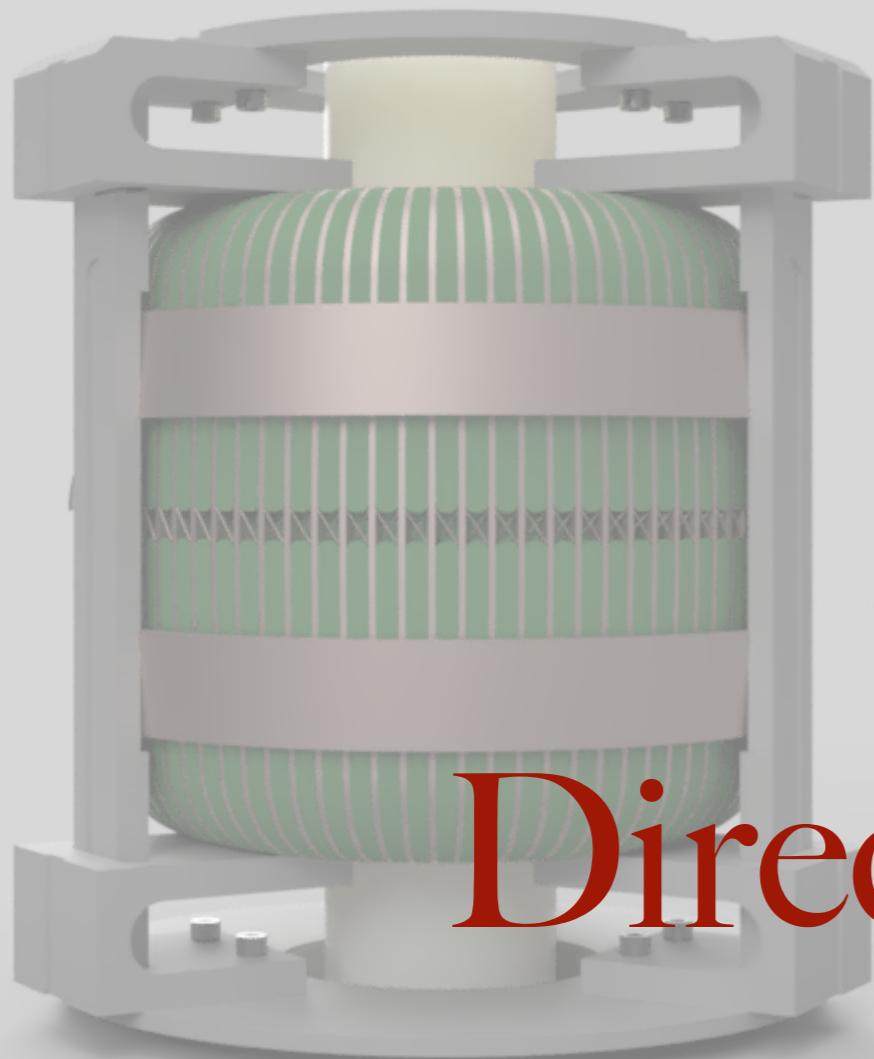
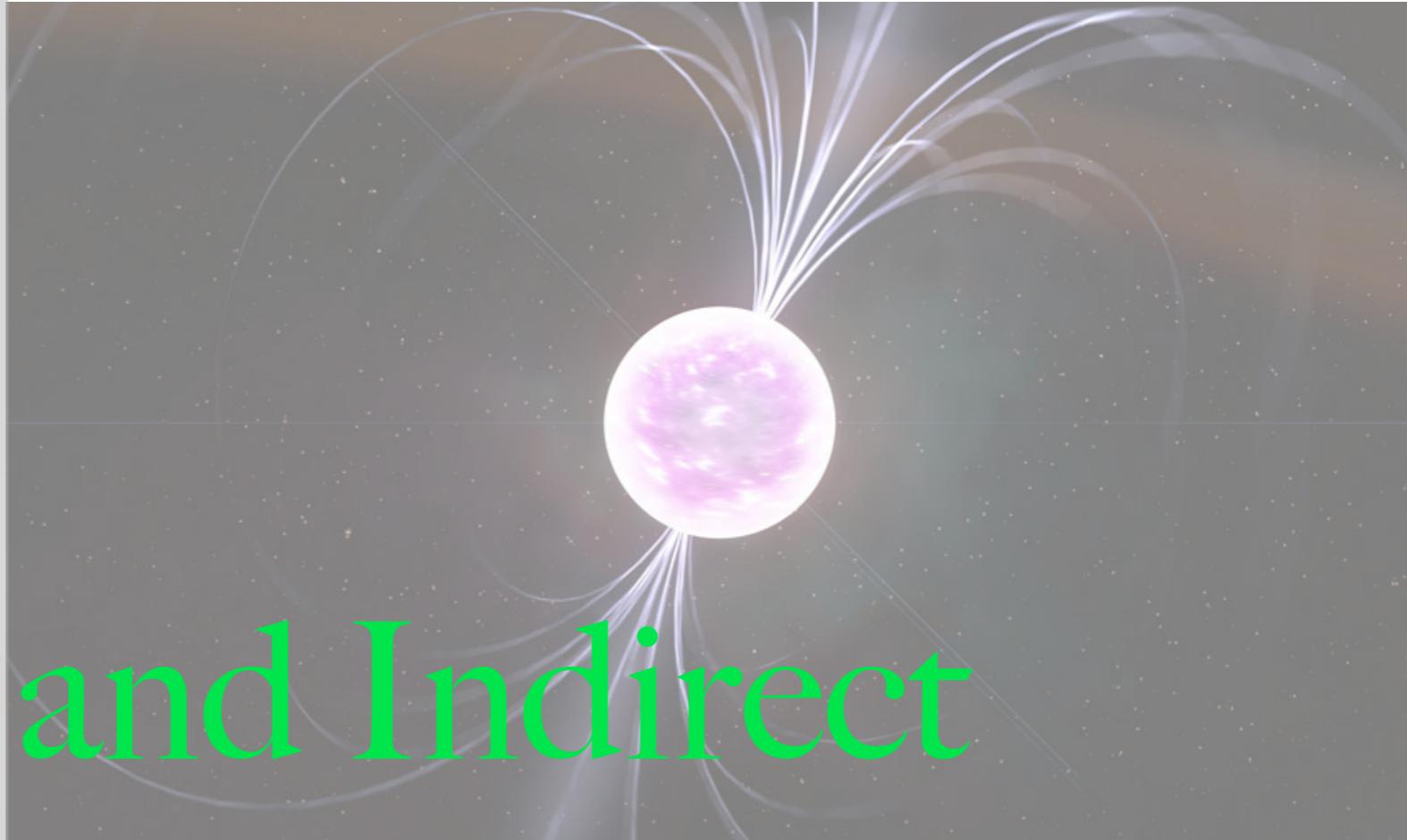


Axion Dark Matter Detection:



Direct and Indirect



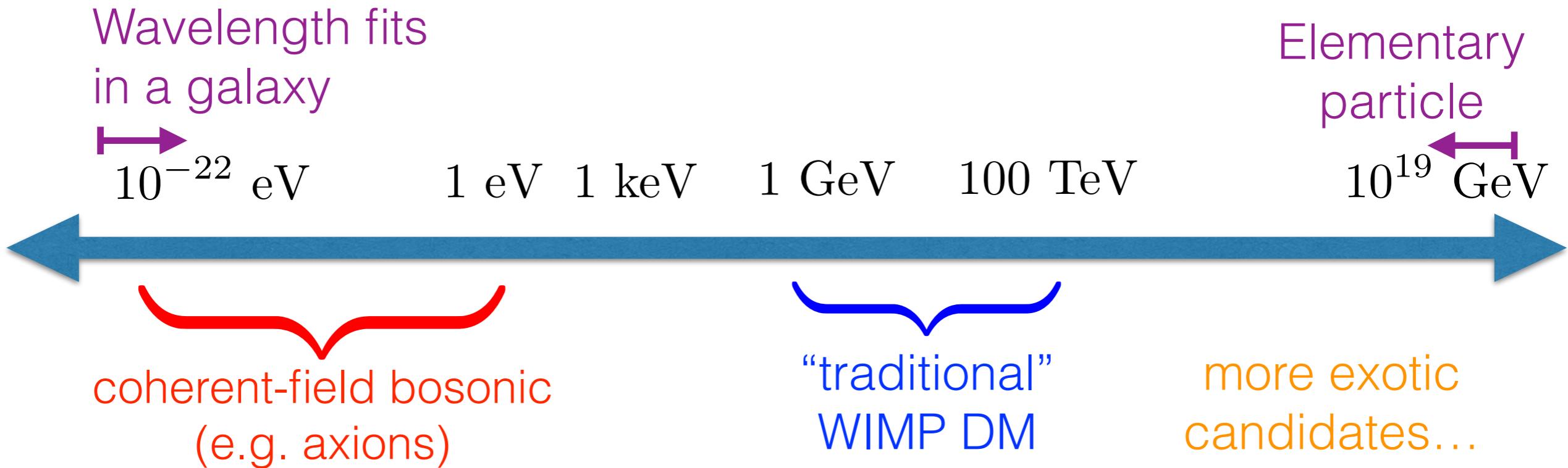
Kavli Institute
for Cosmological Physics
at The University of Chicago

