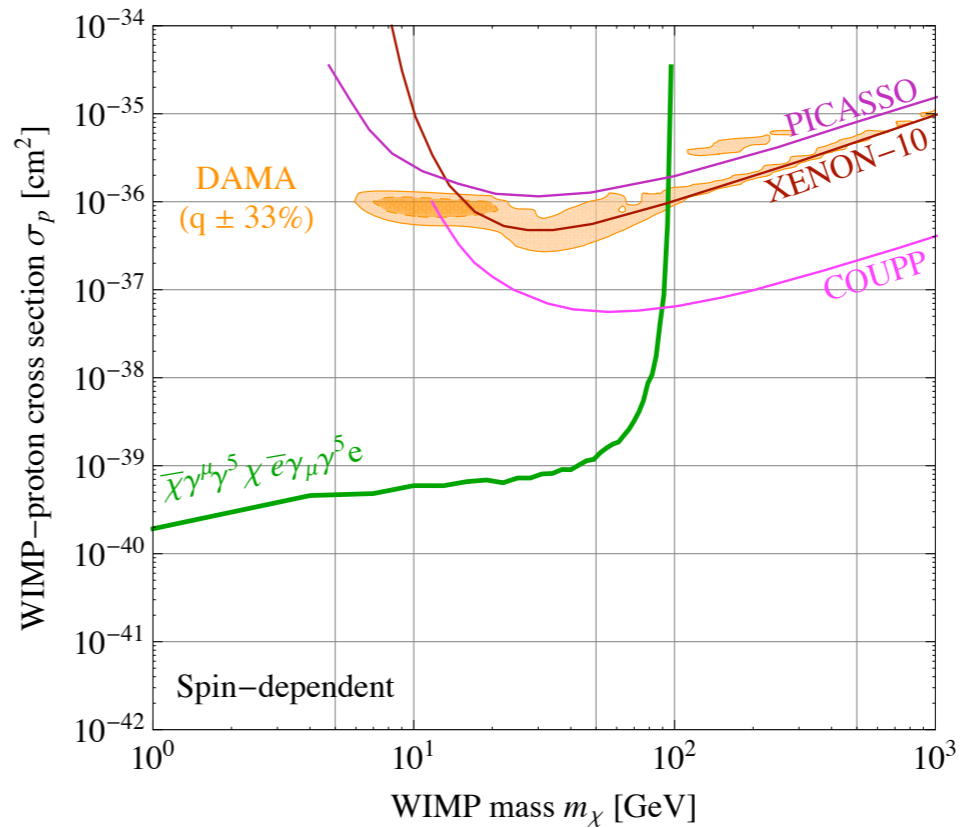
DM has equal coupling to all leptons

# Leptophilic dark matter

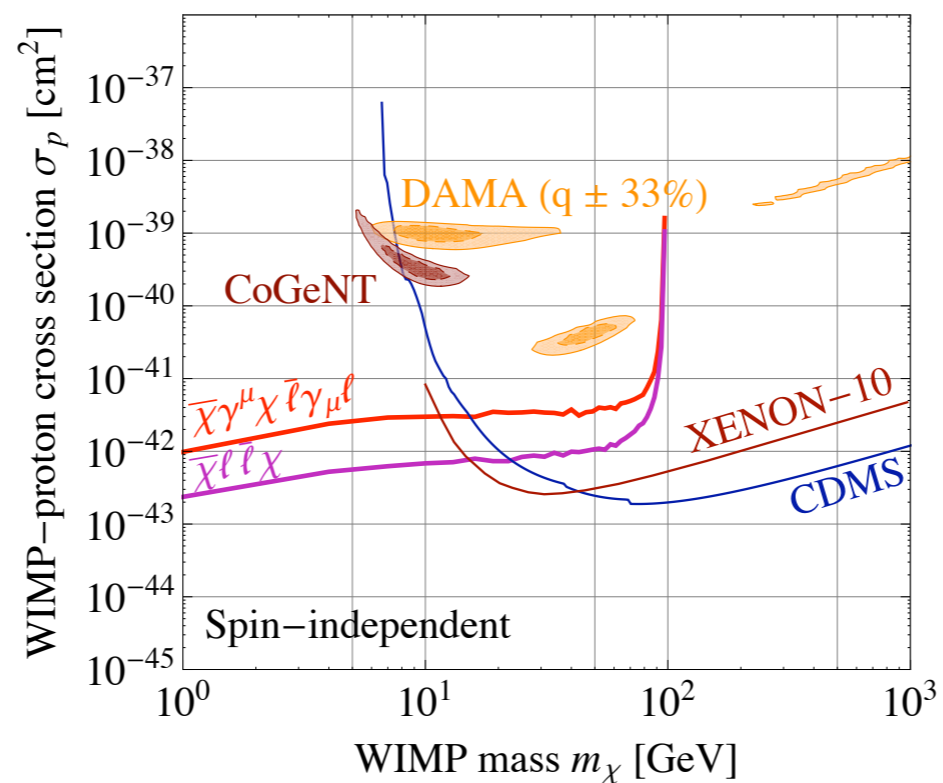
Equal couplings to all SM fermions



Equal couplings to all SM fermions



Couplings to leptons only



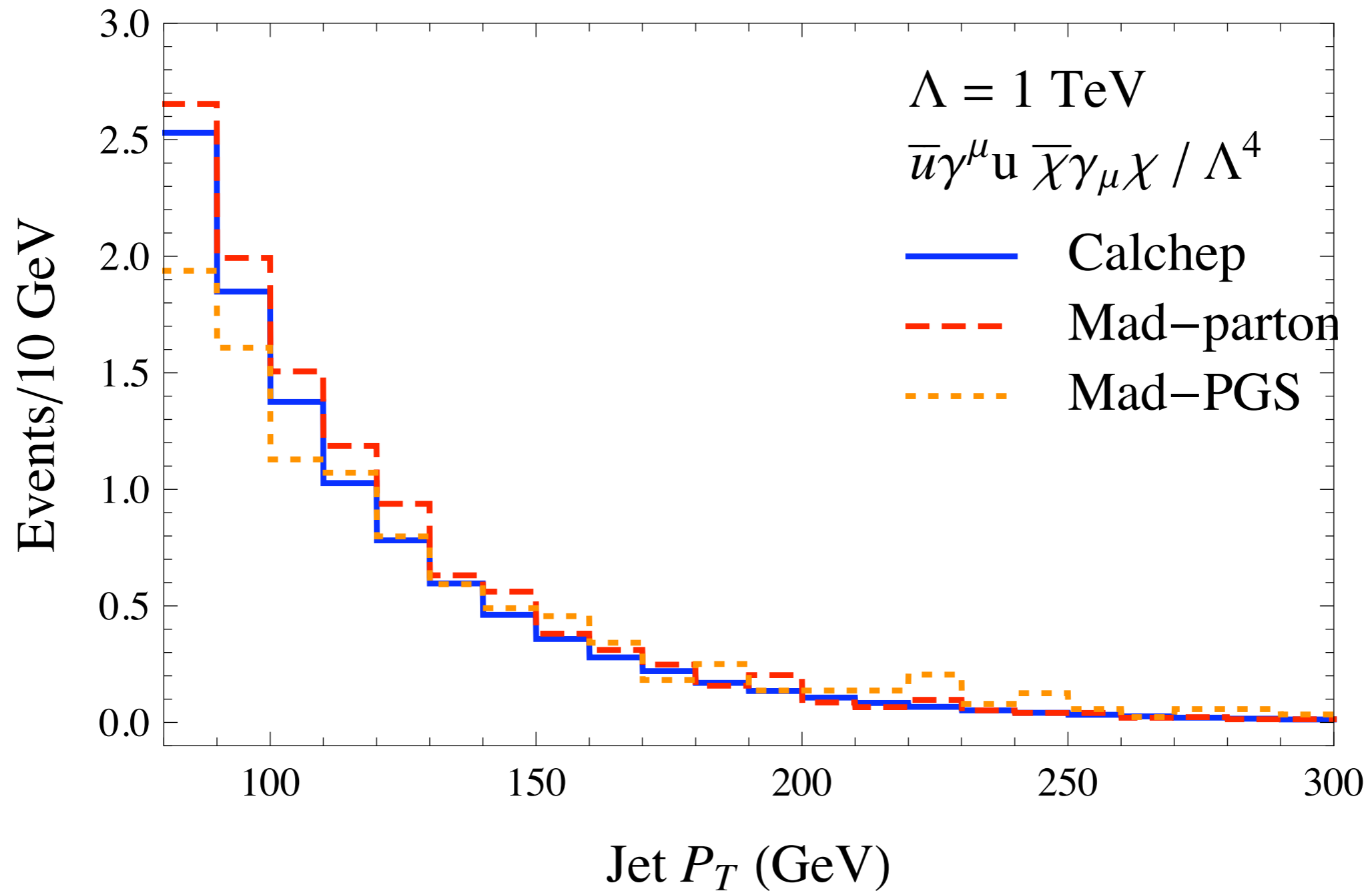
# Conclusions

- Mono-jet searches at the Tevatron already place strong constraints on dark matter
- Competitive with direct detection searches
  - Light DM
  - Spin dependent
  - Non-standard DM e.g. iDM, exoDM, MDDDM
- Independent of all astrophysics uncertainties
- Shape information, reduce theory errors,...
- Light mediators weaken collider bounds
- If we see a DD signal in a region ruled out by colliders we have discovered 2 particles