Yoni Kahn
KICP/UIUC

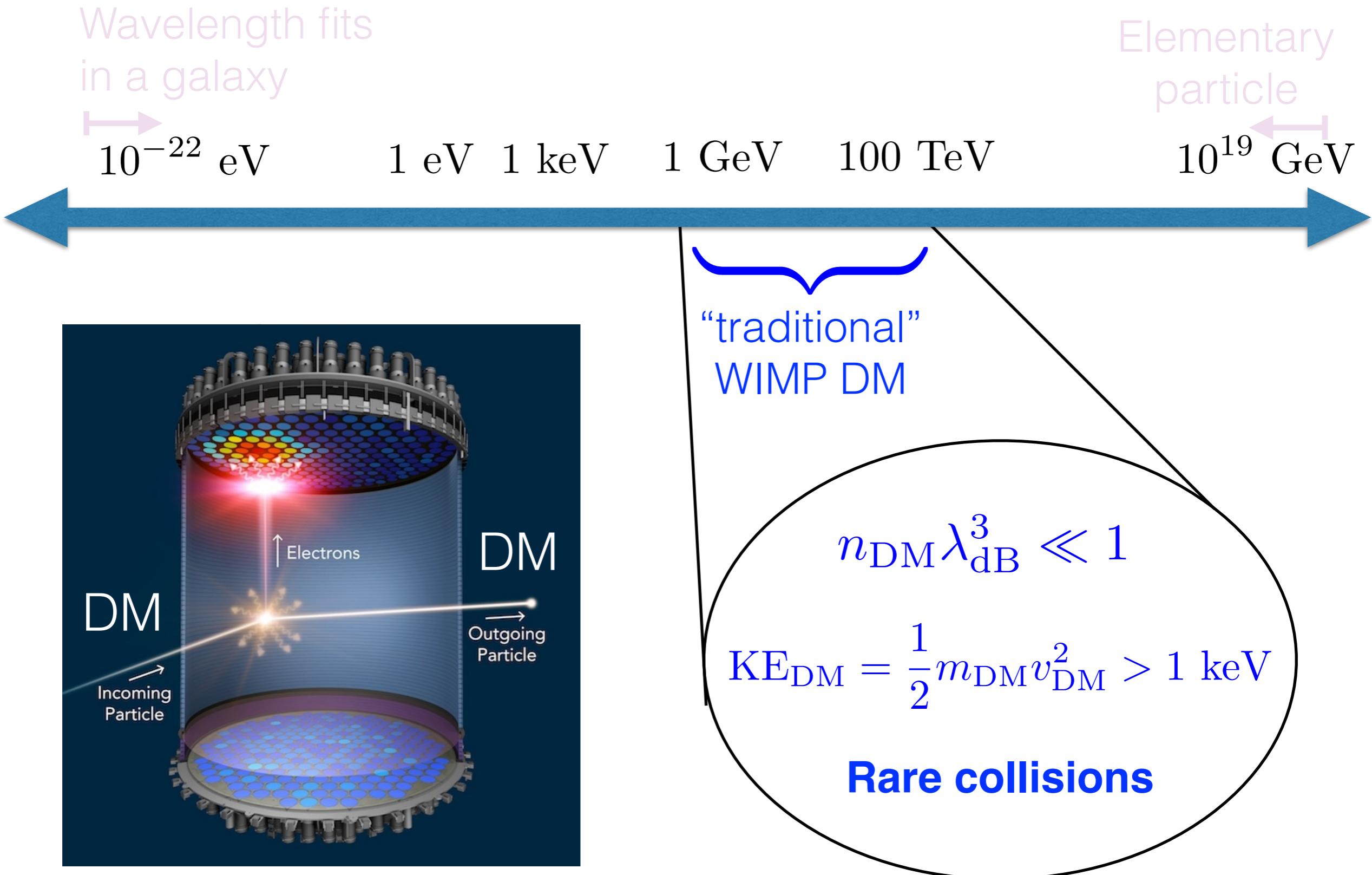


TRIUMF workshop “Signals of Dark Matter in its Natural Habitat”

50 orders of magnitude



50 orders of magnitude



50 orders of magnitude

Wavelength fits
in a galaxy



10^{-22} eV

1 eV

1 keV

1 GeV

100 TeV

10^{19} GeV

Elementary
particle

coherent-field bosonic
(e.g. axions)

“traditional”
WIMP DM

$$n_{\text{DM}} \lambda_{\text{dB}}^3 \gg 1$$
$$\text{KE}_{\text{DM}} = \frac{1}{2} m_{\text{DM}} v_{\text{DM}}^2 \ll 1 \mu\text{eV}$$

Behaves as classical field

New paradigm for
dark matter detection!
Focus of this talk

Axion DM: here and now

$$a(\mathbf{x}, t) = \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t + \mathcal{O}(v_{\text{DM}})\mathbf{x})$$

amplitude set by local DM density oscillates at frequency set by DM mass

e.g. $m_a = 10^{-9}$ eV
 $\lambda_{\text{Comp}} \sim$ km
 $\tau_{\text{Comp}} \sim \mu\text{s}$

Local DM velocity \rightarrow Spatial coherence \rightarrow Temporal coherence

$$\Delta v_{\text{DM}} \sim v_{\text{DM}} \sim 10^{-3}$$

$$\lambda_{\text{dB}} = \frac{\lambda_{\text{Comp}}}{v_{\text{DM}}}$$

$$\tau_{\text{coh}} = \frac{\tau_{\text{Comp}}}{v_{\text{DM}}^2}$$

Classical physics is fine: $m_a = 10^{-9}$ eV $\implies N_a \sim 10^{18}/\text{cm}^3$

What should we measure?

$$a(\mathbf{x}, t) = \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t + \mathcal{O}(v_{\text{DM}})\mathbf{x})$$

In axion DM background, get oscillating observables:

$$\nabla \times \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \left(\mathbf{E}_0 \times \nabla a - \mathbf{B}_0 \frac{\partial a}{\partial t} \right)$$
$$\nabla \cdot \mathbf{E}_a = -g_{a\gamma\gamma} \mathbf{B}_0 \cdot \nabla a$$

Oscillating response from static fields

$$H_N \supset g_{aNN} \nabla a \cdot \vec{\sigma}_N$$

Spin-dependent force

$$d_n = g_d a$$

Time-varying EDM

Note: $\nabla a \sim v_{\text{DM}} \sim 10^{-3}$ so some are easier than others

Axion direct detection

$$\nabla \times \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \left(\mathbf{E}_0 \times \cancel{\nabla} a - \mathbf{B}_0 \frac{\partial a}{\partial t} \right)$$

Goal: detect axion DM **on Earth**
through interactions with laboratory B-fields

Axion searches with magnetic fields

$$\nabla \times \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

Cavity regime: $\lambda_{\text{Comp}} \sim R_{\text{exp}}$
e.g. ADMX

$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

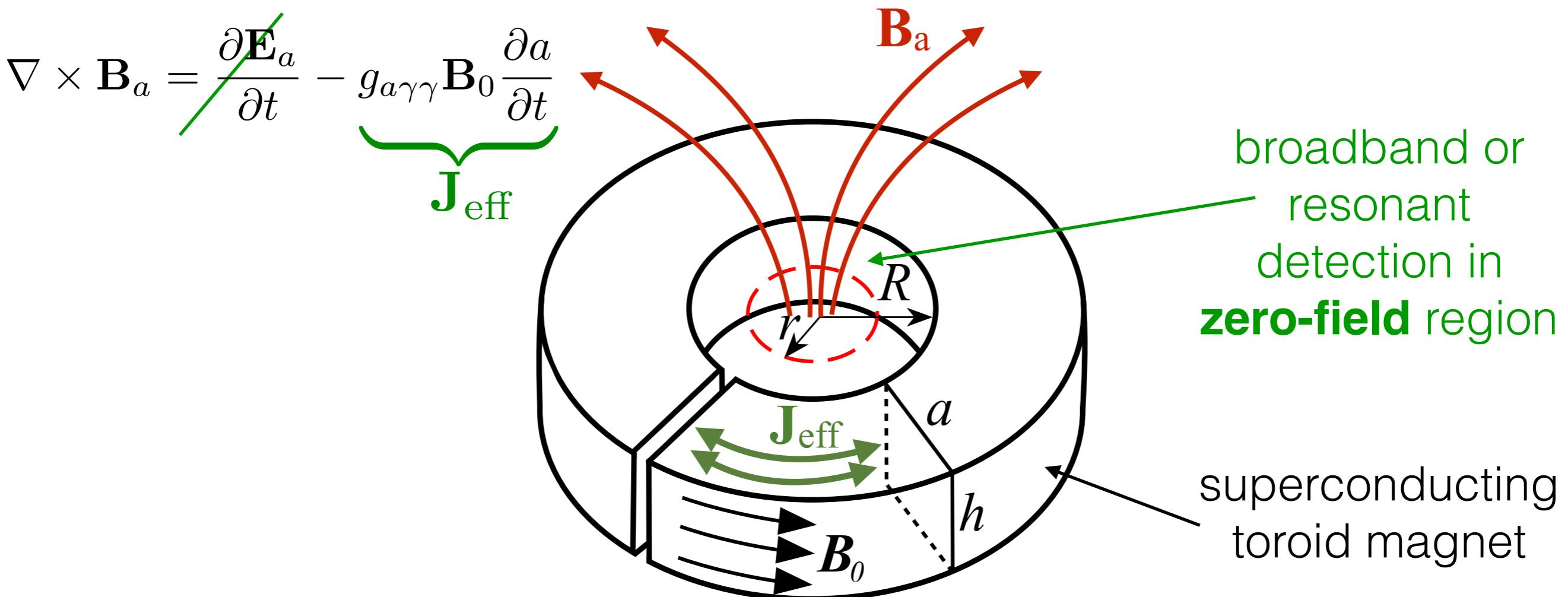
Quasistatic regime: $\lambda_{\text{Comp}} \gg R_{\text{exp}}$
e.g. ABRACADBRA

$$\nabla \times \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

Radiation regime: $\lambda_{\text{Comp}} \ll R_{\text{exp}}$
e.g. MADMAX

Quasistatic regime: ABRACADABRA

A Broadband/Resonant Approach to Cosmic Axion Detection
with an Amplifying B-field Ring Apparatus



$$\langle \Phi(t)^2 \rangle \sim g_{a\gamma\gamma}^2 \rho_{\text{DM}} B_0^2 V^{5/3}$$

Volume enhancement at low masses



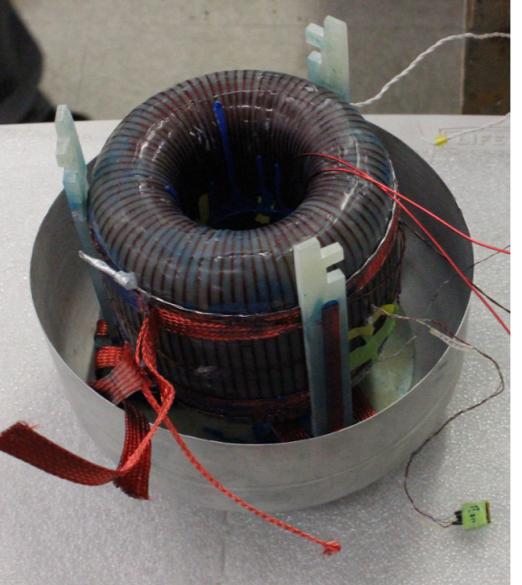
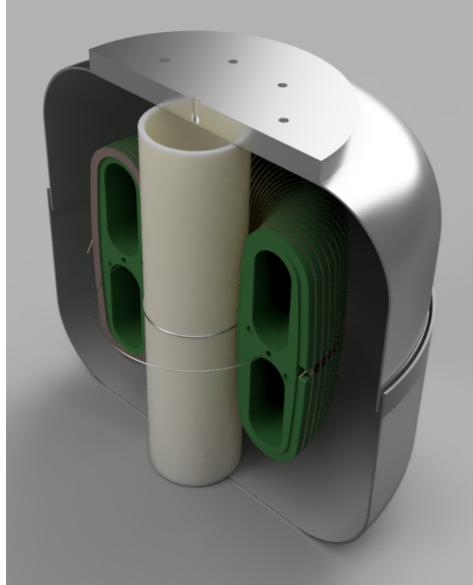
ABRACADABRA-10cm:

hot off the presses

MIT

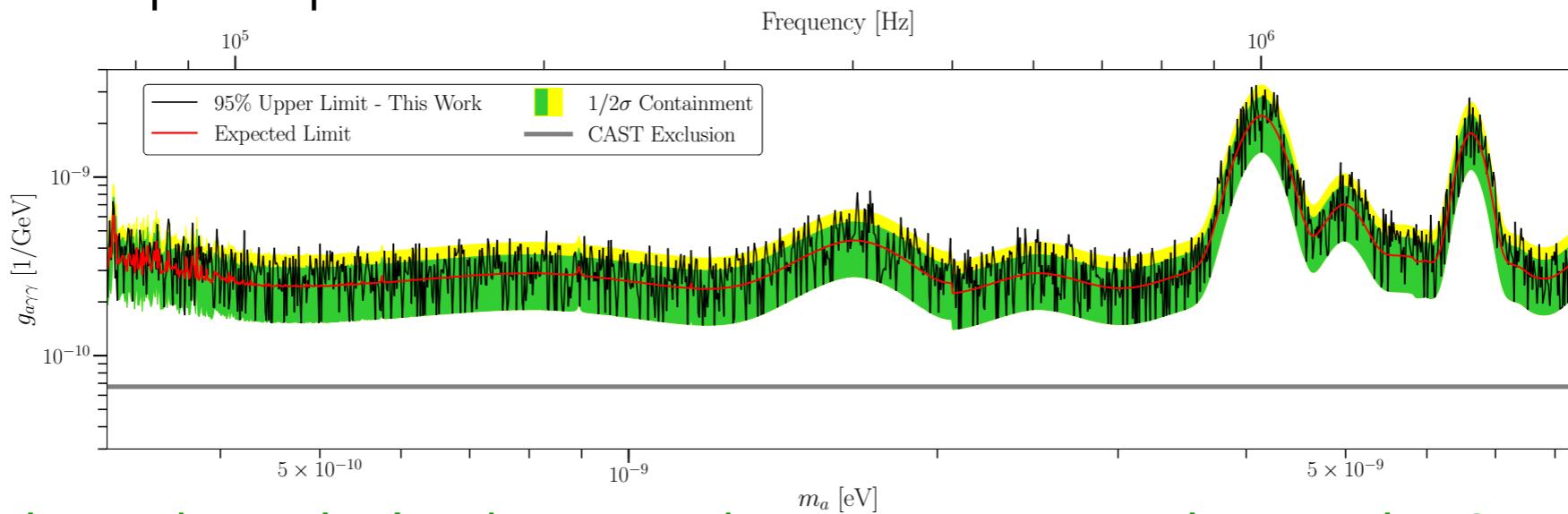
THE UNIVERSITY
of NORTH CAROLINA
at CHAPEL HILL

MICHIGAN



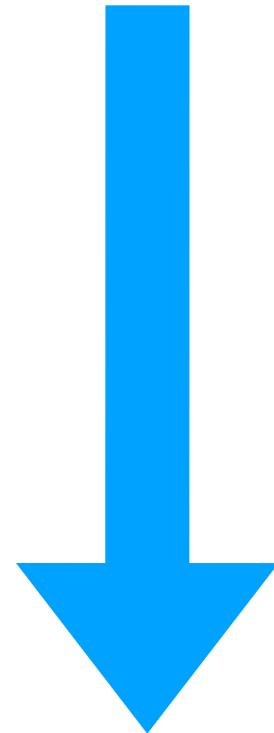
superconducting magnet
and pickup loop

inside a 150 mK
dilution refrigerator

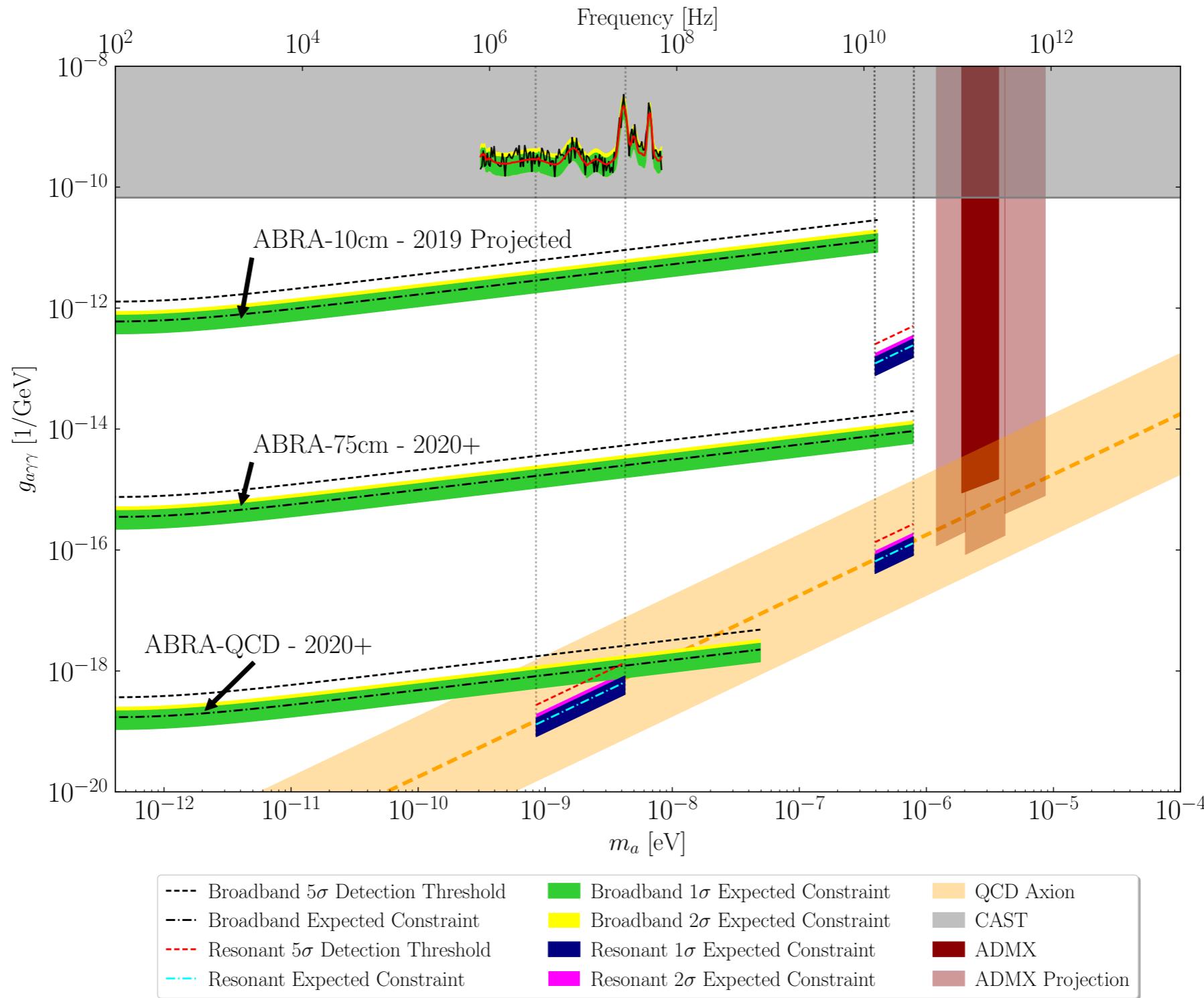


This is a theorist's dream: theory to experiment in 2 years flat!

Future ABRA reach



Bigger
magnets,
better
shielding,
upgraded
readout



Axion indirect detection

Goal: detect photons produced by
axions interacting with **astrophysical** B-fields

Axion-photon mixing

$a(t) = a_0 e^{i(kz - \omega t)}$
 axion wavepacket

$B_t(z)$
 $\mathbf{B}(z)$
 E
 $A_{\parallel}(t) \propto e^{i(k' z - \omega t)}$

$$-\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} \implies$$

$$\left[\omega^2 + \partial_z^2 + \begin{pmatrix} -\omega_p^2 & -g_{a\gamma\gamma} B_t \omega \\ -g_{a\gamma\gamma} B_t \omega & -m_a^2 \end{pmatrix} \right] \begin{pmatrix} A_{\parallel} \\ a \end{pmatrix} = 0$$

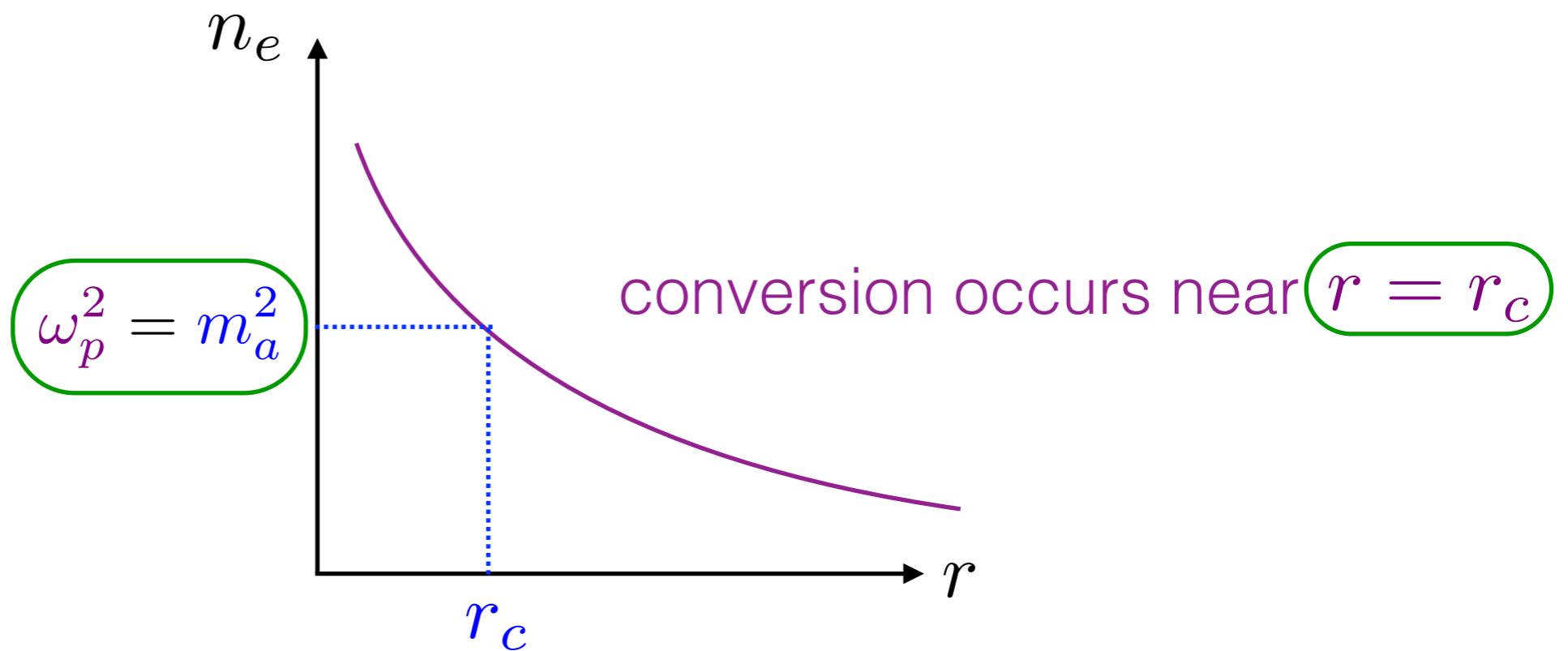
plasma frequency
 only transverse field contributes
 only mixes with one polarization