**Mono-jet + mono-photon analyses important**

# Backup Slides





# ADD analysis

Phys.Rev.Lett.101:181602,2008

arXiv:0807.3132

## Data Selection:

- Central Photon  $E_T > 50$  GeV
- Missing  $E_T > 50$  GeV
- No jets with  $E_T > 15$  GeV
- No tracks with  $P_t > 10$  GeV
- At least 3 low  $P_t$  COT tracks

## Background Predictions:

CDF RunII Preliminary, $2.0 \text{ fb}^{-1}$		
Channel	$\gamma E_T > 50 \text{ GeV}$	$\gamma E_T > 90 \text{ GeV}$
$W \rightarrow e \rightarrow \gamma$	$47.3 \pm 5.1$	$2.6 \pm 0.4$
$W \rightarrow \mu/\tau \rightarrow \gamma$	$19.1 \pm 4.2$	$1.0 \pm 0.2$
$W\gamma \rightarrow \mu\gamma \rightarrow \gamma$	$33.1 \pm 10.2$	$1.7 \pm 1.2$
$W\gamma \rightarrow e\gamma \rightarrow \gamma$	$8.0 \pm 3.0$	$0.8 \pm 0.7$
$W\gamma \rightarrow \tau\gamma \rightarrow \gamma$	$17.6 \pm 1.6$	$2.5 \pm 0.2$
$\gamma\gamma \rightarrow \gamma$	$18.9 \pm 2.3$	$2.3 \pm 0.6$
cosmics	$36.4 \pm 2.5$	$9.8 \pm 1.3$
$Z\gamma \rightarrow \nu\nu\gamma$	$99.7 \pm 9.5$	$25.2 \pm 2.8$
Total	$280.1 \pm 15.7$	$46.7 \pm 3.0$
Data	280	40

## • Optimized Search for LED:

- Leading Jet  $E_T > 150$  GeV
- Event Missing  $E_T > 120$  GeV
- Allow 2nd Jet with  $E_T < 60$  GeV
- No 3rd Jet with  $E_T > 20$  GeV

## • Results:

- Background Predictions:

Background	Number of Events
Z $\rightarrow$ nu nu	390 +/- 30
W $\rightarrow$ tau nu	187 +/- 14
W $\rightarrow$ mu nu	117 +/- 9
W $\rightarrow$ e nu	58 +/- 4
Z $\rightarrow$ ll	6 +/- 1
QCD	23 +/- 20
Gamma plus Jet	17 +/- 5
Non-Collision	10 +/- 10
Total Predicted	808 +/- 62
Data Observed	809