Resonant conversion

Need momentum matching

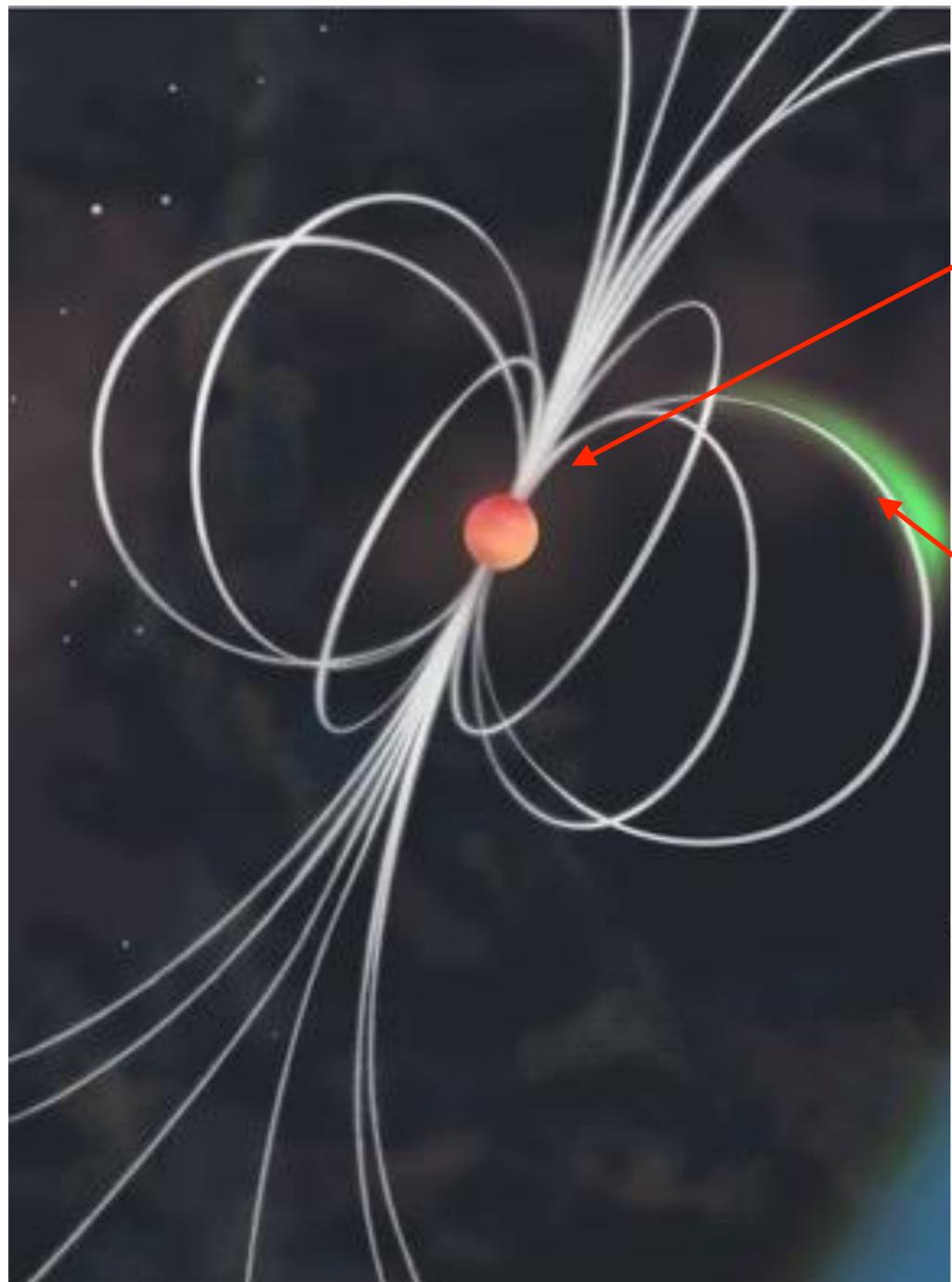
Option 1: B-field has significant spatial variations at axion wavelength (hard to obtain: $m_a = 10^{-6}$ eV $\Rightarrow \lambda_a \sim$ m)

Option 2: B-field is approximately homogeneous, photon dispersion changes with plasma density



Neutron stars: ideal candidates!

2 key (related) ingredients for axion indirect detection:



1. Strong B-fields

$$B_\theta = \frac{B_0}{2} \left(\frac{r_{\text{NS}}}{r} \right)^3 \sin \theta$$

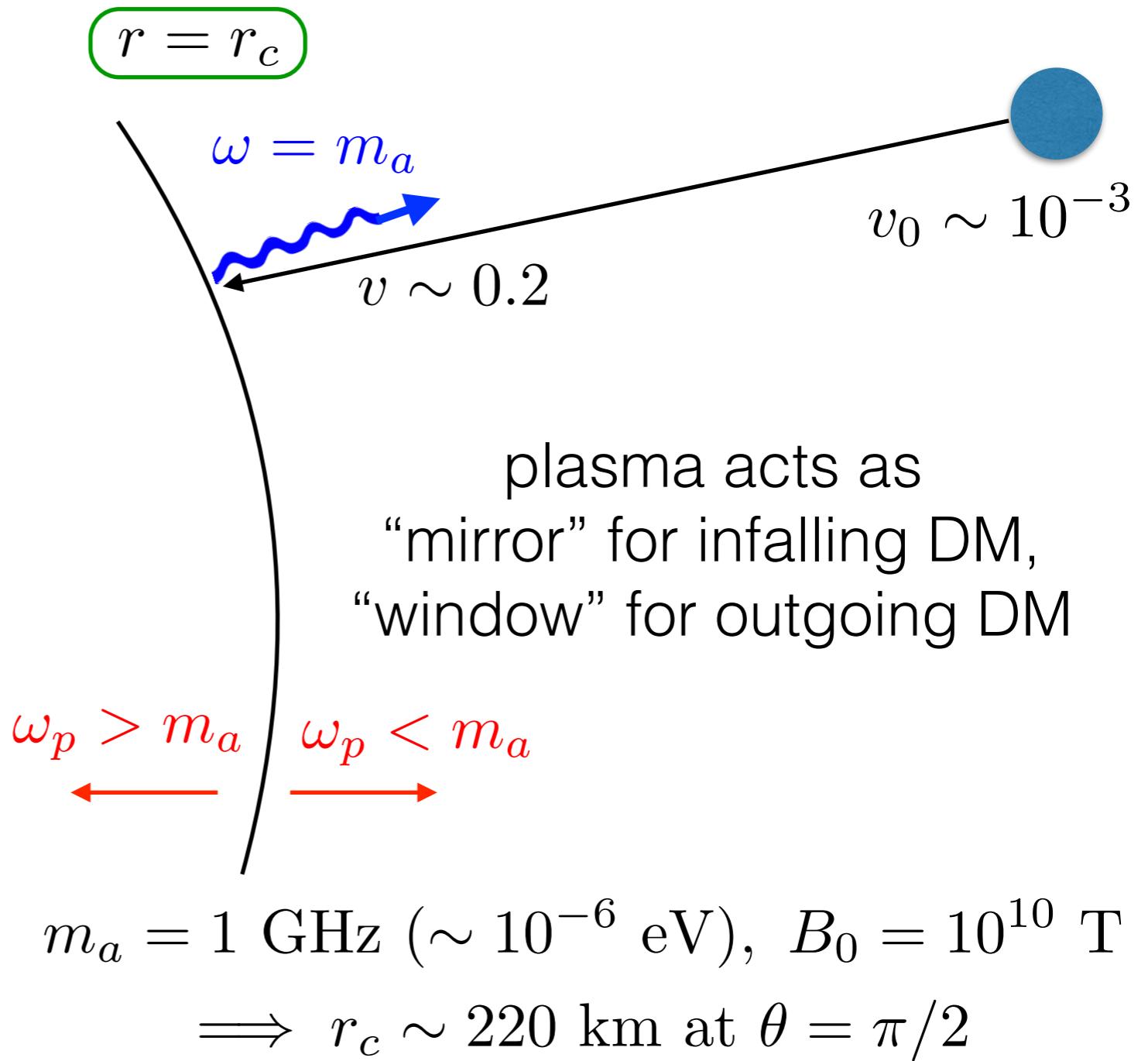
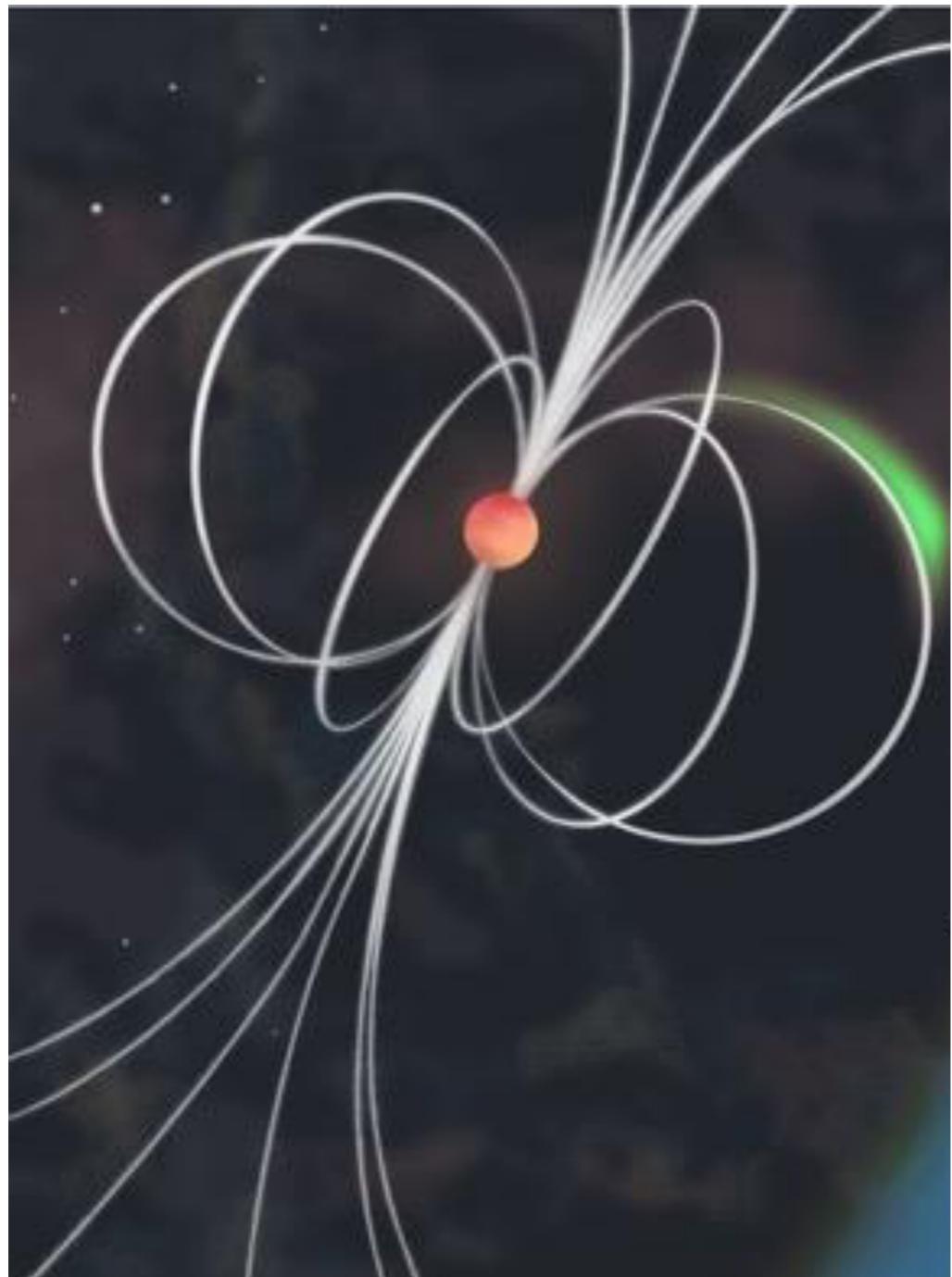
$$B_0 \sim 10^{10} \text{ T}$$

2. Goldreich-Julian model relates plasma frequency in “lobes” to dipole B-field:

$$\omega_p \propto \sqrt{n_e} \propto \sqrt{B_0 \left(\frac{r_{\text{NS}}}{r} \right)^3 (3 \cos^2 \theta - 1)}$$

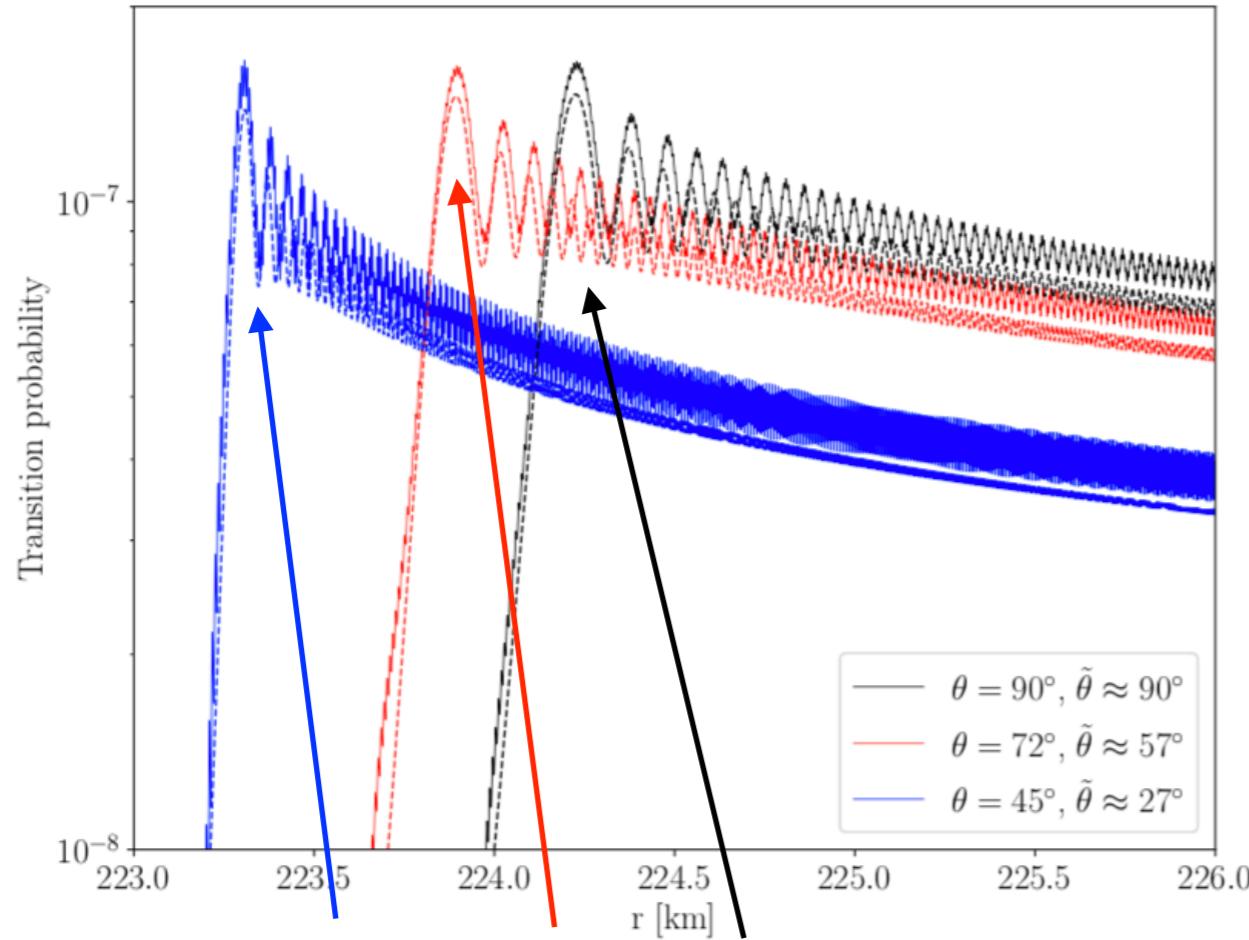
Monotonically decreasing,
can always solve $\omega_p = m_a$.

Infalling axion DM conversion



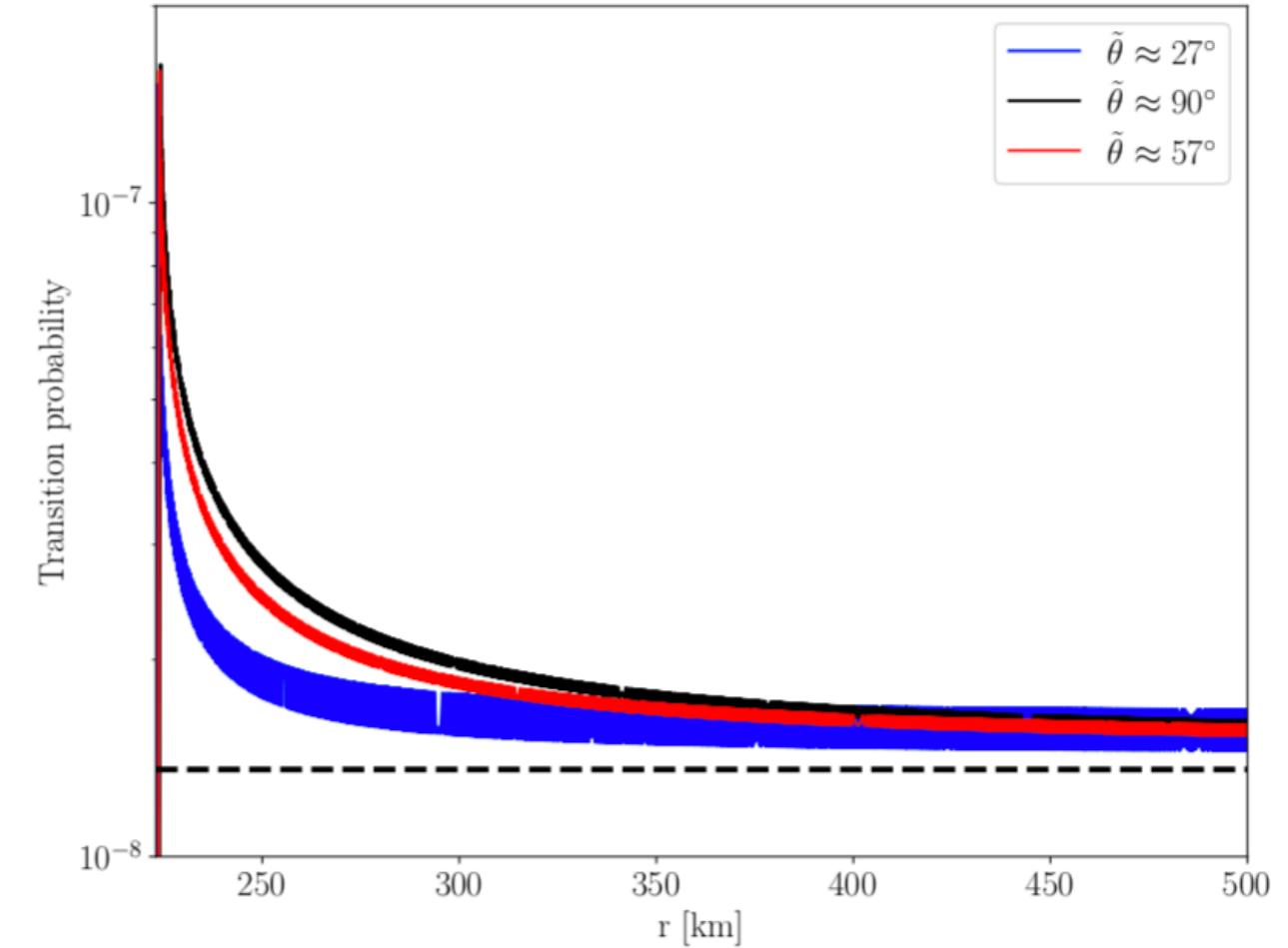
gravitational acceleration + Liouville = enhanced DM density

Photon transition probability



transition prob. peaks at conversion radius, happens over distance

$$L = \sqrt{\frac{2\pi r_c v_c}{3m_a}}$$



outgoing photon wave **damped by plasma** (like ocean waves):

$$p_{a\gamma}^{\infty} \approx \frac{1}{2v_c} g_{a\gamma\gamma}^2 B(r_c)^2 L^2$$

Expected photon flux

$$\frac{d\mathcal{P}}{d\Omega} \approx 2 \times p_{a\gamma}^\infty \rho_{\text{DM}}^{r_c} v_c r_c^2$$

incoming+outgoing

$$= \rho_{\text{DM}}^\infty \frac{2}{\sqrt{\pi}} \frac{1}{v_0} \sqrt{\frac{GM}{r_c}} \approx \sqrt{\frac{2GM_{\text{NS}}}{r_c}}$$

$$\frac{d\mathcal{P}(\theta, \theta_m, t)}{d\Omega} = \frac{d\mathcal{P}(\theta = \frac{\pi}{2}, \theta_m = 0)}{d\Omega} \times \frac{3 (\hat{\mathbf{m}} \cdot \hat{\mathbf{r}})^2 + 1}{|3 \cos \theta \hat{\mathbf{m}} \cdot \hat{\mathbf{r}} - \cos \theta_m|^{4/3}}$$

time-dependent! misalignment angle

Total power $\sim 10^{10}$ W for QCD axion, local DM density

Radio bump hunt

$$\frac{F}{B} \sim \frac{P}{4\pi d^2} \frac{1}{10^{-6}\omega}$$

$$S_{\min} = \text{SNR}_{\min} \frac{2 \text{ Jy}}{\sqrt{n_{\text{pol}} B \Delta t_{\text{obs}}}}$$

\downarrow

$2 \text{ Jy} \sim \frac{2 \times 10^{15} \text{ W}}{\text{kHz kpc}^2}$

\nearrow

$\sim 10^{-6}\omega$

SEFD



[Arecibo]

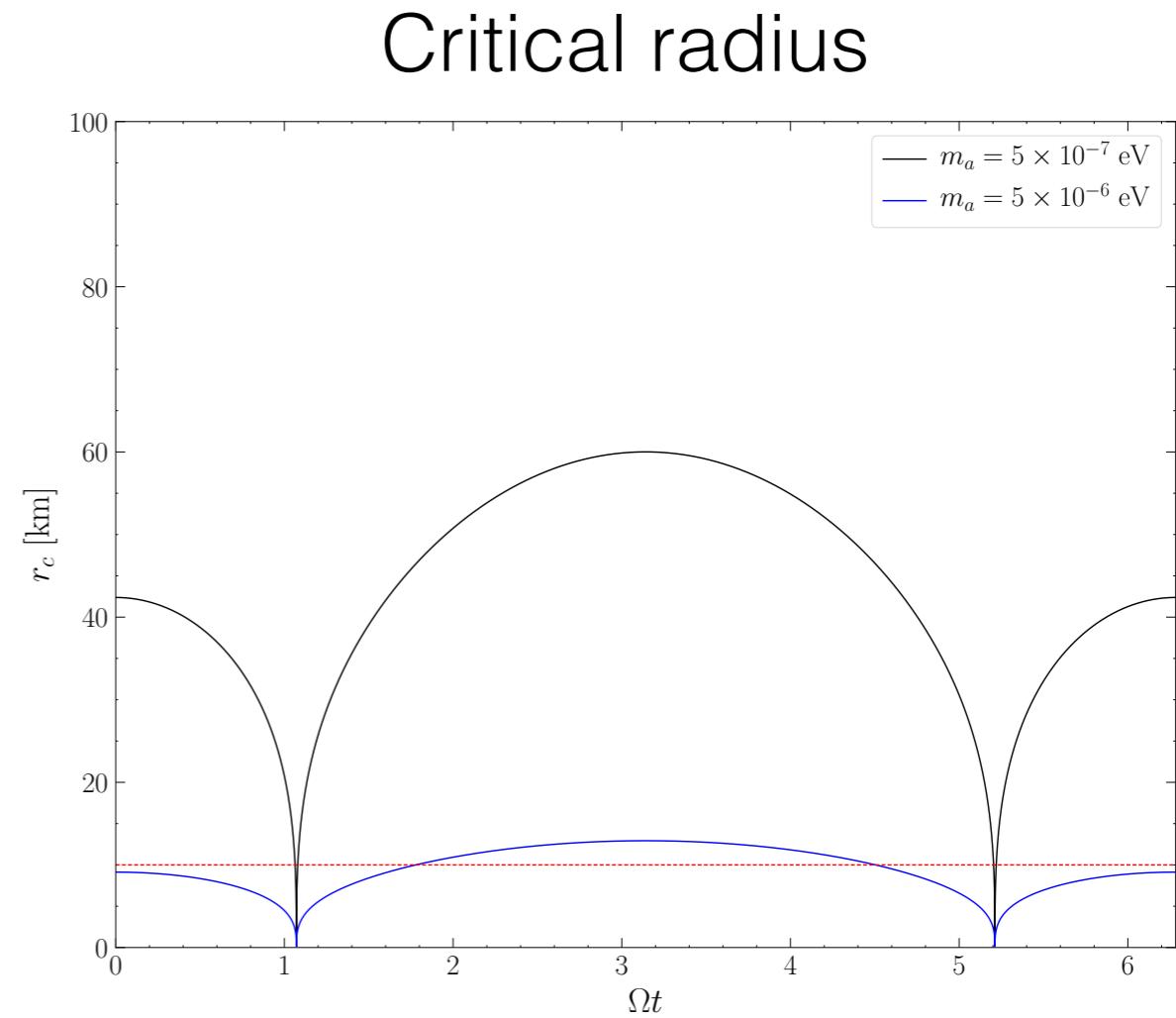
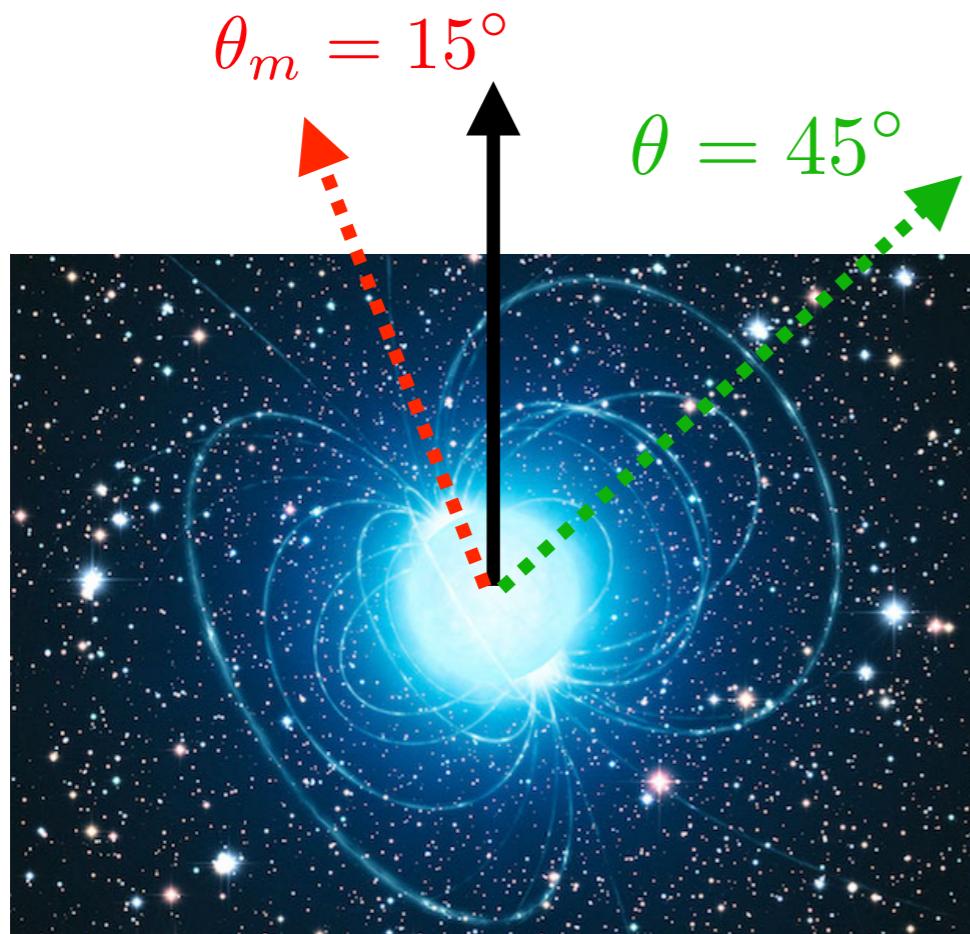
Desired characteristics:

- Radio-quiet (negligible foreground)
- Low DM velocity dispersion (dwarfs)
- Close by (< kpc), **or**
- DM-rich (Galactic center, dwarfs)

Added bonus: energy conservation keeps line narrow

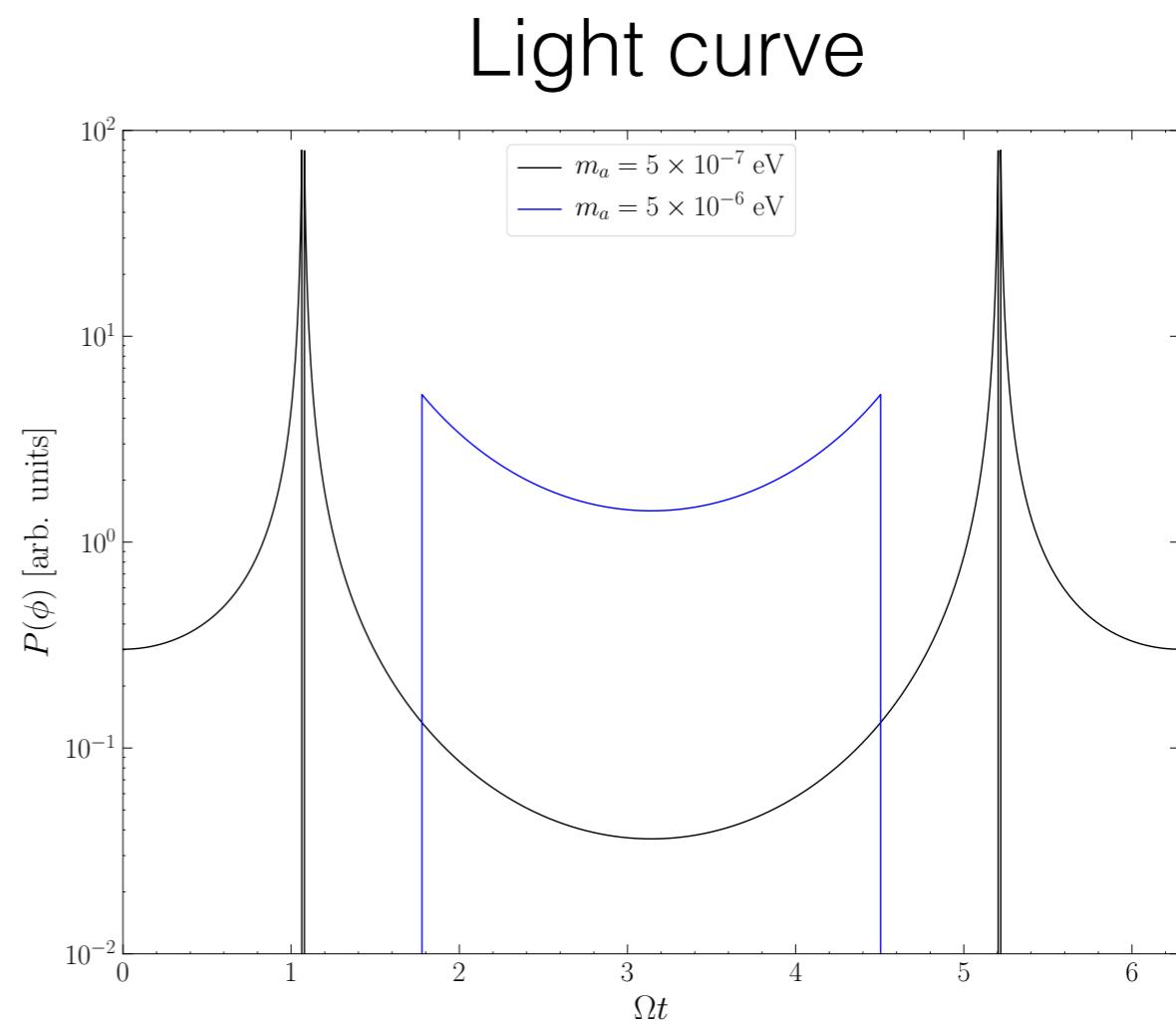
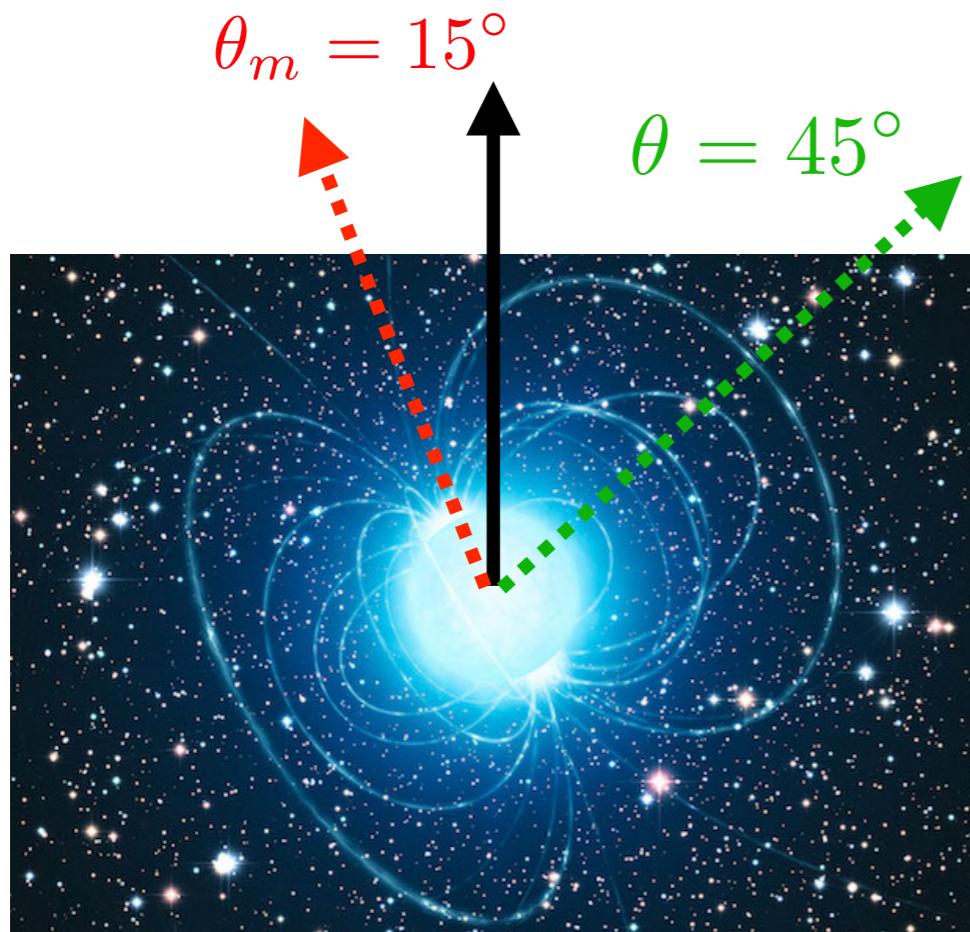
Time-dependent signal

Misaligned neutron stars have rotation axis misaligned from magnetic poles:
strong time dependence of plasma freq. in GJ model



Time-dependent signal

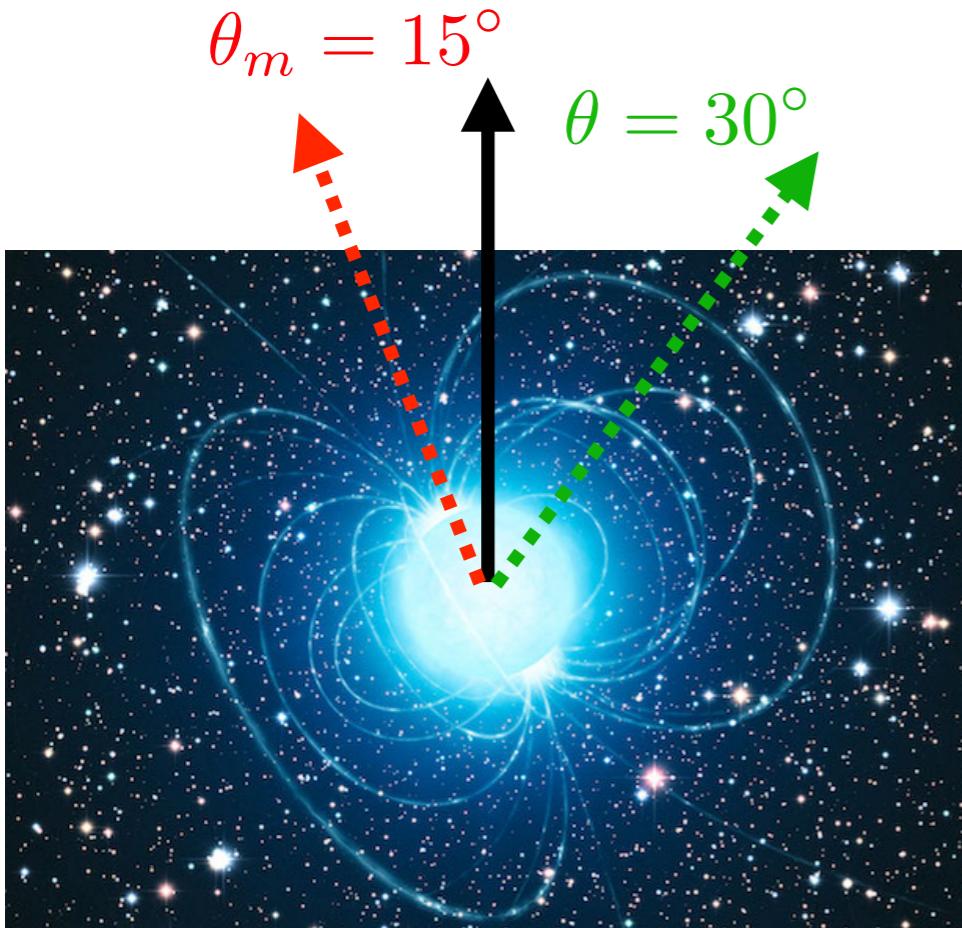
Misaligned neutron stars have rotation axis misaligned from magnetic poles:
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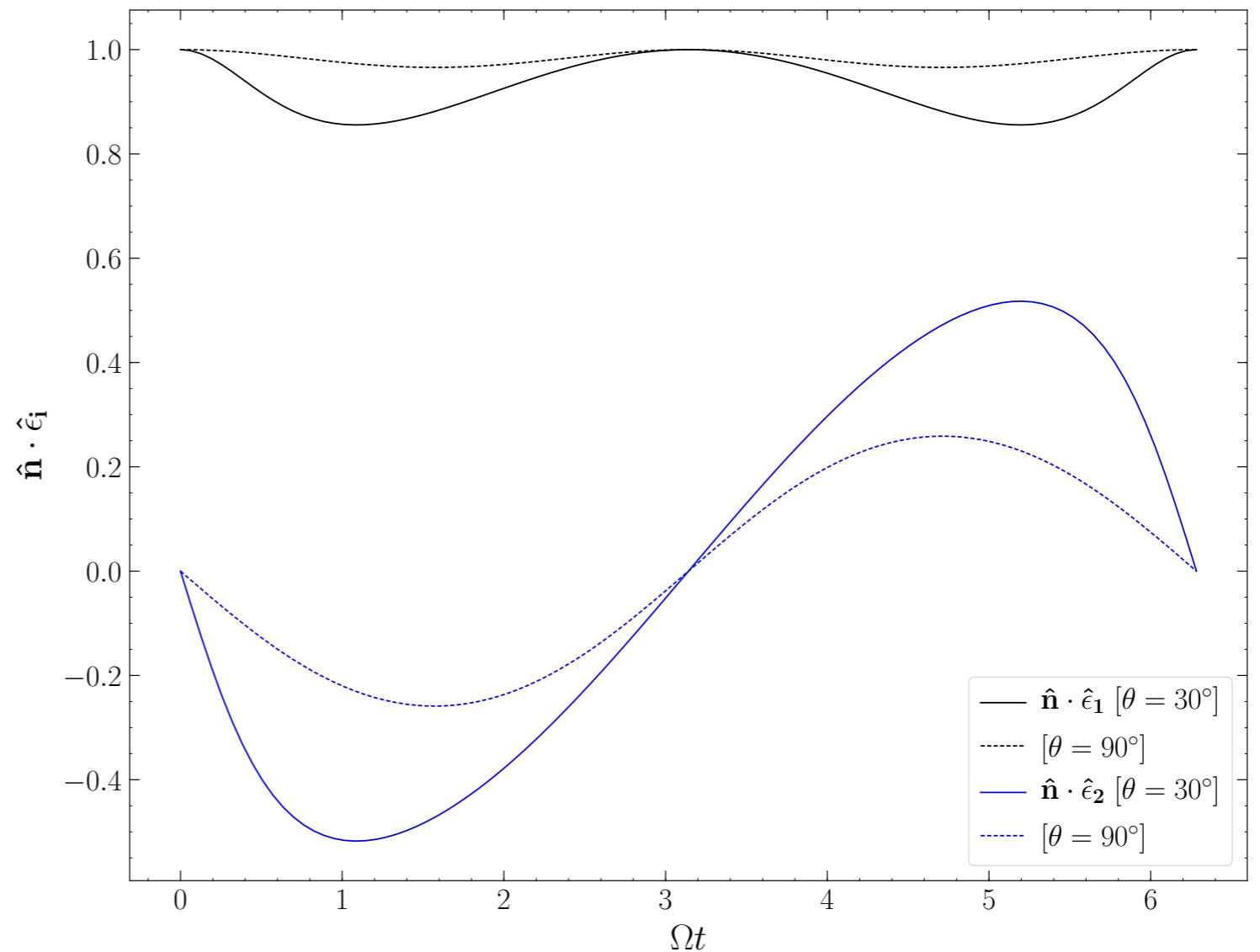
Striking signal: looks promising!

Polarization signal

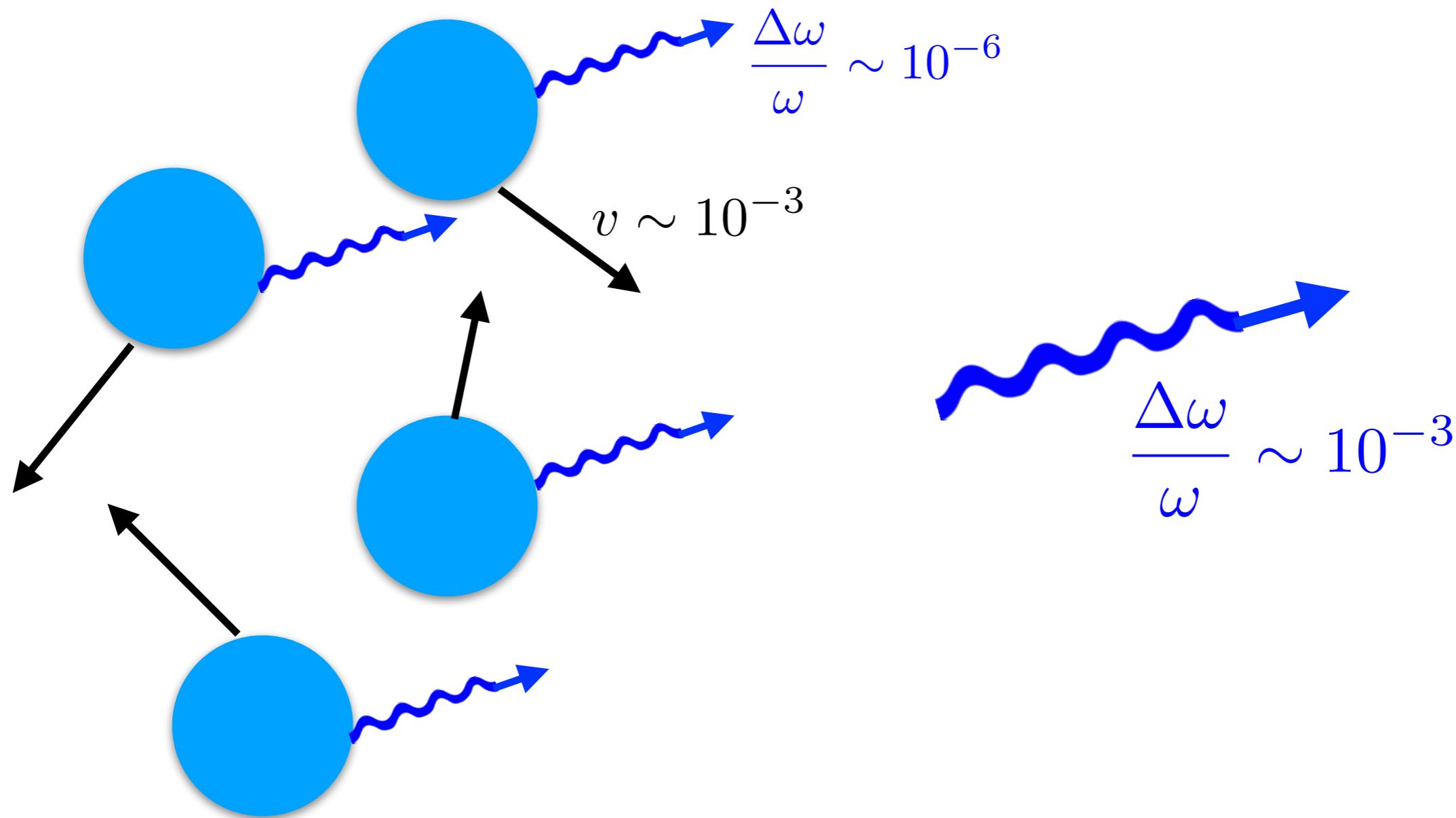
EM wave always polarized along
direction of B-field at conversion radius



$$(\hat{\epsilon}_1 = \hat{\theta}, \hat{\epsilon}_2 = \hat{\phi})$$



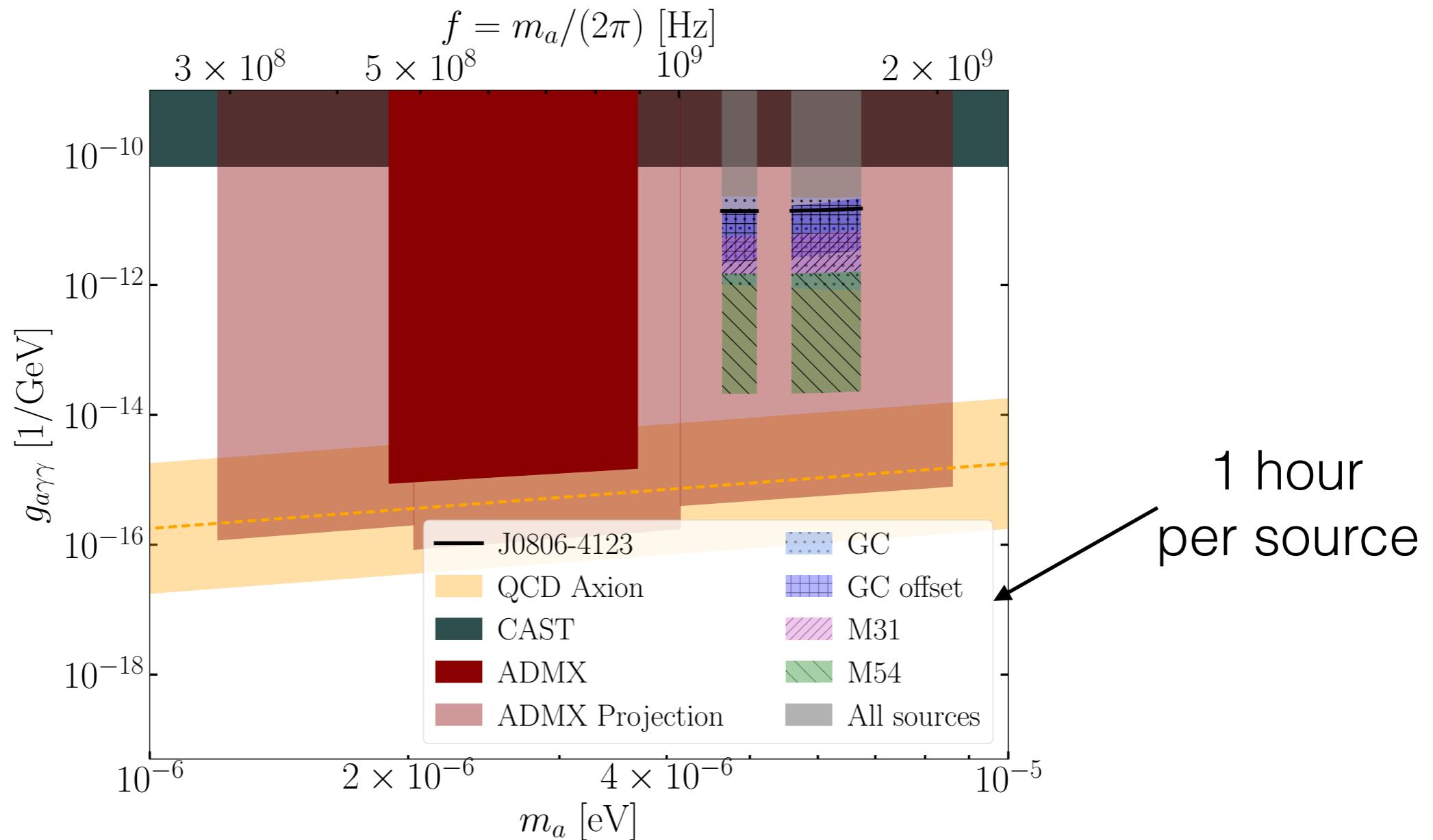
Neutron star populations



Doppler broadening: larger bandwidth
But signal adds incoherently!

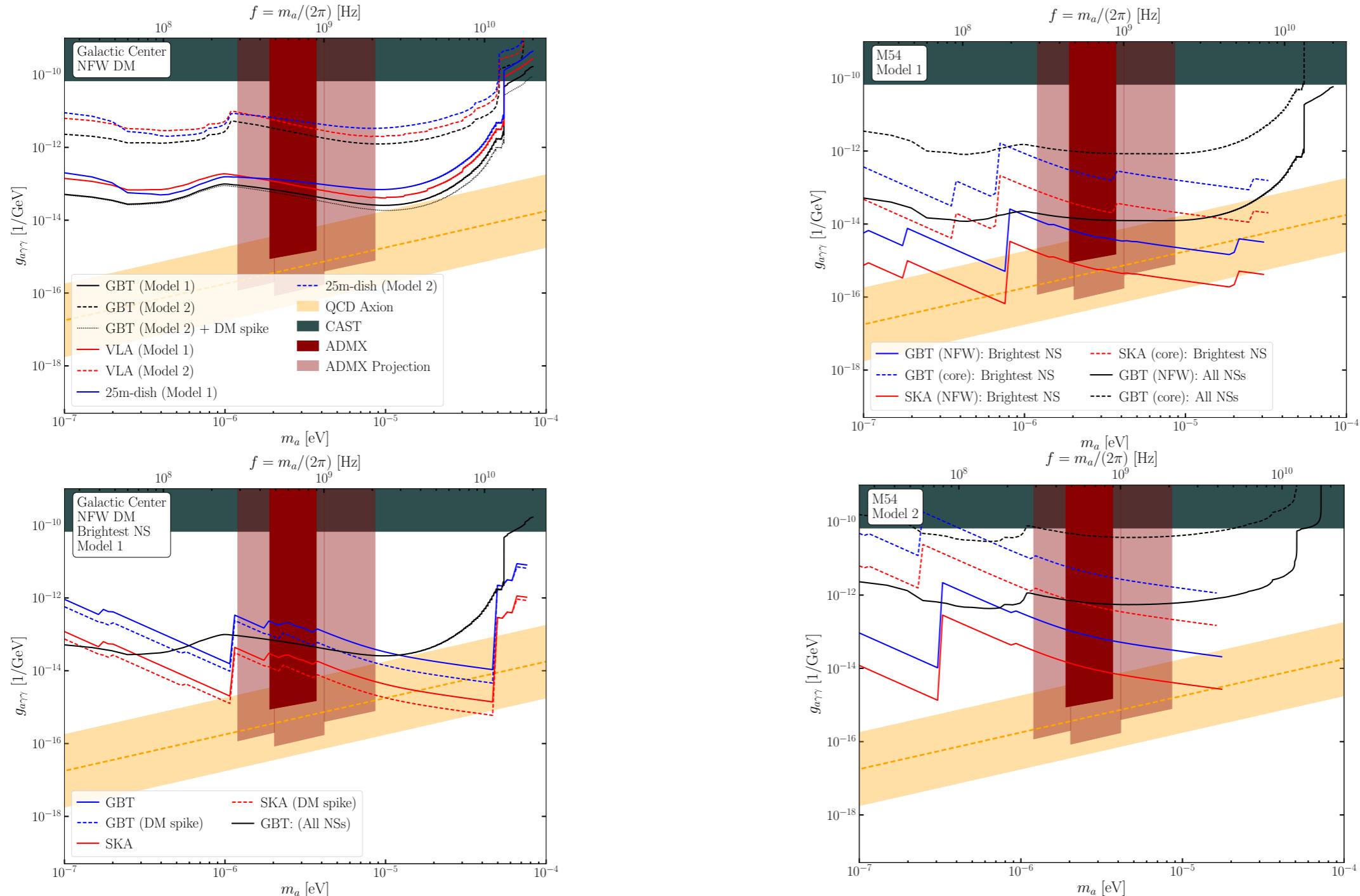
If $N_{NS} > 1000$, still win

NS populations at Green Bank Telescope



Can be competitive with ADMX with 1 hour of observation!

Future radio searches



QCD axion might be within our reach!